Cosmological detectors of UHF gravitational waves

Camilo A. Garcia Cely

Alexander von Humboldt Fellow



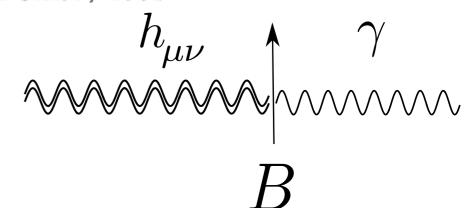
October 14, 2021

Revisiting Gertsenhstein's ideas

SOVIET PHYSICS JETP

VOLUME 14, NUMBER 1

JANUARY, 1962



WAVE RESONANCE OF LIGHT AND GRAVITIONAL WAVES

M. E. GERTSENSHTEĬN

Submitted to JETP editor July 29, 1960

J. Exptl. Theoret. Phys. (U.S.S.R.) 41, 113-114 (July, 1961)

The energy of gravitational waves excited during the propagation of light in a constant magnetic or electric field is estimated.

SOVIET PHYSICS JETP

VOLUME 16, NUMBER 2

FEBRUARY, 1963

ON THE DETECTION OF LOW FREQUENCY GRAVITATIONAL WAVES

M. E. GERTSENSHTEĬN and V. I. PUSTOVOĬT

Submitted to JETP editor March 3, 1962

J. Exptl. Theoret! Phys. (U.S.S.R.) 43, 605-607 (August, 1962)

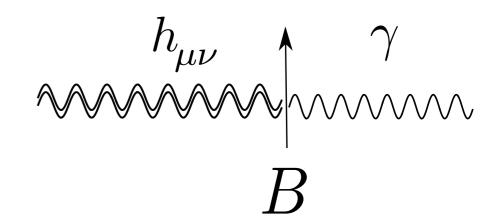
It is shown that the sensitivity of the electromechanical experiments for detecting gravitational waves by means of piezocrystals is ten orders of magnitude worse than that estimated by Weber. [1] In the low frequency range it should be possible to detect gravitational waves by the shift of the bands in an optical interferometer. The sensitivity of this method is investigated.

Terrestrial interferometers

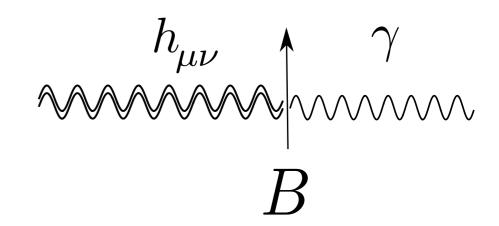


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- The process is strictly analogous to axionphoton conversion. Raffelt, Stodolski'89

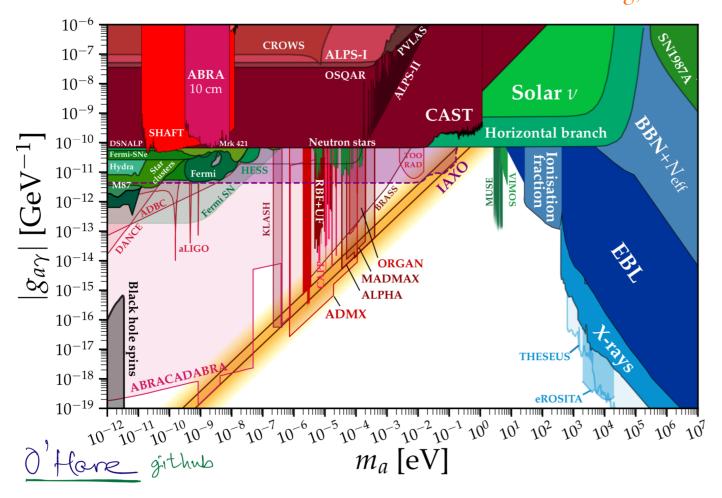


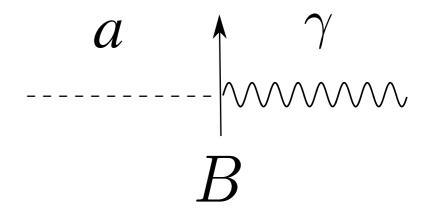
Axion dark matter

• The strong CP problem: experiments put a strong upper bound on the electric dipole of the neutron

$$\mathcal{L} = \frac{g_s^2 \theta}{32\pi^2} F_{\mu\nu}^a \tilde{F}^{a\mu\nu} \longrightarrow \frac{g_s^2 a}{32\pi^2 f} F_{\mu\nu}^a \tilde{F}^{a\mu\nu} \qquad \text{Peccei, Quinn 1977}$$

• Excellent dark matter candidate Weinberg, Wilczek 1978





Axion dark matter

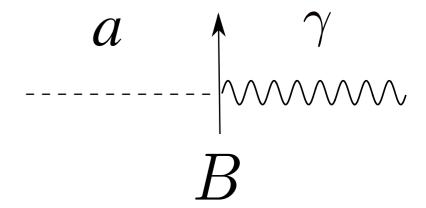
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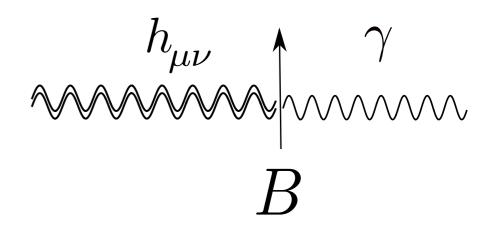
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Ideas and techniques developed for axions can be adapted to gravitational waves



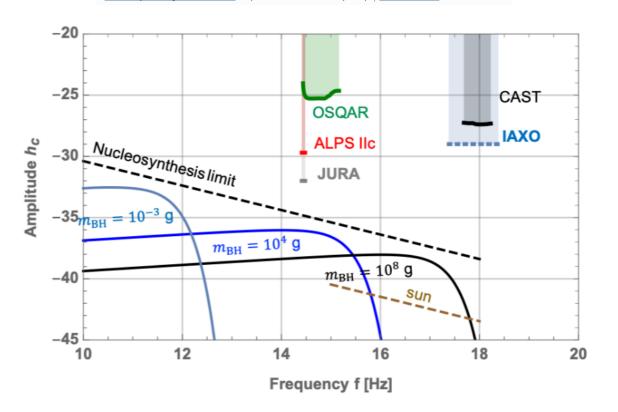
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Upper limits on the amplitude of ultra-high-frequency gravitational waves from graviton to photon conversion

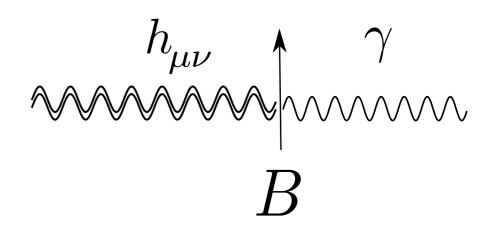
A. Ejlli [™], D. Ejlli, A. M. Cruise, G. Pisano & H. Grote



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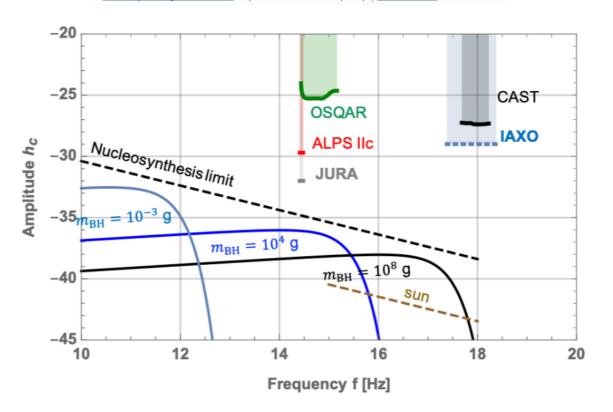
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- Distortions of the CMB

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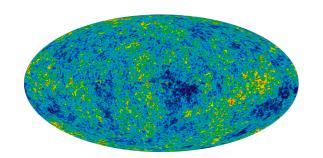


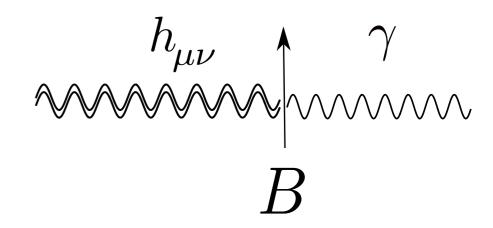
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Domcke, CGC 2021

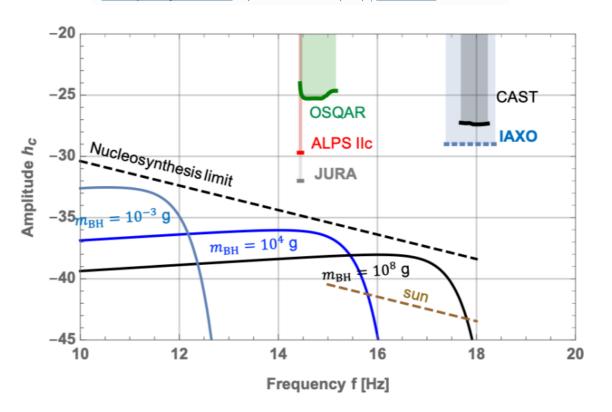
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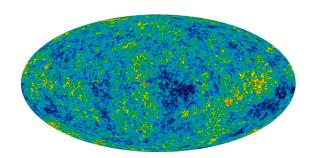


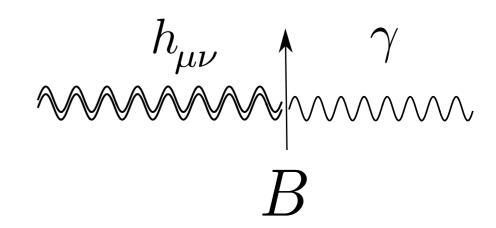
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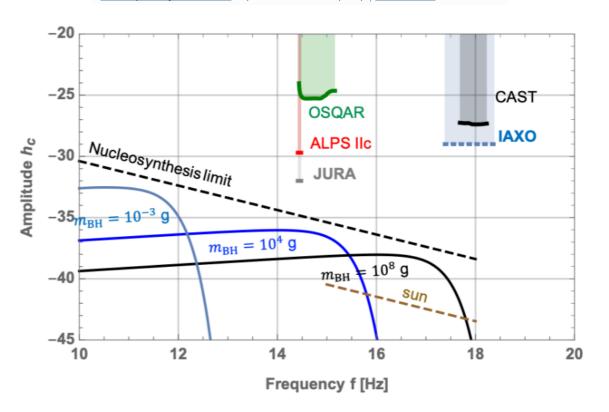
Dolgov, Ejlli 2012 Pshirkov, Baskaran 2009 Chen 1995





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Cosmic magnetic fields and multimessenger astronomy

Cosmic magnetic fields

Durrer, Neronov, 2013

$$\langle \mathbf{B}_{i}(\mathbf{x})\mathbf{B}_{j}(\mathbf{x}')\rangle = \frac{1}{(2\pi)^{3}a(t)^{4}} \int d^{3}k e^{i\mathbf{k}\cdot(\mathbf{x}'-\mathbf{x})} \left(\left(\delta_{ij} - \hat{k}_{i}\hat{k}_{j}\right) P_{B}(k) - i\epsilon_{ijk}\hat{k}_{k} P_{aB}(k) \right) ,$$

The adiabatic evolution of the magnetic field due to cosmic expansion is determined by the scale factor.

Cosmic magnetic fields

Durrer, Neronov, 2013

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The adiabatic evolution of the magnetic field due to cosmic expansion is determined by the scale factor.

$$\langle B^2 \rangle = \frac{1}{\pi^2 a(t)^4} \int_0^\infty dk k^2 P_B(k) = \int_{-\infty}^\infty d\log \lambda \, B_\lambda^2$$
where $B_\lambda^2 \equiv \frac{8\pi}{\lambda^3 a(t)^4} P_B\left(\frac{2\pi}{\lambda}\right)$,

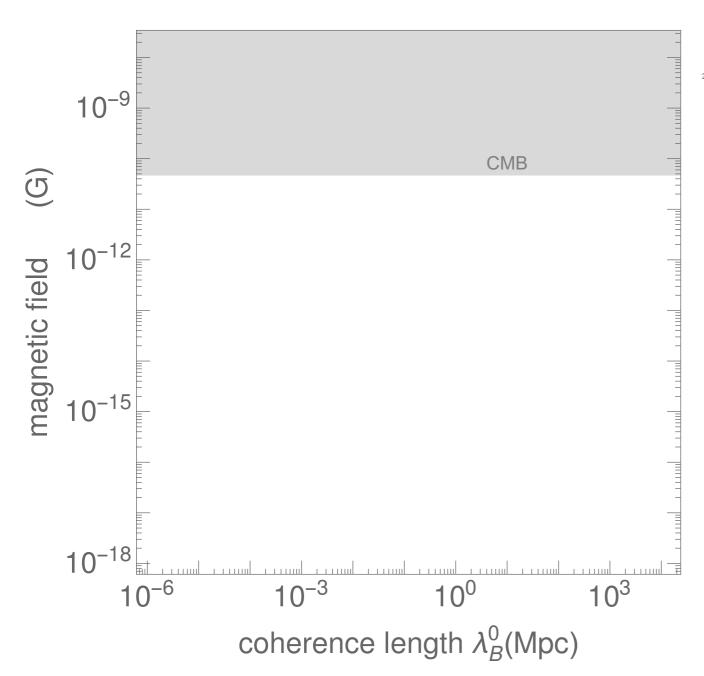
average magnetic field

$$\lambda_B = \int_0^\infty d\lambda \frac{B_\lambda^2}{\langle B^2 \rangle}$$

the coherence length

Cosmic magnetic fields in 2021

PHYSICAL REVIEW LETTERS 123, 021301 (2019)



Stringent Limit on Primordial Magnetic Fields from the Cosmic Microwave Background Radiation

Karsten Jedamzik^{1,*} and Andrey Saveliev^{2,3,†}

¹Laboratoire Univers et Particules de Montpellier, UMR5299-CNRS, Université de Montpellier, 34095 Montpellier, France ²Institute of Physics, Mathematics and Information Technology, Immanuel Kant Baltic Federal University, 236016 Kaliningrad, Russia ³Faculty of Computational Mathematics and Cybernetics, Lomonosov Moscow State University, 119991 Moscow, Russia

(Received 8 May 2018; revised manuscript received 13 September 2018; published 10 July 2019)

Primordial magnetic fields (PMFs), being present before the epoch of cosmic recombination, induce small-scale baryonic density fluctuations. These inhomogeneities lead to an inhomogeneous recombination process that alters the peaks and heights of the large-scale anisotropies of the cosmic microwave background (CMB) radiation. Utilizing numerical compressible MHD calculations and a Monte Carlo Markov chain analysis, which compares calculated CMB anisotropies with those observed by the *WMAP* and *Planck* satellites, we derive limits on the magnitude of putative PMFs. We find that the *total* remaining present day field, integrated over all scales, cannot exceed 47 pG for scale-invariant PMFs and 8.9 pG for PMFs with a violet Batchelor spectrum at 95% confidence level. These limits are more than one order of magnitude more stringent than any prior stated limits on PMFs from the CMB, which have not accounted for this effect.

Cosmic magnetic fields in 2021

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10^{-9} **CMB** magnetic field Blazars

Durrer, Neronov, 2013 coherence length $\lambda_B^0(\text{Mpc})$

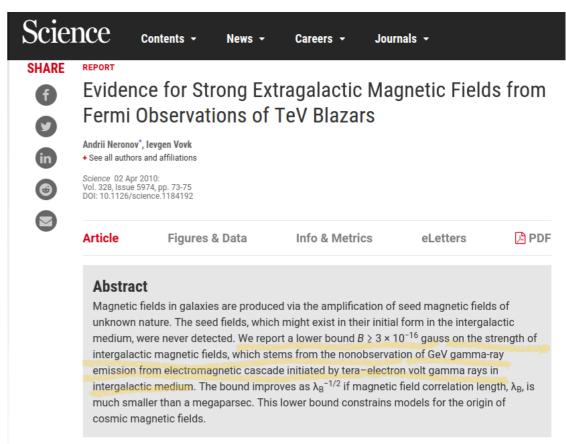
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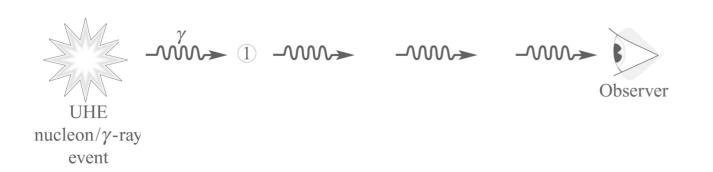
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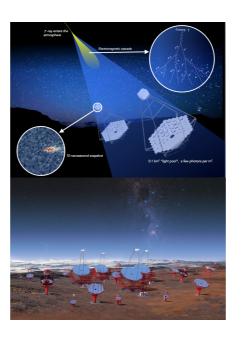
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Kronberg , 2016 Cambridge University Press

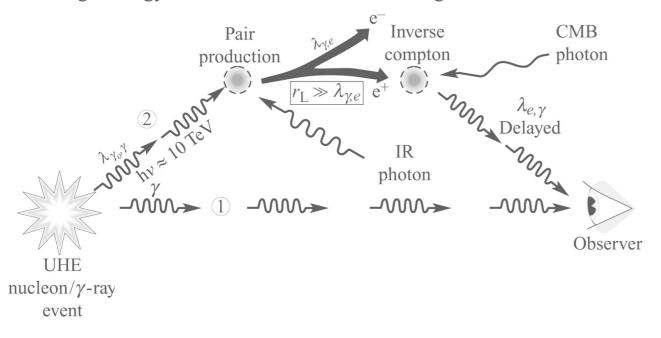


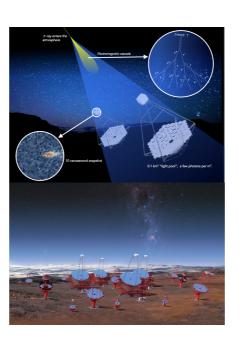


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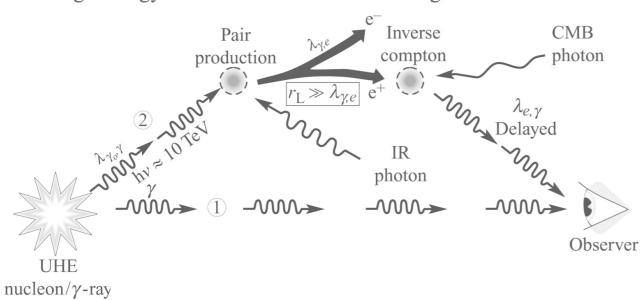
High energy $hv - e^+e^-$ cascades in the intergalactic medium



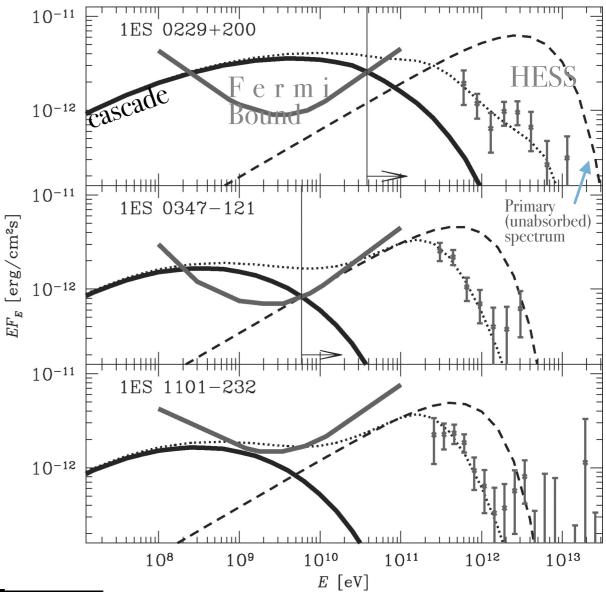


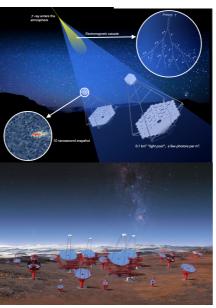
Kronberg , 2016 Cambridge University Press

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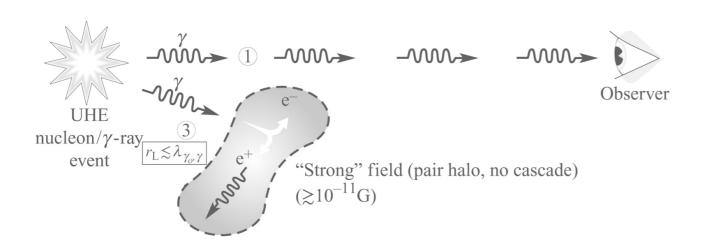


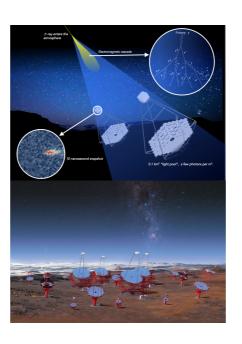
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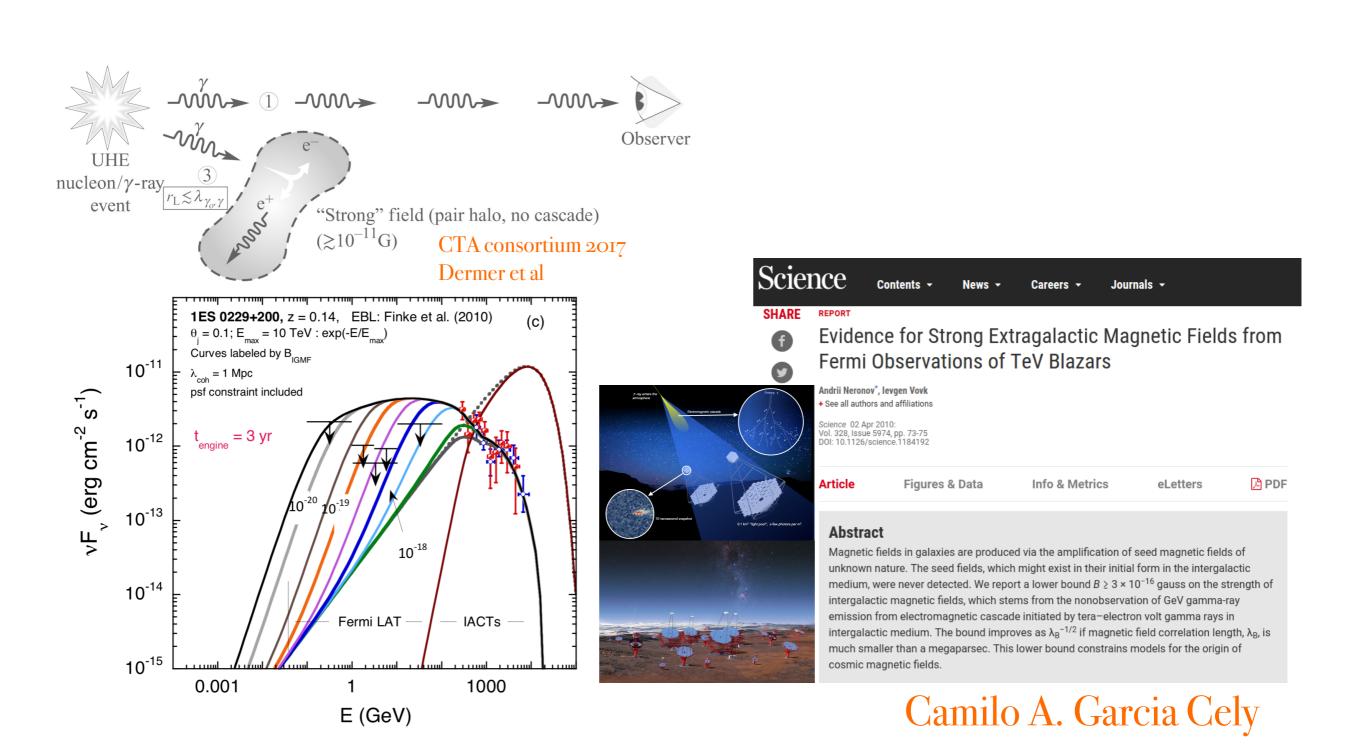


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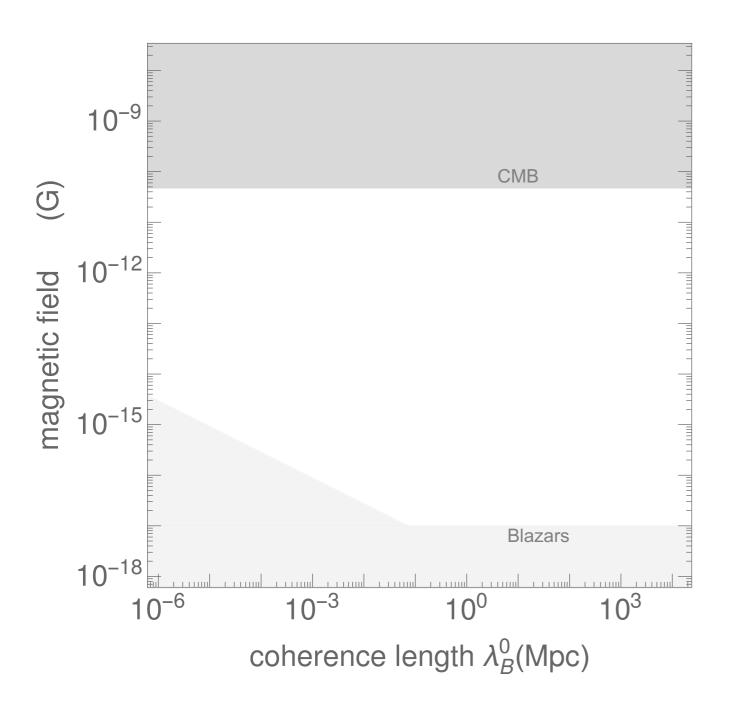


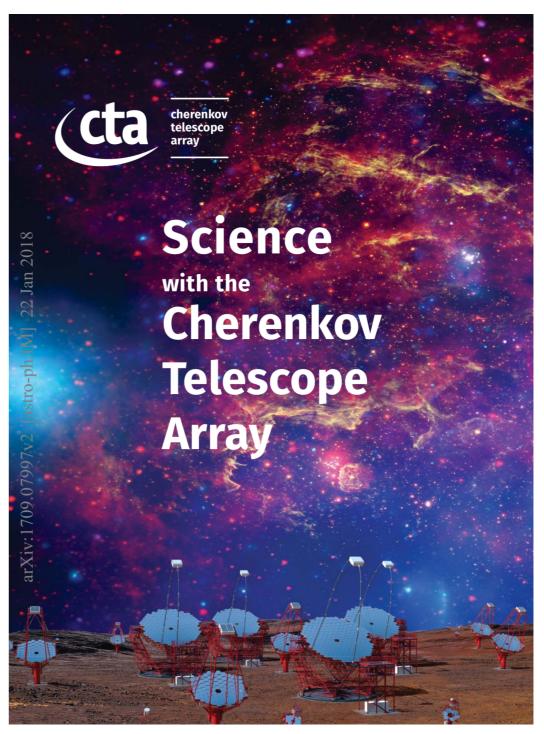
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leads to the appearance of an extended emission with an IGMF-dependent size. If the IGMF strength is in the range, $B \sim 10^{-16} - 10^{-12}$ G, the spatially-extended emission may be detectable and resolvable by CTA by virtue of its high sensitivity and angular resolution; e.g., for a source at a distance of 100 Mpc,

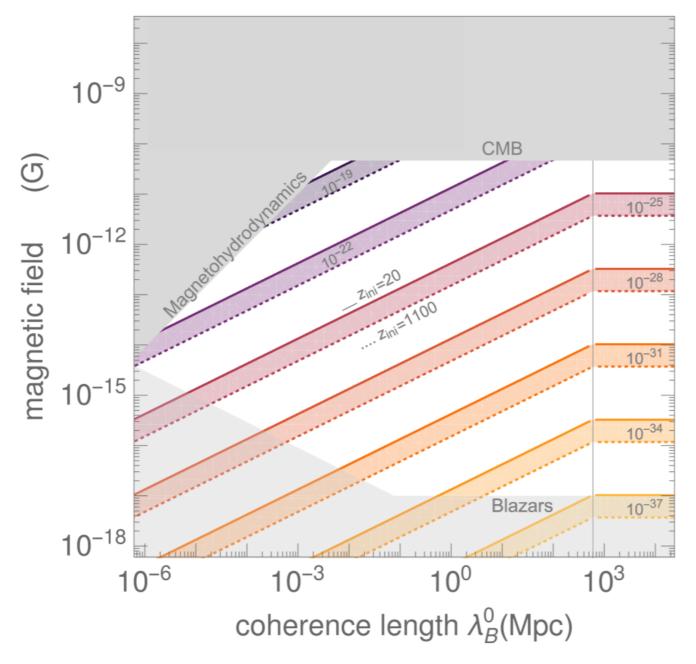
Domcke, CGC 2021





Cosmic magnetic fields in 2021

Domcke, CGC 2021



PHYSICAL REVIEW LETTERS **123**, 021301 (2019)

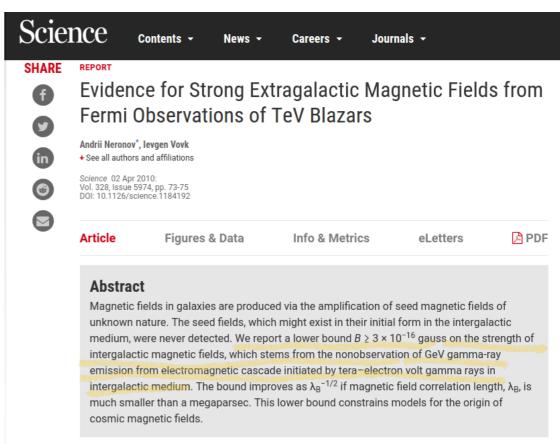
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$$\left(\Box + \omega_{\rm pl}^2\right) A_{\lambda} = -B \partial_{\ell} h_{\lambda}$$

$$\Box h_{\lambda} = 16\pi GB \, \partial_{\ell} A_{\lambda}$$

$$\omega_{\rm pl} = \sqrt{e^2 n_e/m_e}$$

The plasma frequency acts as an effective mass term

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Although cosmic magnetic fields are not expected to be perfectly homogeneous, coherent oscillations take place in highly homogeneous patches.

$$\mathcal{\ell}_{\rm OSC} = 4\omega/(1+z)^2 X_e(z) \omega_{\rm pl,0}^2 \ll 1\,{\rm pc}$$

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Cosmic magnetic fields in 2021

Domcke, CGC 2021

10^{-9} **CMB** 10-25 10^{-12} magnetic field 10-28 10-31 10^{-15} 10^{-34} Blazars 10^{-18} 10^{-6} 10^{-3} 10^{0} 10^{3} coherence length $\lambda_B^0(Mpc)$ $\langle \Gamma_{g \leftrightarrow \gamma} \rangle dt = \int_{0}^{z_{\text{ini}}} \frac{\langle \Gamma_{g \leftrightarrow \gamma} \rangle}{(1+z)H} dz$

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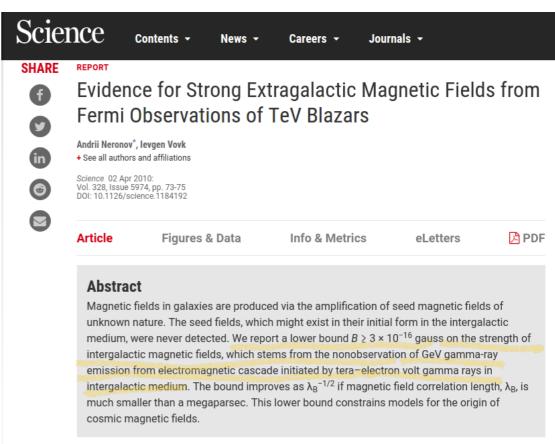
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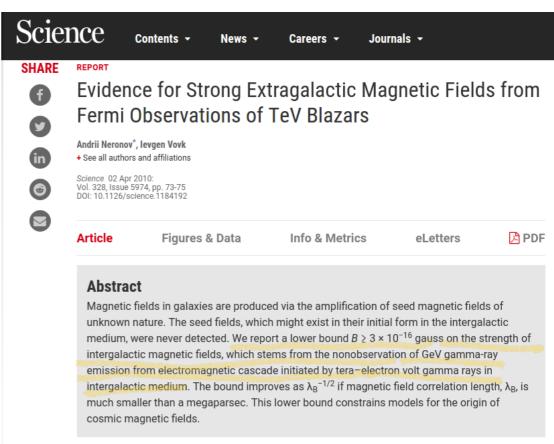
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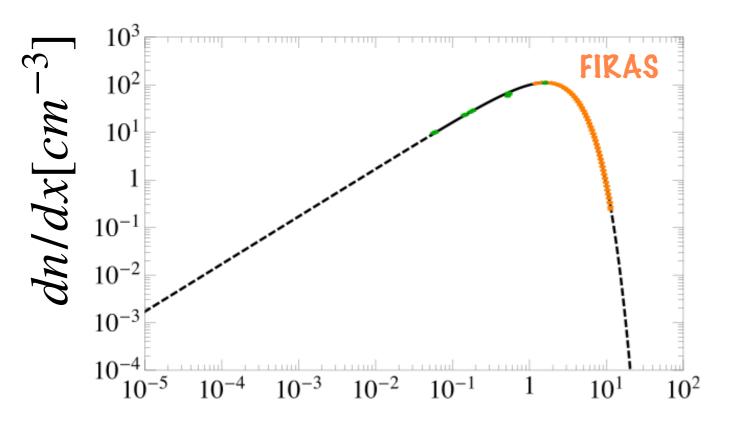
(Received 8 May 2018; revised manuscript received 13 September 2018; published 10 July 2019)

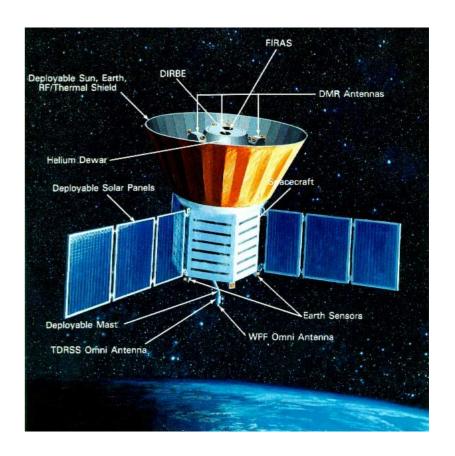
Primordial magnetic fields (PMFs), being present before the epoch of cosmic recombination, induce small-scale baryonic density fluctuations. These inhomogeneities lead to an inhomogeneous recombination process that alters the peaks and heights of the large-scale anisotropies of the cosmic microwave background (CMB) radiation. Utilizing numerical compressible MHD calculations and a Monte Carlo Markov chain analysis, which compares calculated CMB anisotropies with those observed by the WMAP and Planck satellites, we derive limits on the magnitude of putative PMFs. We find that the total remaining present day field, integrated over all scales, cannot exceed 47 pG for scale-invariant PMFs and 8.9 pG for PMFs with a violet Batchelor spectrum at 95% confidence level. These limits are more than one order of magnitude more stringent than any prior stated limits on PMFs from the CMB, which have not accounted for this effect.



CMB observations

CMB distortions





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THE COSMIC MICROWAVE BACKGROUND SPECTRUM FROM THE FULL $COBE^1$ FIRAS DATA SET

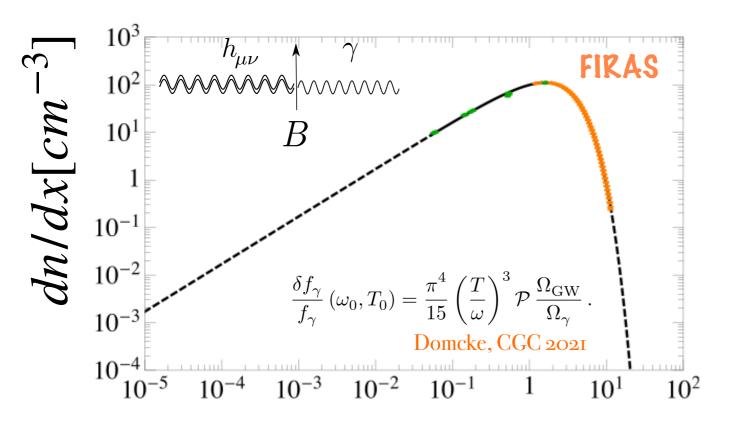
D. J. Fixsen,² E. S. Cheng,³ J. M. Gales,² J. C. Mather,³ R. A. Shafer,³ and E. L. Wright⁴
Received 1996 January 19; accepted 1996 July 11

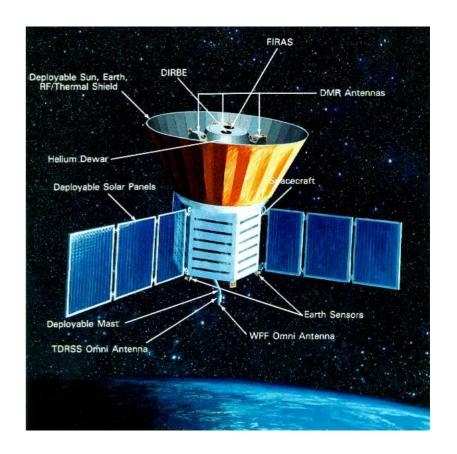
ABSTRACT

We have refined the analysis of the data from the FIRAS (Far-InfraRed Absolute Spectrophotometer) on board the COBE (COsmic Background Explorer). The FIRAS measures the difference between the cosmic microwave background and a precise blackbody spectrum. We find new, tighter upper limits on general deviations from a blackbody spectrum. The rms deviations are less than 50 parts per million of the peak of the cosmic microwave background radiation. For the Comptonization and chemical potential, we find $|y| < 15 \times 10^{-6}$ and $|\mu| < 9 \times 10^{-5}$ (95% confidence level [CL]). There are also refinements in the absolute temperature, 2.728 ± 0.004 K (95% CL), the dipole direction, (ℓ , b) = (264°14 \pm 0.30, 48°26 \pm 0.30) (95% CL), and the amplitude, 3.372 ± 0.014 mK (95% CL). All of these results agree with our previous publications.

Subject headings: cosmic microwave background — cosmology: observations

CMB distortions





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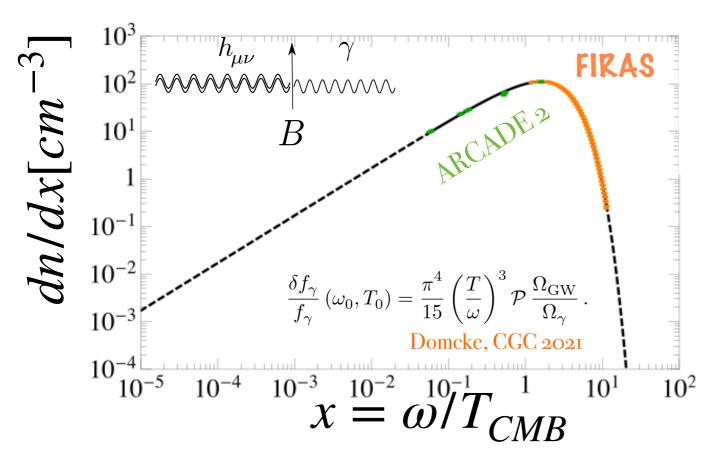
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Rayleigh-Jeans Tail

THE ASTROPHYSICAL JOURNAL



• Largely unexplored with upcoming advances in radio astronomy probing it in the near future.

ARCADE 2 MEASUREMENT OF THE ABSOLUTE SKY BRIGHTNESS AT 3-90 GHz

D. J. Fixsen¹, A. Kogut², S. Levin³, M. Limon⁴, P. Lubin⁵, P. Mirel⁶, M. Seiffert³, J. Singal⁷, E. Wollack², T. Villela⁸ + Show full author list
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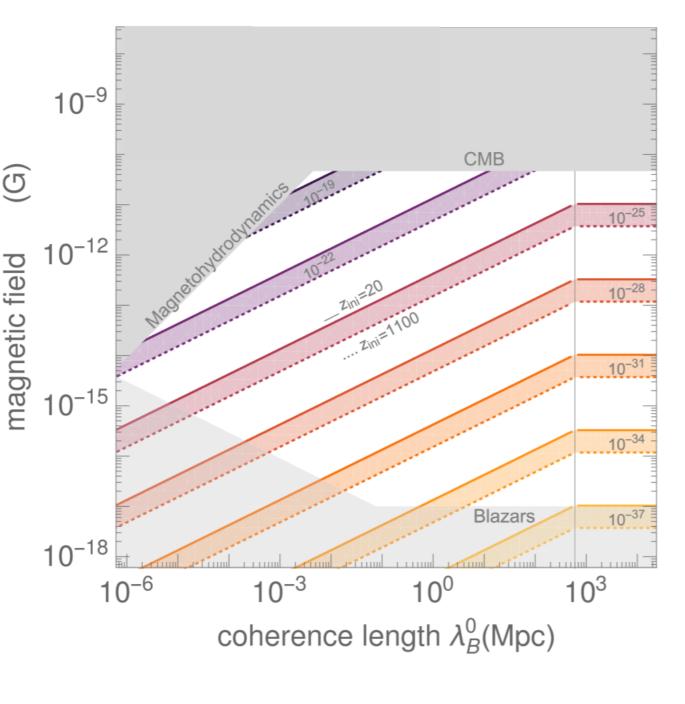
The Astrophysical Journal, Volume 734, Number 1

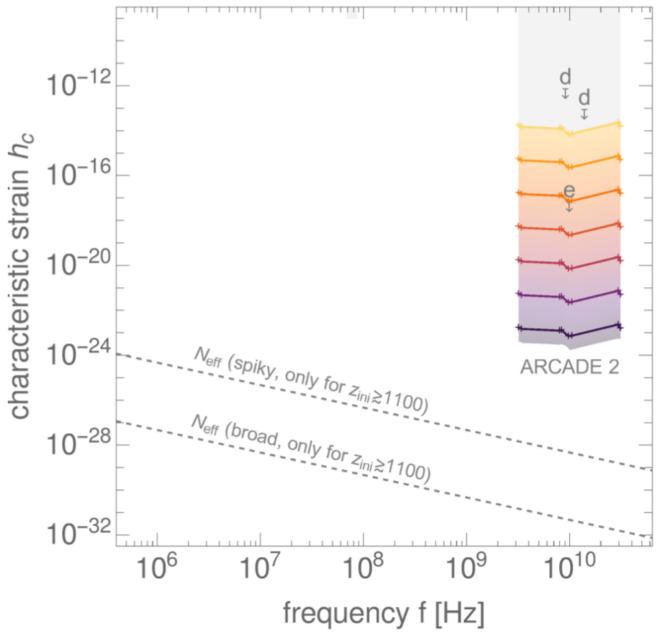
Upper bounds on stochastic gravitational waves

PHYSICAL REVIEW LETTERS 126, 021104 (2021)

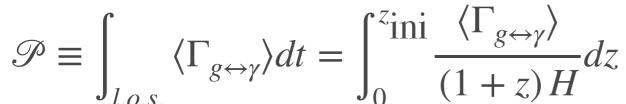
Potential of Radio Telescopes as High-Frequency Gravitational Wave Detectors

Valerie Domcke^{1,2,3,*} and Camilo Garcia-Cely^{1,†}





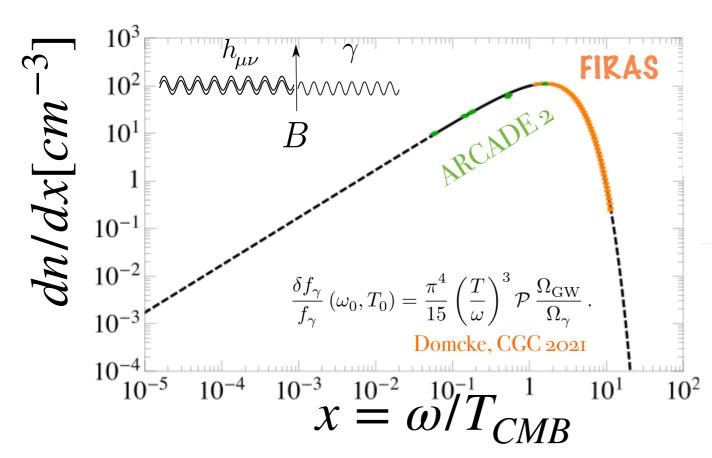
existing laboratory bounds from



- d) magnon detector Ito, Soda '04
- e) magnetic conversion detector Cruise et al '12

Rayleigh-Jeans Tail

THE ASTROPHYSICAL JOURNAL



- Largely unexplored with upcoming advances in radio astronomy probing it in the near future.
- Puzzling signal by EDGES.
 (Experiment to Detect the Global Epoch of Reionization Signature)

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nature

An absorption profile centred at 78 megahertz in the sky-averaged spectrum

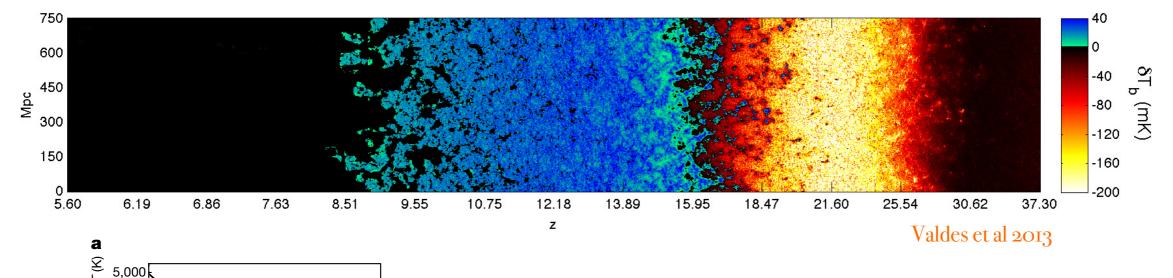
Judd D. Bowman ☑, Alan E. E. Rogers, Raul A. Monsalve, Thomas J. Mozdzen & Nivedita Mahesh

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Camilo A. Garcia Cely

Expectations for a 21 cm signal



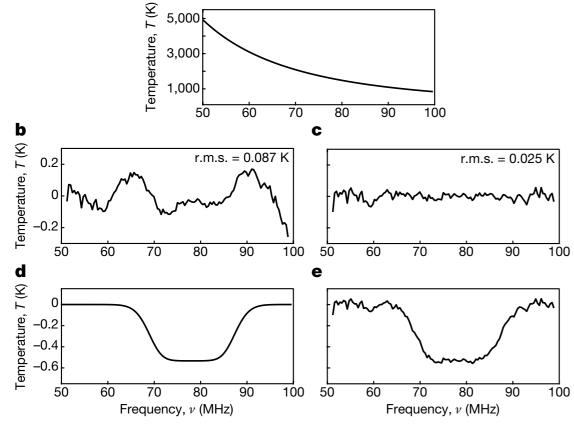


Figure 1 | **Summary of detection. a**, Measured spectrum for the reference dataset after filtering for data quality and radio-frequency interference. The spectrum is dominated by Galactic synchrotron emission. **b**, **c**, Residuals after fitting and removing only the foreground model (**b**) or the foreground and 21-cm models (**c**). **d**, Recovered model profile of the 21-cm absorption, with a signal-to-noise ratio of 37, amplitude of 0.53 K, centre frequency of 78.1 MHz and width of 18.7 MHz. **e**, Sum of the 21-cm model (**d**) and its residuals (**c**).

nature

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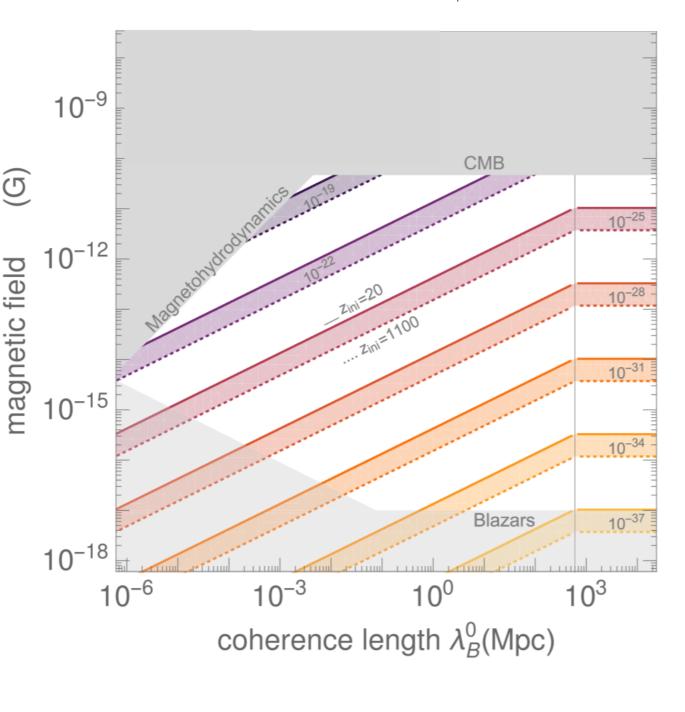
Camilo A. Garcia Cely

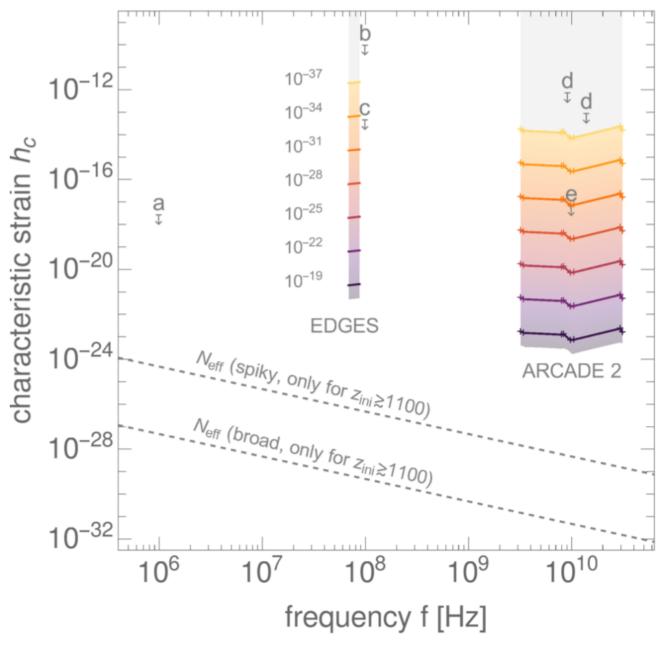
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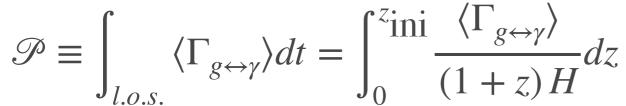
Valerie Domcke^{1,2,3,*} and Camilo Garcia-Cely^{1,†}





existing laboratory bounds from

- a) superconducting parametric converter Reece et al '84
- b) waveguide Cruise Ingley '06
- c) 0.75 m interferometer Akutsu '08
- d) magnon detector Ito, Soda '04
- e) magnetic conversion detector Cruise et al '12



Conclusions

• The Gertsenshtein effect during the dark ages provides a powerful way to probe gravitational waves in the MHz-GHz range from distortions of the Rayleigh-Jeans CMB tail.

• With upcoming advances in 21cm astronomy targeting precisely this frequency range with increasing accuracy, it becomes conceivable to push the limits derived from radio telescopes below the cosmological bound constraining the total energy in gravitational waves.