

# The CMB as a detector of high-frequency gravitational waves

**Camilo A. Garcia Cely**

*Alexander von Humboldt fellow*



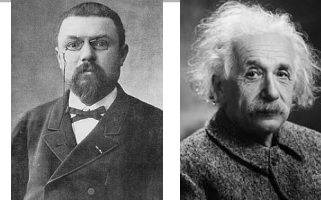
Quantum Universe Day

DESY

June 23, 2020

In collaboration with Valerie Domcke,  
based on arXiv:2006.01161 [astro-ph.CO]

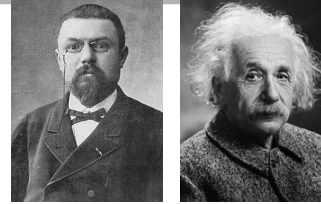
# Gravitational Waves



- Predicted by Poincaré (1905).
- Einstein provided a firm theoretical ground for them (1916).

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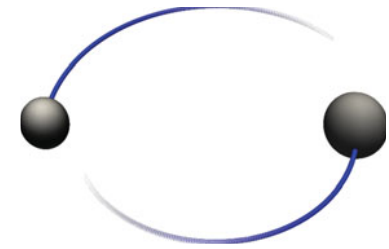
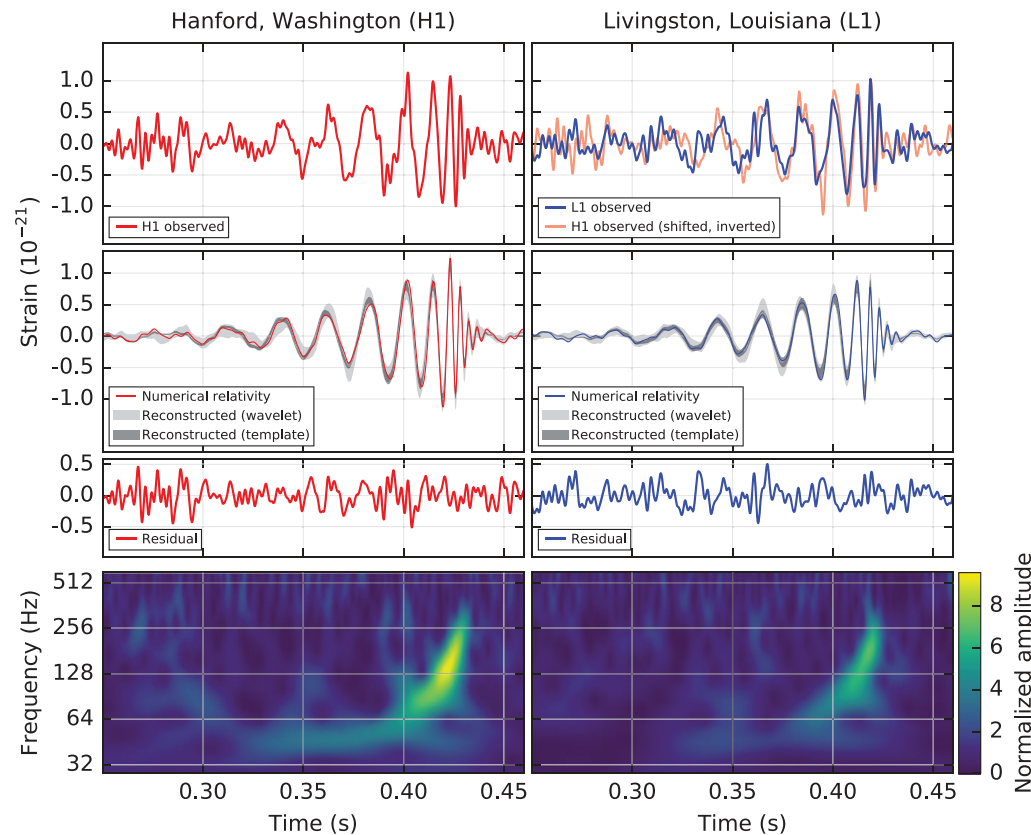
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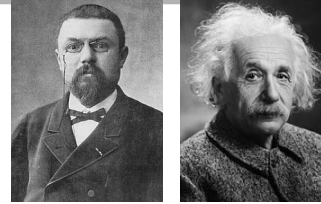
PRL **116**, 061102 (2016) PHYSICAL REVIEW LETTERS week ending  
12 FEBRUARY 2016



Terrestrial  
interferometers



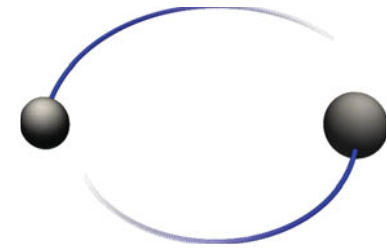
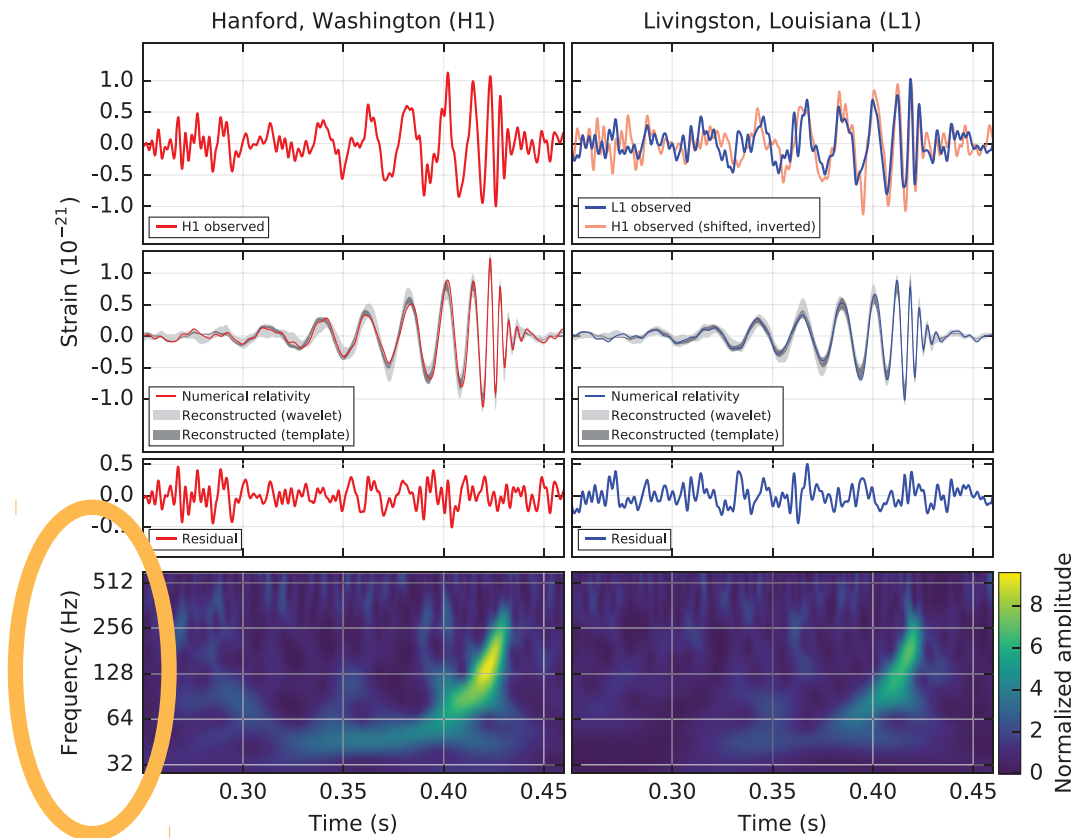
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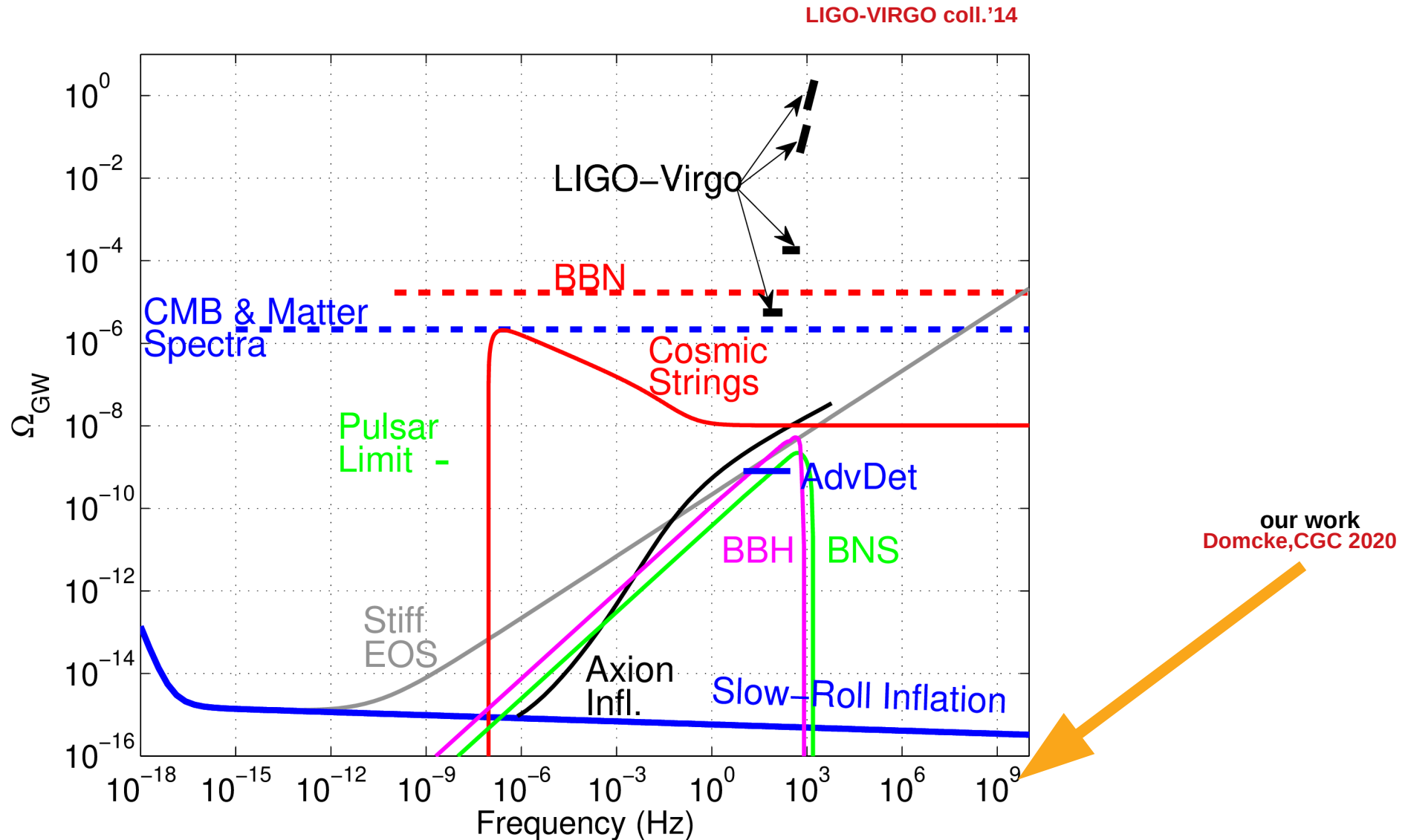
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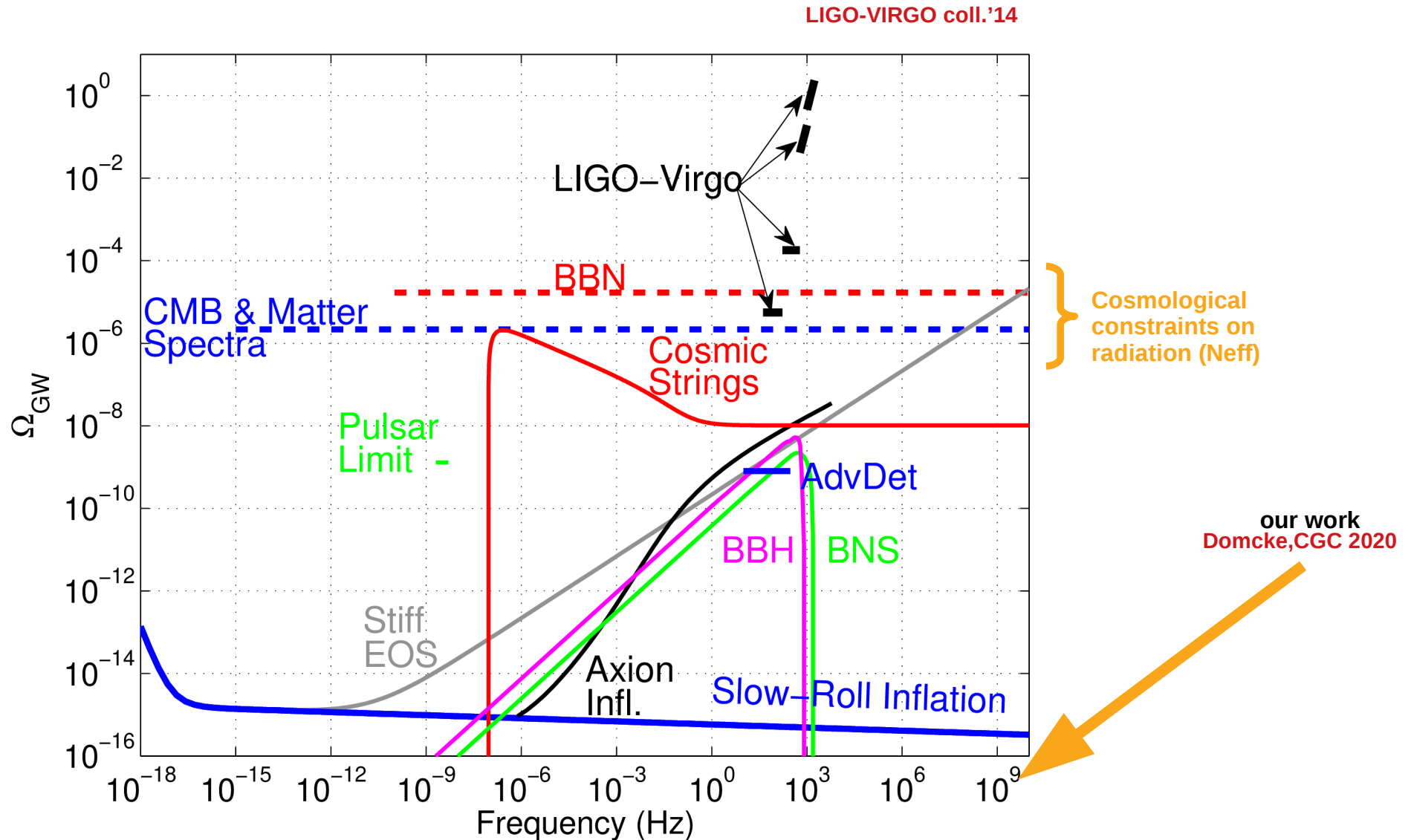
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# Gravitational waves spectrum



# Gravitational waves spectrum



# Revisiting Gertsenhstein's ideas

SOVIET PHYSICS JETP

VOLUME 16, NUMBER 2

FEBRUARY, 1963

## ON THE DETECTION OF LOW FREQUENCY GRAVITATIONAL WAVES

M. E. GERTSENHSTEIN and V. I. PUSTOVOIT

Submitted to JETP editor March 3, 1962

J. Exptl. Theoret. Phys. (U.S.S.R.) 43, 605-607 (August, 1962)

It is shown that the sensitivity of the electromechanical experiments for detecting gravitational waves by means of piezocrystals is ten orders of magnitude worse than that estimated by Weber.<sup>[1]</sup> In the low frequency range it should be possible to detect gravitational waves by the shift of the bands in an optical interferometer. The sensitivity of this method is investigated.

Terrestrial  
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SOVIET PHYSICS JETP

VOLUME 14, NUMBER 1

JANUARY, 1962

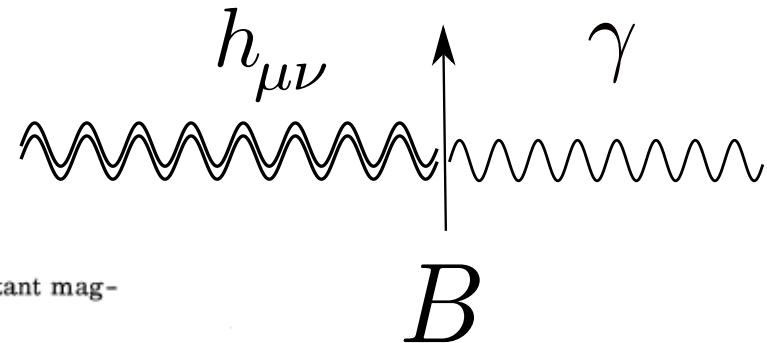
## WAVE RESONANCE OF LIGHT AND GRAVITATIONAL WAVES

M. E. GERTSENSHTEĬN

Submitted to JETP editor July 29, 1960

J. Exptl. Theoret. Phys. (U.S.S.R.) **41**, 113-114 (July, 1961)

The energy of gravitational waves excited during the propagation of light in a constant magnetic or electric field is estimated.



SOVIET PHYSICS JETP

VOLUME 16, NUMBER 2

FEBRUARY, 1963

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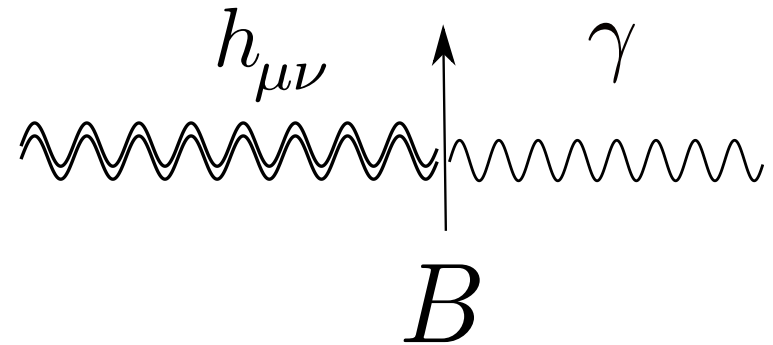
Terrestrial  
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# The Gertsenhstein Effect

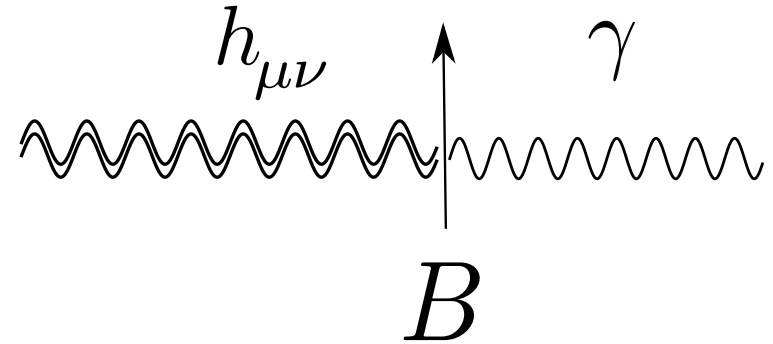
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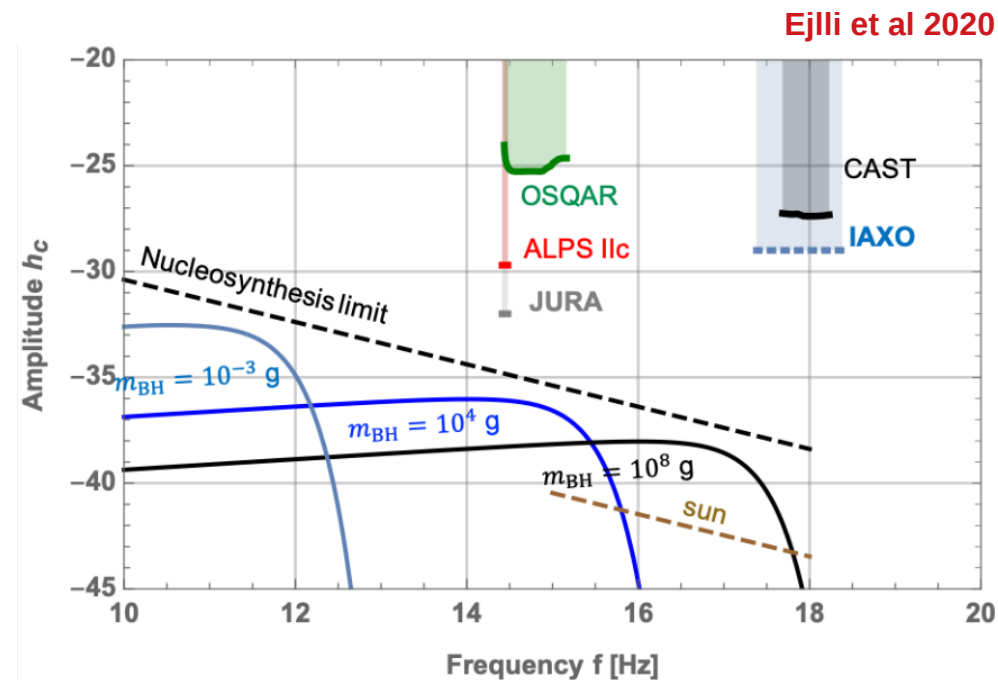
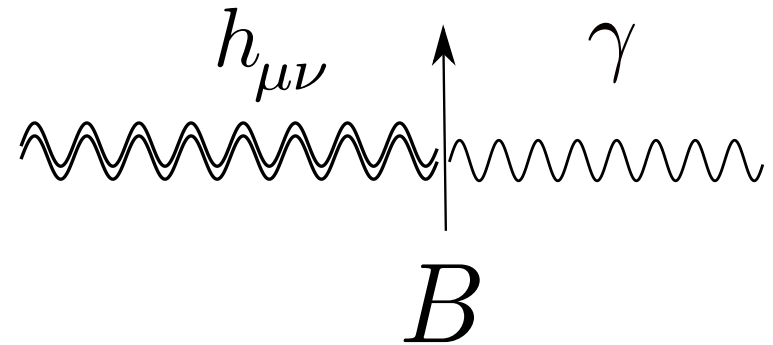
Raffelt, Stodolsk '89



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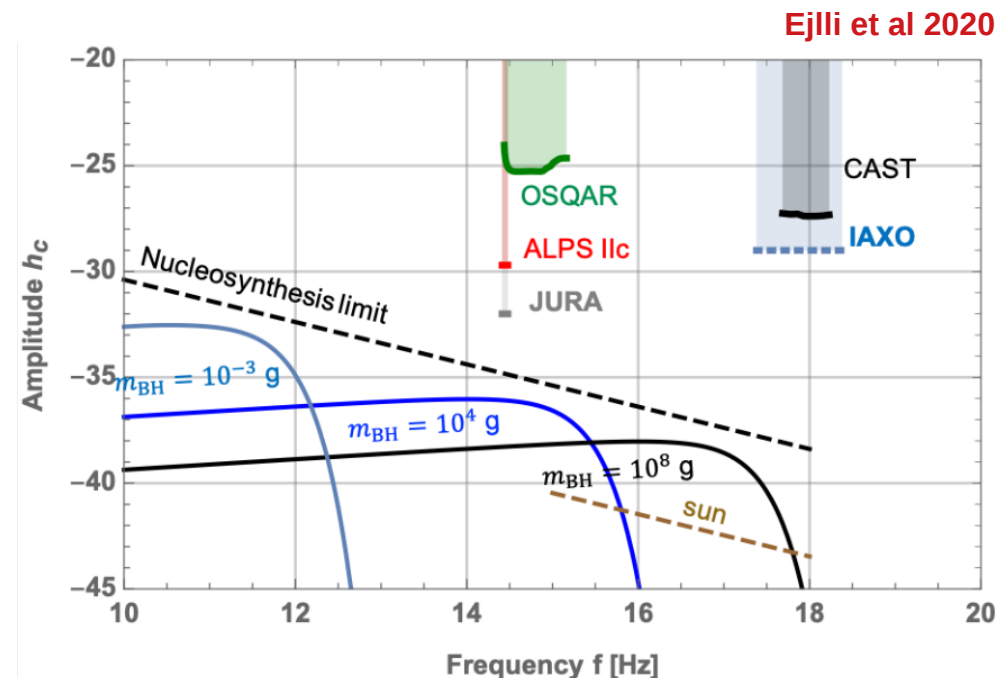
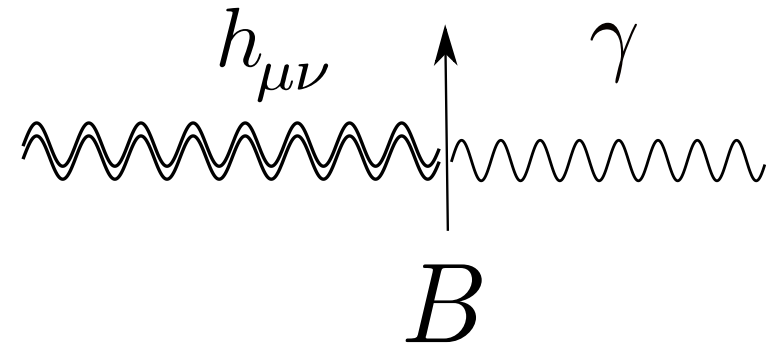
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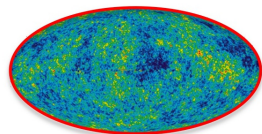
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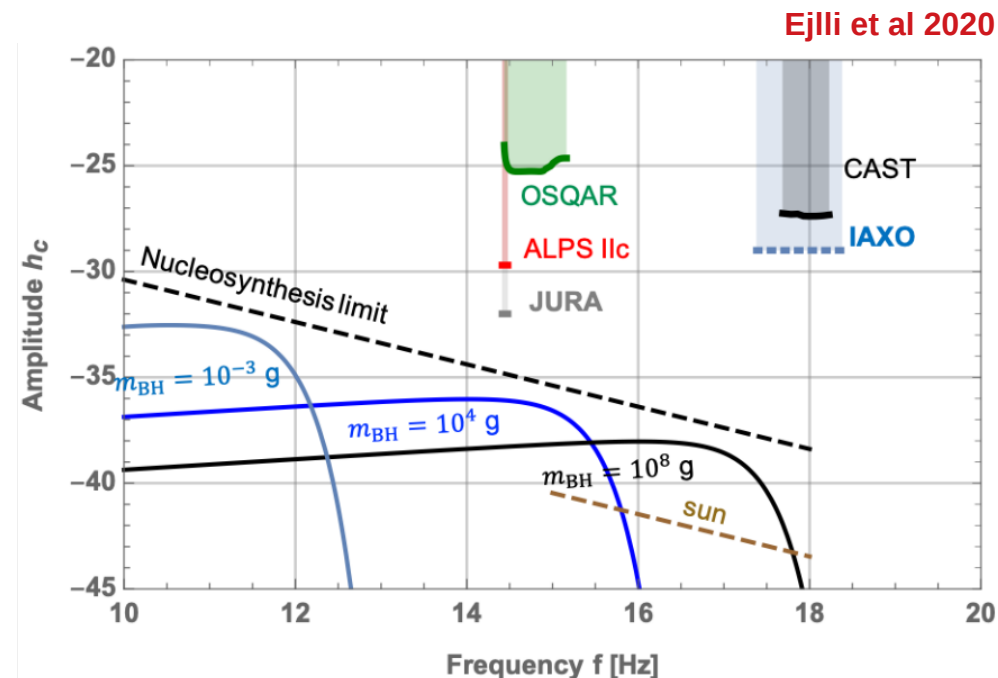
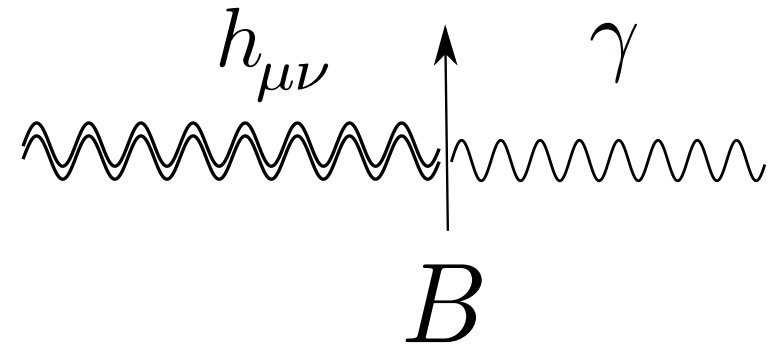
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Distortions of the CMB?

Domcke,CGC 2020



The CMB as a detector of high-frequency GWs



Camilo A. Garcia Cely (DESY)

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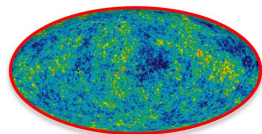
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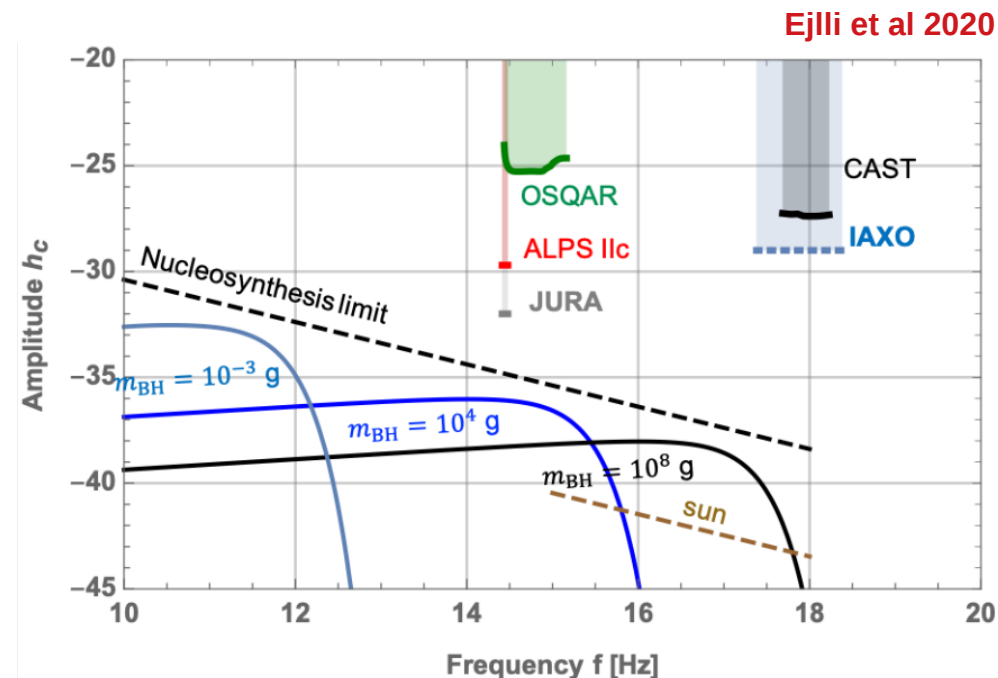
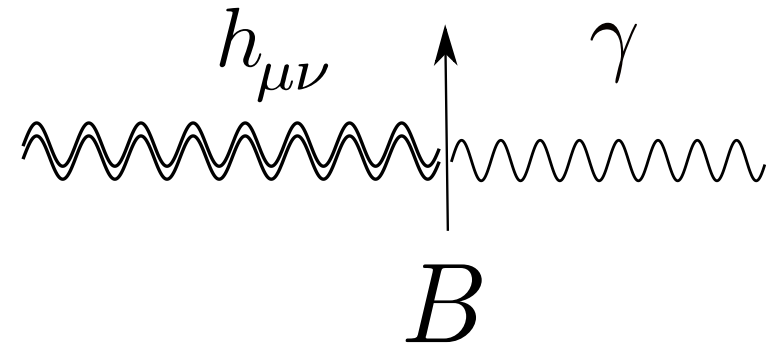
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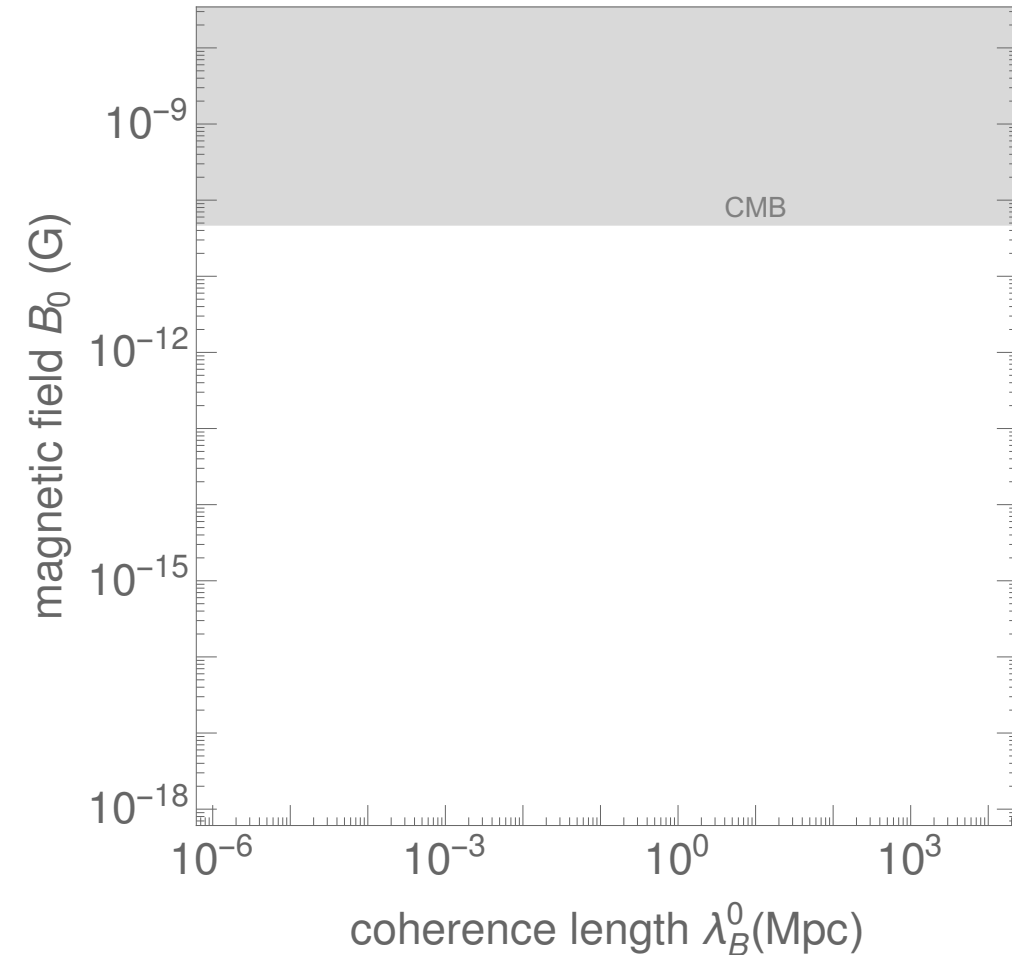


Domcke, CGC 2020  
Dolgov, Ejlli 2012  
Pshirkov, Baskaran 2009  
Chen 1995  
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# Cosmic Magnetic Fields in 2020

Domcke, CGC 2020



PHYSICAL REVIEW LETTERS **123**, 021301 (2019)

## Stringent Limit on Primordial Magnetic Fields from the Cosmic Microwave Background Radiation

Karsten Jedamzik<sup>1,\*</sup> and Andrey Saveliev<sup>2,3,†</sup>

<sup>1</sup>Laboratoire Univers et Particules de Montpellier, UMR5299-CNRS, Université de Montpellier, 34095 Montpellier, France

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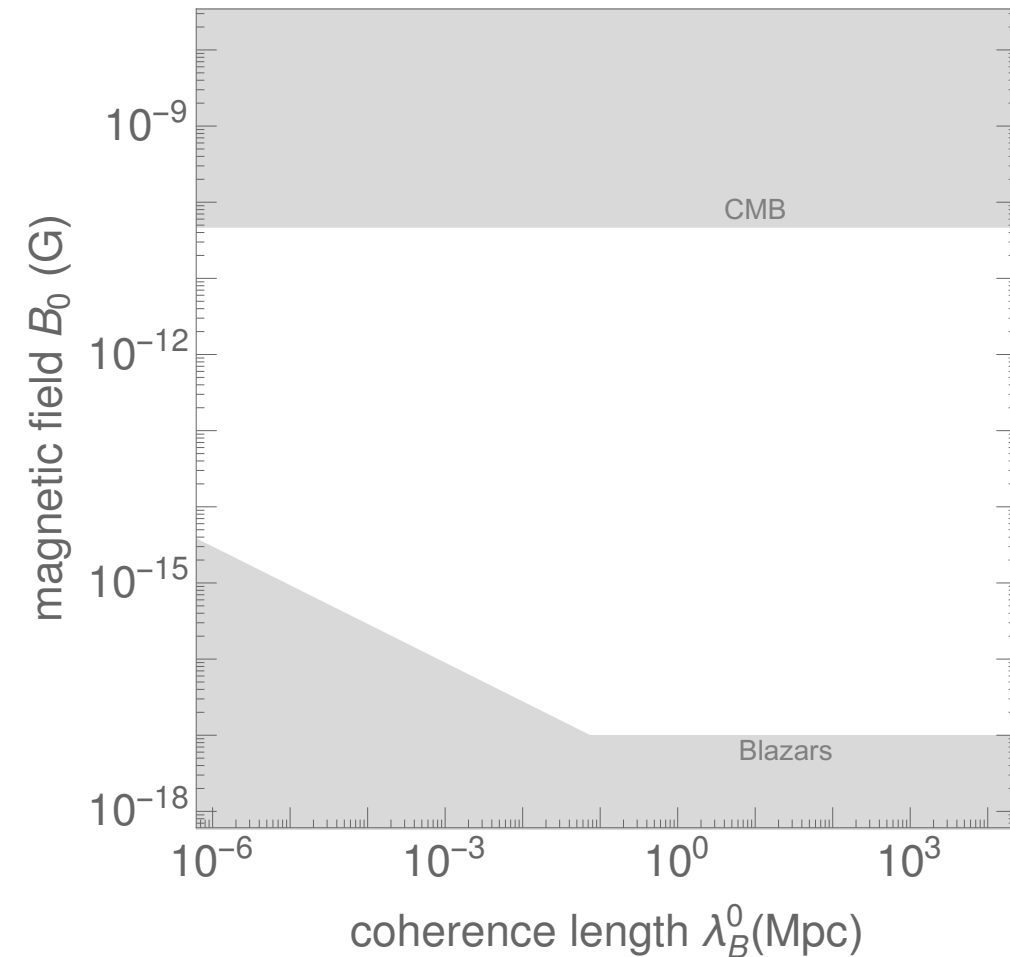


(Received 8 May 2018; revised manuscript received 13 September 2018; published 10 July 2019)

Primordial magnetic fields (PMFs), being present before the epoch of cosmic recombination, induce small-scale baryonic density fluctuations. These inhomogeneities lead to an inhomogeneous recombination process that alters the peaks and heights of the large-scale anisotropies of the cosmic microwave background (CMB) radiation. Utilizing numerical compressible MHD calculations and a Monte Carlo Markov chain analysis, which compares calculated CMB anisotropies with those observed by the *WMAP* and *Planck* satellites, we derive limits on the magnitude of putative PMFs. We find that the *total* remaining present day field, integrated over all scales, cannot exceed 47 pG for scale-invariant PMFs and 8.9 pG for PMFs with a violet Batchelor spectrum at 95% confidence level. These limits are more than one order of magnitude more stringent than any prior stated limits on PMFs from the CMB, which have not accounted for this effect.

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## Evidence for Strong Extragalactic Magnetic Fields from Fermi Observations of TeV Blazars

Andrii Neronov\*, Ievgen Vovk

+ See all authors and affiliations

Science 02 Apr 2010;  
Vol. 328, Issue 5974, pp. 73-75  
DOI: 10.1126/science.1184192

Article

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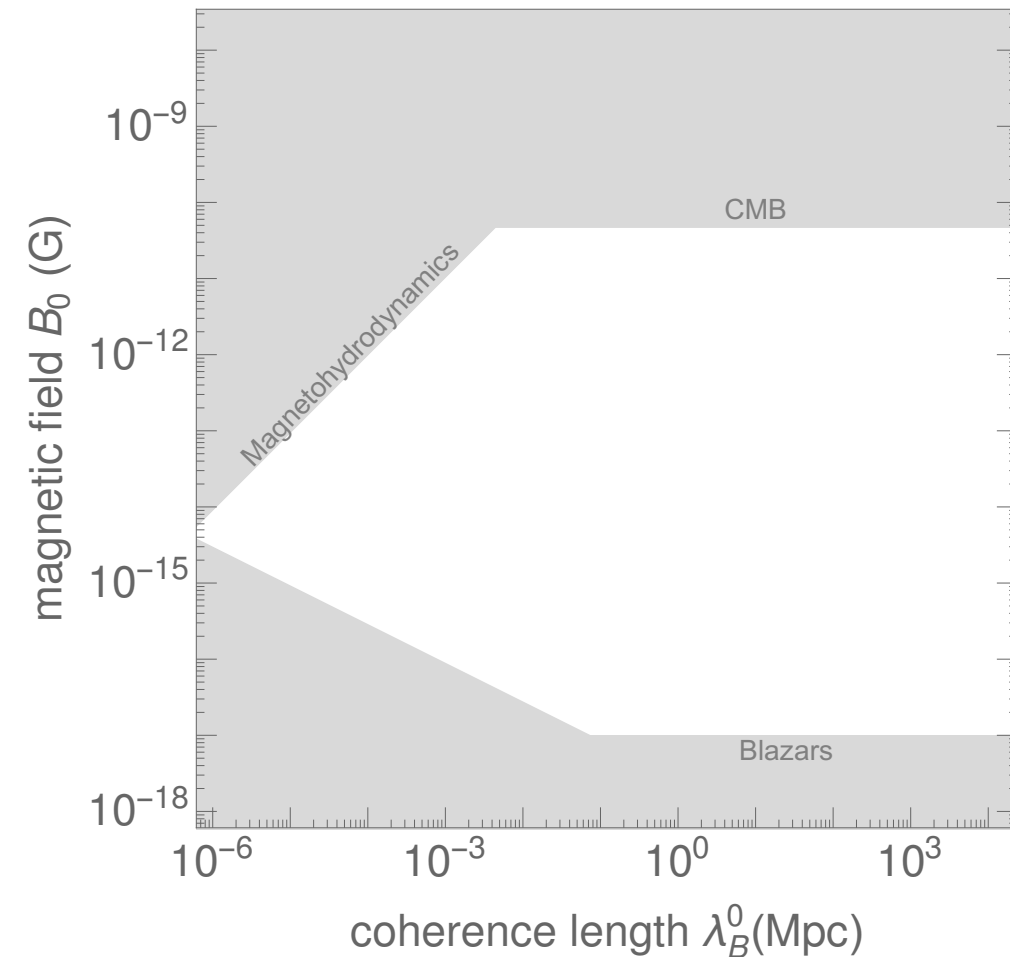
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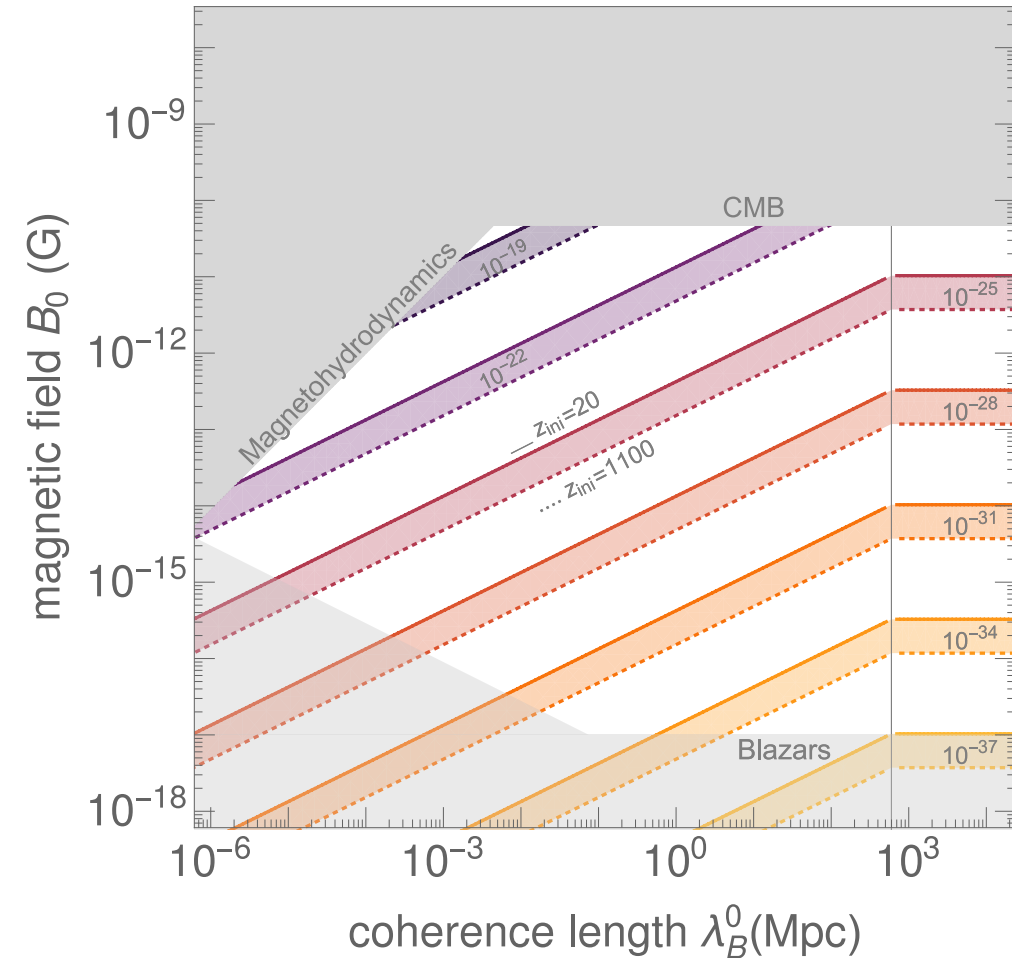
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$$\mathcal{P} \equiv \int_{l.o.s.} \langle \Gamma_{g \leftrightarrow \gamma} \rangle dt$$

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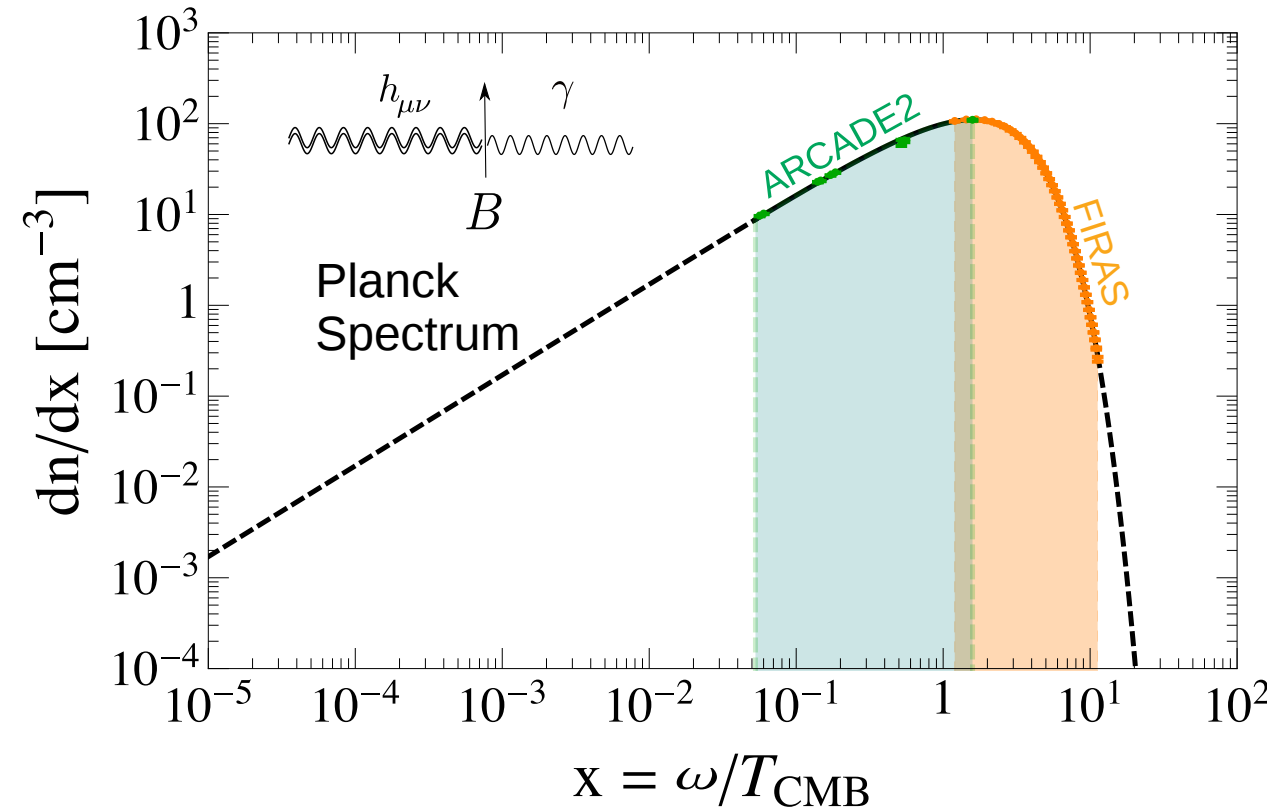
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Camilo A. Garcia Cely (DESY)

# CMB distortions in 2020



Absolute Radiometer for Cosmology,  
Astrophysics, and Diffuse Emission

THE ASTROPHYSICAL JOURNAL

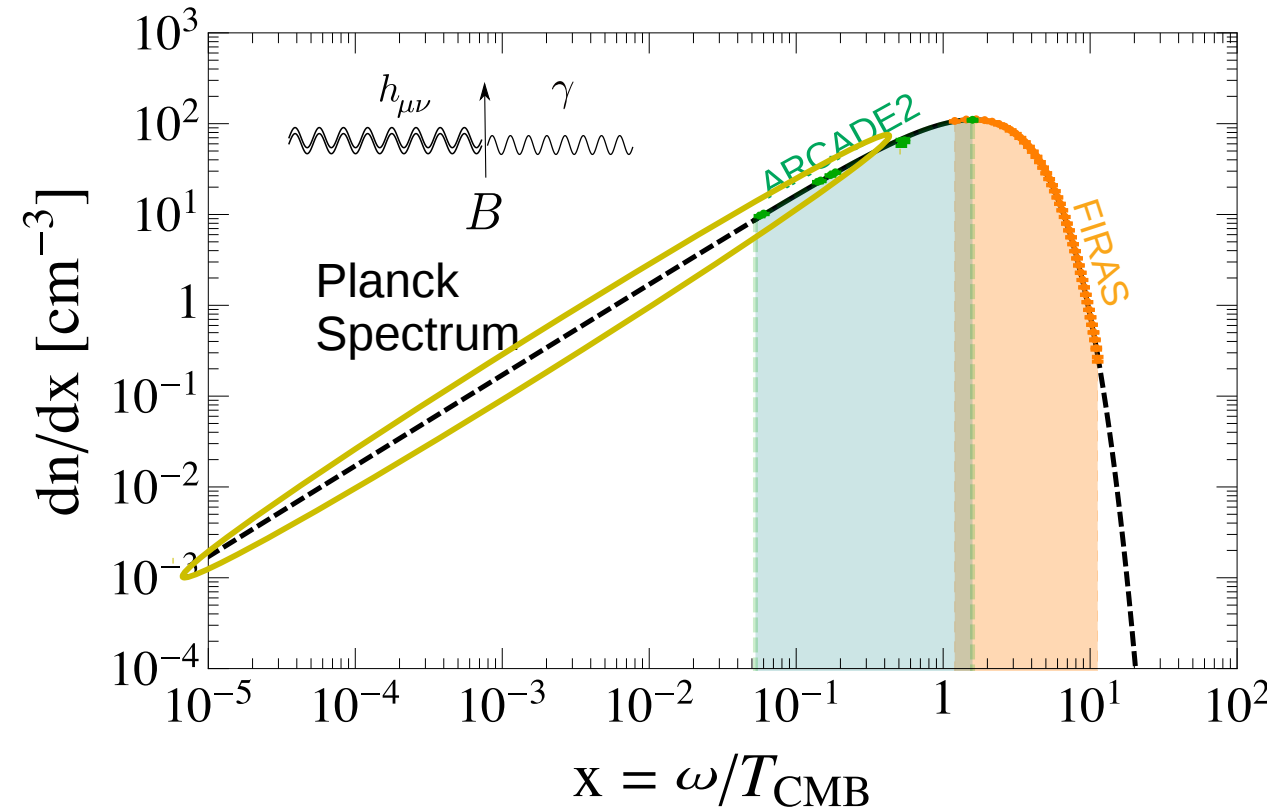
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T. Villella<sup>8</sup> [+Show full author list](#)

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[The Astrophysical Journal, Volume 734, Number 1](#)

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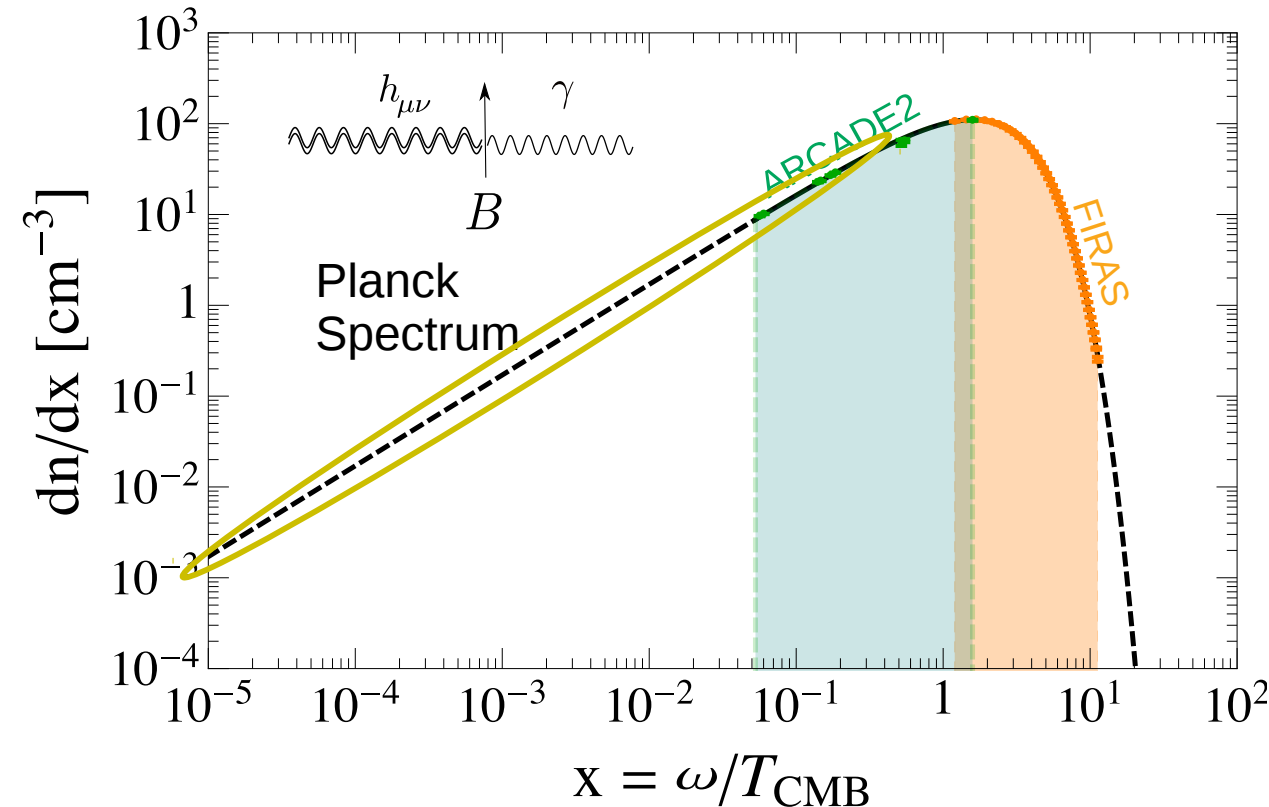
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## Rayleigh-Jeans tail

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nature

Published: 01 March 2018

### An absorption profile centred at 78 megahertz in the sky-averaged spectrum

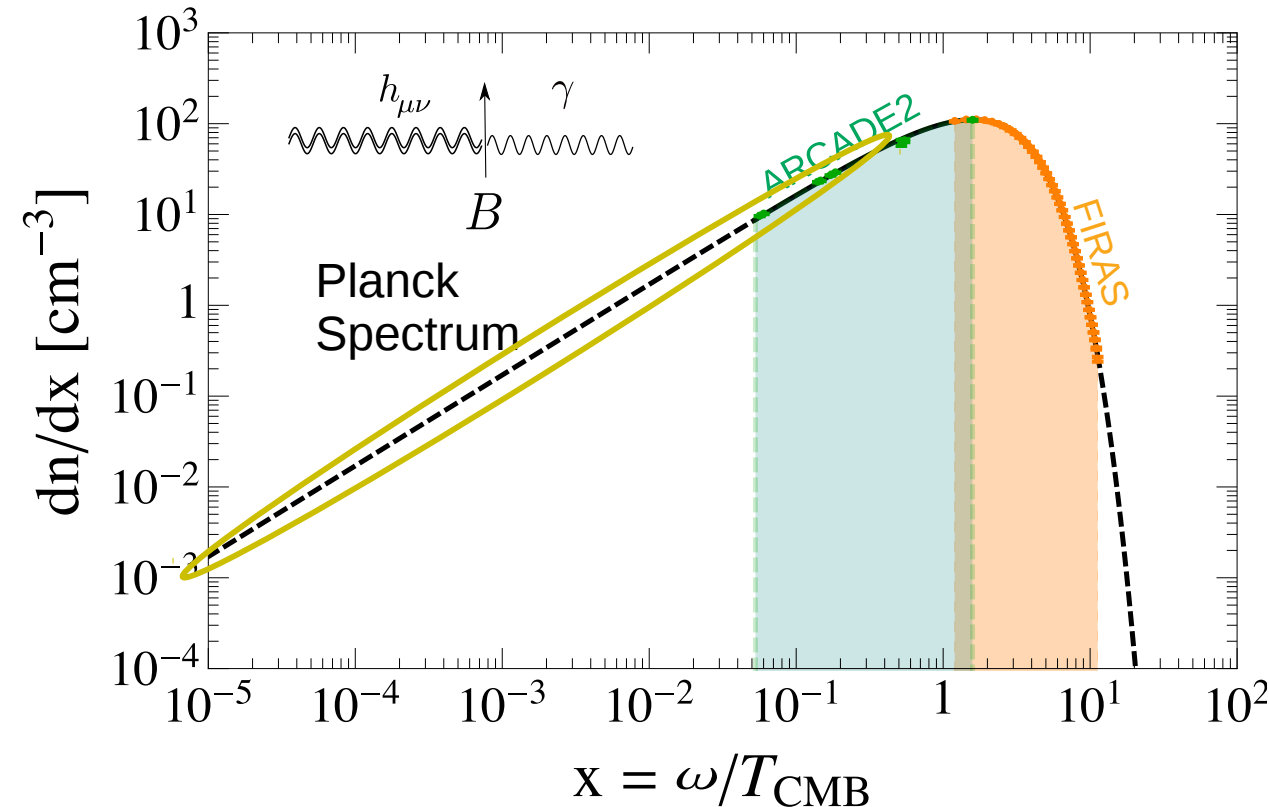
Judd D. Bowman [✉](#), Alan E. E. Rogers, Raul A. Monsalve, Thomas J. Mozdzen &  
Nivedita Mahesh

Experiment to Detect the Global  
Epoch of Reionization Signature



Camilo A. Garcia Cely (DESY)

# CMB distortions in 2020



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- They may conceivably push these bounds below the Neff constraint.

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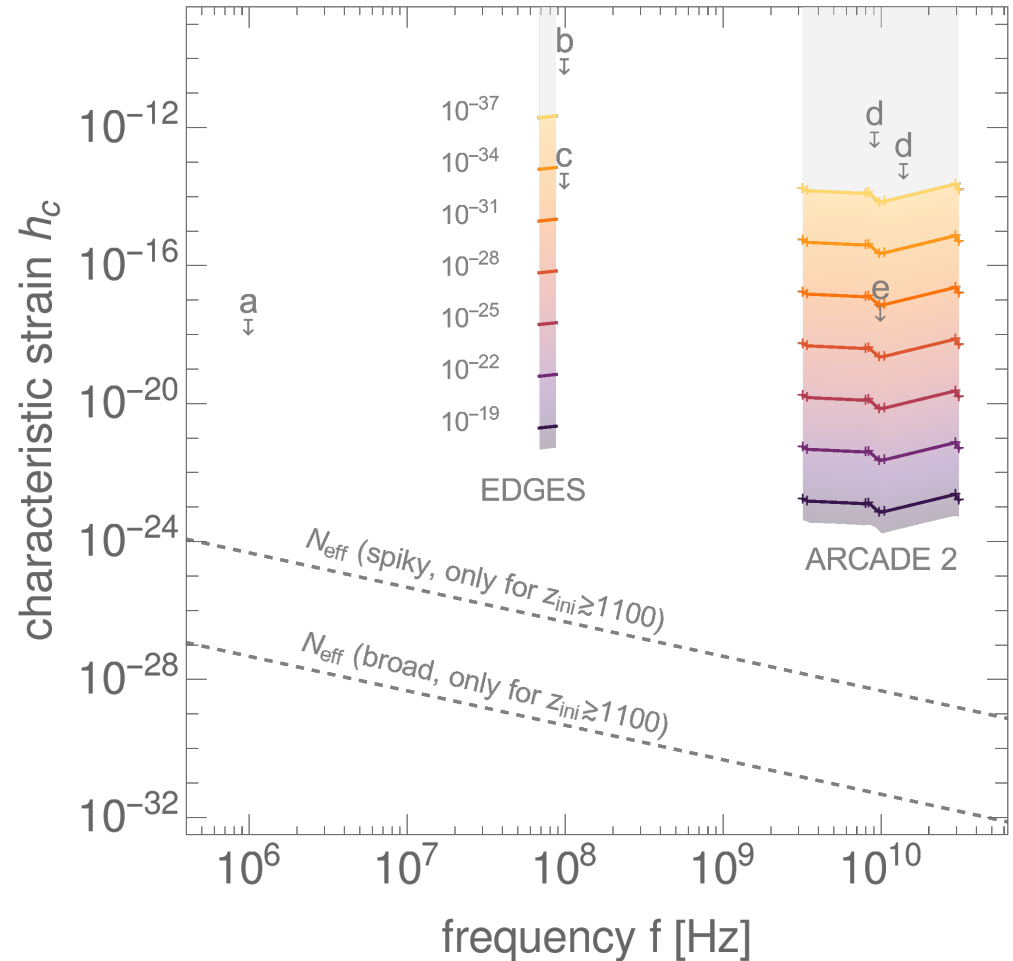
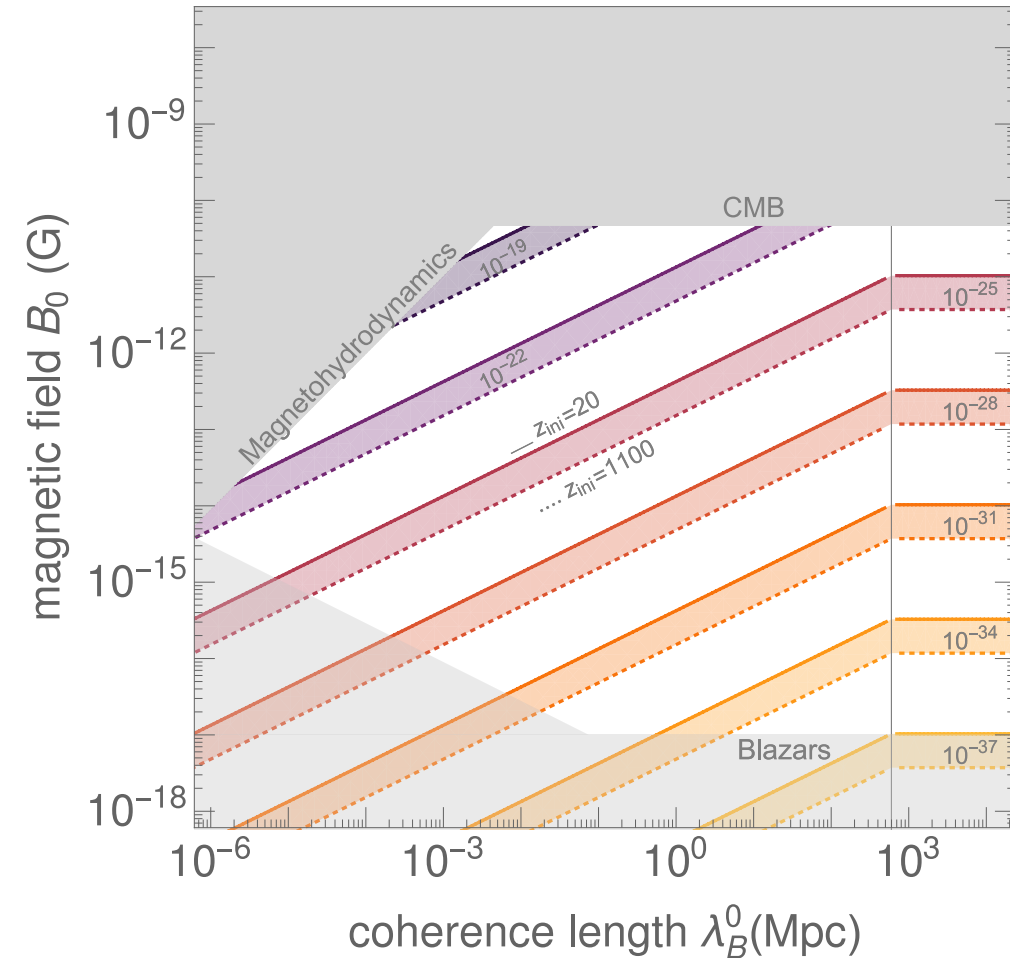


Camilo A. Garcia Cely (DESY)



# Upper bounds on stochastic gravitational waves

Domcke, CGC 2020



existing laboratory bounds from

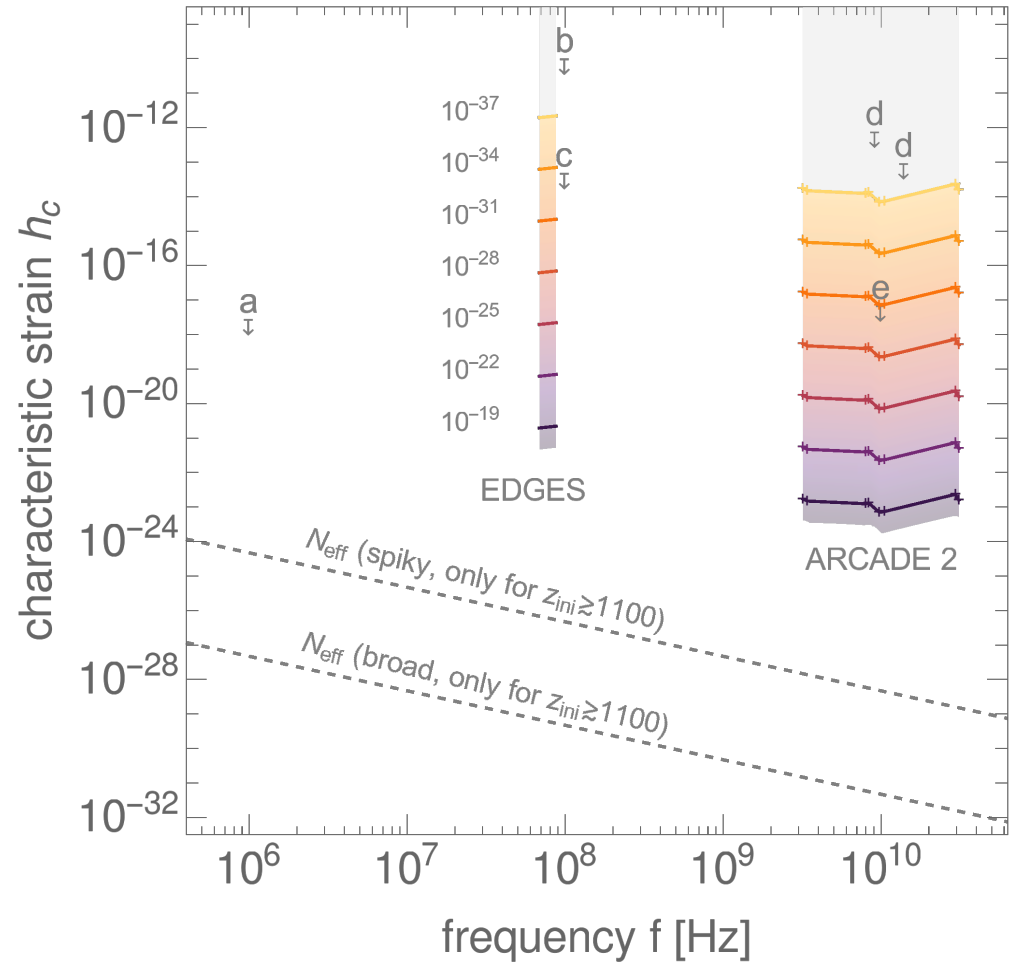
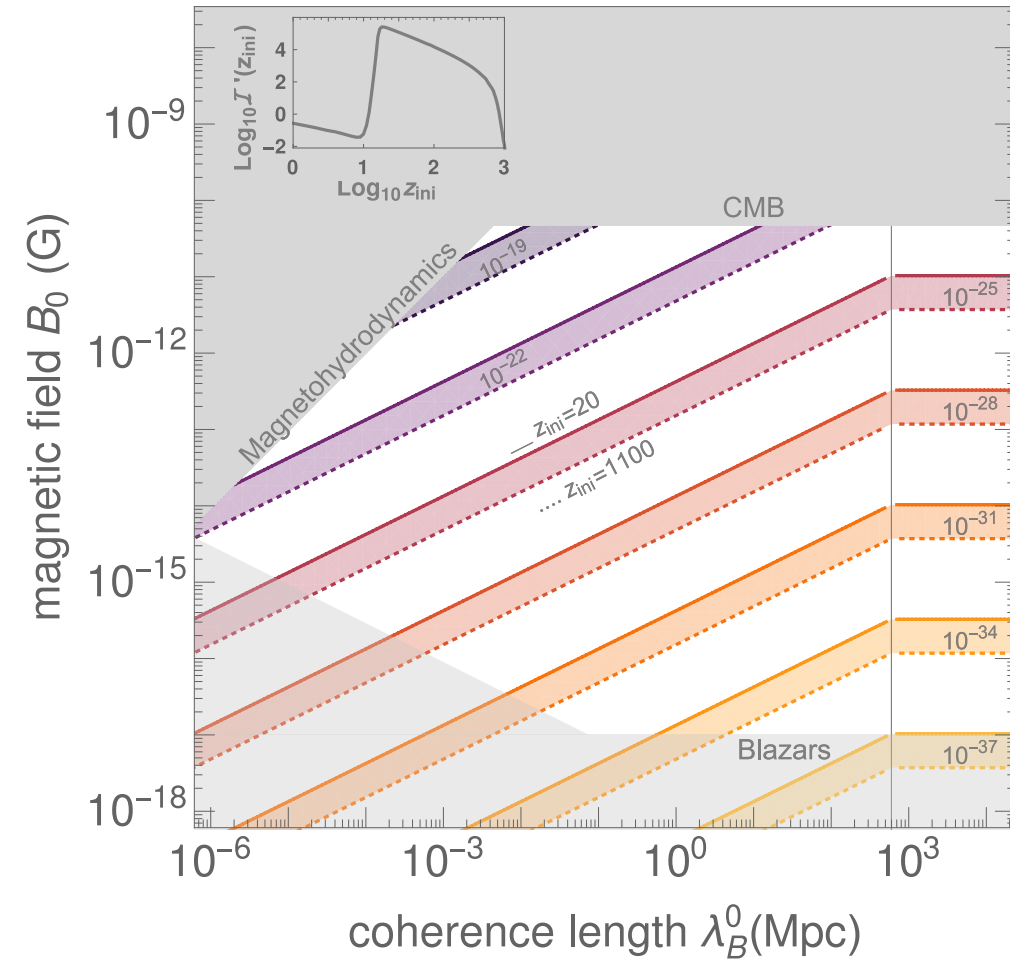
- a) superconducting parametric converter Reece et al '84
- b) waveguide Cruise Ingley '06
- c) 0.75 m interferometer Akutsu '08
- d) magnon detector Ito, Soda '04
- e) magnetic conversion detector Cruise et al '12

The CMB as a detector of high-frequency GWs

Camilo A. Garcia Cely (DESY)

# Upper bounds on stochastic gravitational waves

Domcke, CGC 2020





# Conclusions

- The Gertsenshtein effect during the dark ages provides a powerful way to probe GWs in the MHz-GHz range from distortions of the Rayleigh-Jeans CMB tail.
- With upcoming advances in 21cm astronomy targeting precisely this frequency range with increasing accuracy, it becomes conceivable to push the limits derived from radio telescopes below the cosmological bound constraining the total energy in GWs.

# Conclusions

- The Gertsenshtein effect during the dark ages provides a powerful way to probe GWs in the MHz-GHz range from distortions of the Rayleigh-Jeans CMB tail.
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Thank you for your attention

Wave equation  $\left(\square + \omega_{\text{pl}}^2/c^2\right) A_\lambda = -B\partial_\ell h_\lambda, \quad \square h_\lambda = \kappa^2 B\partial_\ell A_\lambda,$

Solution

$$\psi(t, \ell) \equiv \begin{pmatrix} \sqrt{\mu} A_\lambda \\ \frac{1}{\kappa} h_\lambda \end{pmatrix} = e^{-i\omega t} e^{iK\ell} \psi(0, 0),$$

$$K = \begin{pmatrix} \frac{\mu}{c} \sqrt{\omega^2 + \left(\frac{\kappa B}{1+\mu}\right)^2} & -i \frac{\sqrt{\mu} \kappa B}{1+\mu} \\ i \frac{\sqrt{\mu} \kappa B}{1+\mu} & \frac{1}{c} \sqrt{\omega^2 + \left(\frac{\kappa B}{1+\mu}\right)^2} \end{pmatrix}.$$

Conversion  
rate

$$\langle \Gamma_{g \leftrightarrow \gamma} \rangle = \frac{c |K_{12}| \ell_{\text{osc}}}{2\Delta\ell}.$$

$$\ell_{\text{osc}}^{-1} = \sqrt{\omega^2(1-\mu)^2/c^2 + \kappa^2 B^2/2}.$$

$$\mathcal{P} \equiv \int \langle \Gamma_{g \leftrightarrow \gamma} \rangle dt = \int^{z_{\text{ini}}} \frac{\langle \Gamma_{g \leftrightarrow \gamma} \rangle}{(1+z)H} dz,$$

CMB distortions

$$\delta f_\gamma(\omega_0, T_0) = (f_g(\omega_{\text{ini}}, T_{\text{ini}}) - f_{\text{eq}}) \mathcal{P} + \mathcal{O}(\mathcal{P}^2),$$

$$\frac{\delta f_\gamma}{f_\gamma}(\omega_0, T_0) = \frac{\pi^4}{15} \left(\frac{T}{\omega}\right)^3 \mathcal{P} \frac{\Omega_{\text{GW}}}{\Omega_\gamma} \quad \text{for } \omega \ll T.$$

$$h_c = \left( \frac{3H_0^2}{4\pi^2} \Omega_{\text{GW}} f^{-2} \right)^{1/2}.$$