# Dark Matter Indirect Detection with Neutrino Lines

#### Camilo Garcia Cely, ULB



 ${\sf Da} \nu {\sf Co}$  Dark Matter,  $\nu {\sf s}$  and their Connections

Aug 29th, 2017

In collaboration with Chaimae El Aisati, Thomas Hambye, Julian Heeck and Laurent Vanderheyden.

Based on arXiv:1706.06600 and JCAP 1703 (2017) no.03, 054

#### Indirect dark matter detection:

Dark matter annihilates or decays into some particles and these in turn produce a flux of  $\gamma$ ,  $e^{\pm}$ , p,  $\overline{p}$ , and (anti-)neutrinos. Subsequently, these propagate from the point where they are produced until they reach the earth.

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#### This talk: DM $\rightarrow \nu \nu$ or DM DM $\rightarrow \nu \nu$ .

- Neutrinos point to the direction where they come from.
- Neutrinos are not subject to energy loses → The observation of a line would allow to infer the DM mass.
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#### What sort of models naturally lead to neutrino lines?

- multi-TeV DM Dark Matter: Thomas Hambye's Talk.
- Majoron DM below a few TeV. CGC, Heeck (2017)

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- Below a few TeV, Majorons decay and mainly produce monochromatic neutrinos.
  - Above a few TeV, three-body and four-body final states dominate the decay. Dudas, Mambrini and Olive (2015).

• Consider  $\sigma = \frac{f + \sigma^0 + iJ}{\sqrt{2}}$  with two units of lepton number:

$$\mathcal{L} = \underbrace{-\overline{L}yN_RH}_{\text{Dirac Mass Term}} - \underbrace{\frac{1}{2}\overline{N}_R^c\lambda N_R\sigma}_{\text{Majorana Mass Term}} + \text{h.c.}$$

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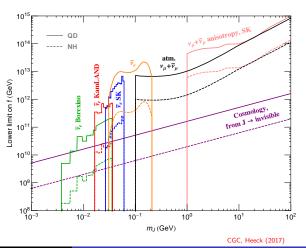
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- Sub-GeV Majorons can be produced via Freeze-in

  See Danielo Teresi's

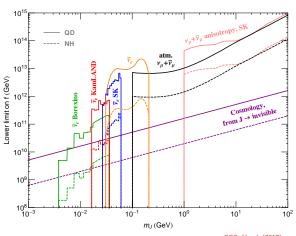
Frigerio, Hambye, Masso (2011) , Heeck, Teresi (2017). See Daniele Teresi's talk.

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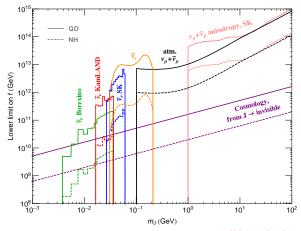
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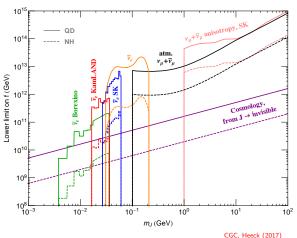
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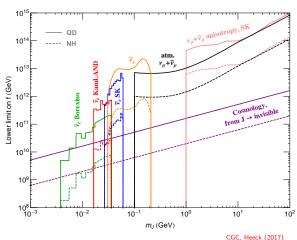
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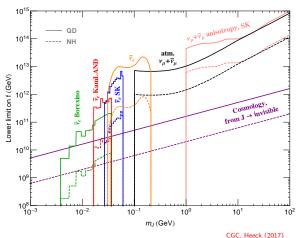
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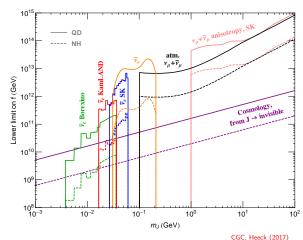
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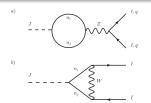


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- Neutrino experiments can be used as DM detectors



# Decay into charged fermion pairs



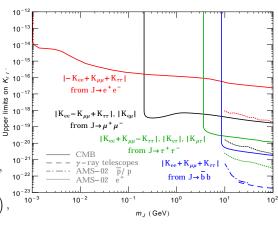
$$\mathcal{L}_{J} = iJ\bar{f}_{1}(g_{Jf_{1}f_{2}}^{S} + g_{Jf_{1}f_{2}}^{P}\gamma_{5})f_{2}$$

$$g^P_{Jqq'} \simeq \frac{m_q}{8\pi^2 v} \delta_{qq'} \, T_3^q \, \text{tr} K \, , \quad g^S_{Jqq'} = 0 \, , \label{eq:gqq}$$

$$g_{J\ell\ell'}^P \simeq rac{m_\ell + m_{\ell'}}{16\pi^2 v} \left( \delta_{\ell\ell'} \, T_3^\ell \, {
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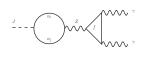
$$g_{J\ell\ell'}^S \simeq \frac{m_{\ell'}-m_\ell}{16\pi^2 v} K_{\ell\ell'}$$
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$$K \equiv \frac{m_D m_D^{\dagger}}{vf}$$



CGC, Heeck (2017)

# Decay into photons (much more difficult)



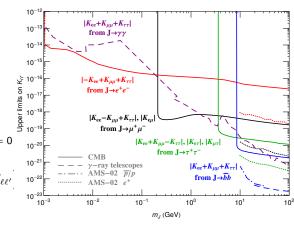
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$$g_{J\ell\ell'}^S \simeq \frac{m_{\ell'}-m_\ell}{16\pi^2 v} K_{\ell\ell'} \,, \label{eq:global_state}$$

$$K \equiv \frac{m_D m_D^{\dagger}}{v_f}$$



CGC, Heeck (2017)

### Astrophysical neutrinos (produced as flavor eigenstates)

$$\rho^{S} = \overbrace{\begin{pmatrix} \alpha_e^S & 0 & 0 \\ 0 & \alpha_\mu^S & 0 \\ 0 & 0 & \alpha_\tau^S \end{pmatrix}}_{\text{flavor-eigenstate basis}},$$

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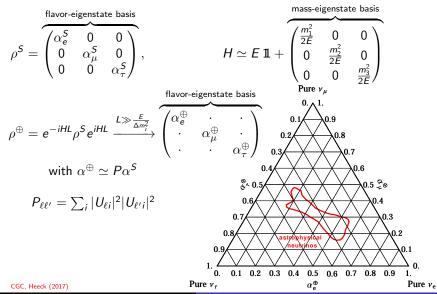
flavor-eigenstate basis

$$\rho^{\oplus} = e^{-iHL} \rho^{S} e^{iHL} \xrightarrow{L \gg \frac{E}{\Delta m_{i}^{2}}} \overbrace{\begin{pmatrix} \alpha_{e}^{\oplus} & \cdot & \cdot \\ \cdot & \alpha_{\mu}^{\oplus} & \cdot \\ \cdot & \cdot & \alpha_{\tau}^{\oplus} \end{pmatrix}}^{C}$$

with 
$$\alpha^{\oplus} \simeq P\alpha^{S}$$

$$P_{\ell\ell'} = \sum_i |U_{\ell i}|^2 |U_{\ell' i}|^2$$

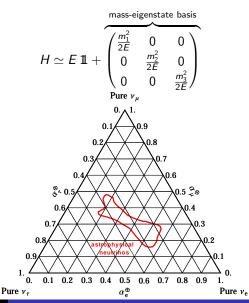
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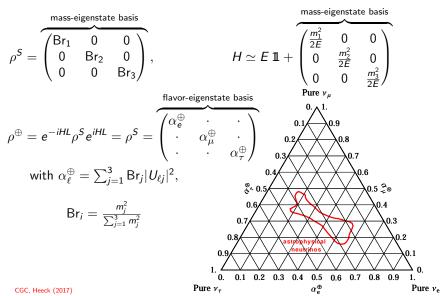
#### Neutrinos from Majoron decay

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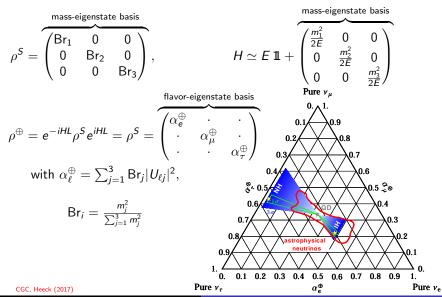
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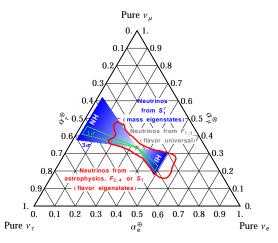


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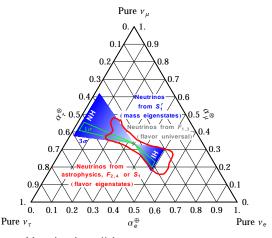
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- Process DM DM  $\rightarrow \nu\nu$
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More on this in Thomas Hambye's talk!

### Conclusions

- Majoron DM is a consistent model that provides neutrino lines as its most important signature.
- Neutrino experiments can be used as DM detectors.
- Another signature: the flavor structure of the lines can not be mimicked by astrophysical neutrinos in certain cases.

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Thanks for your attention!