

# Polynomial spectral features from dark matter

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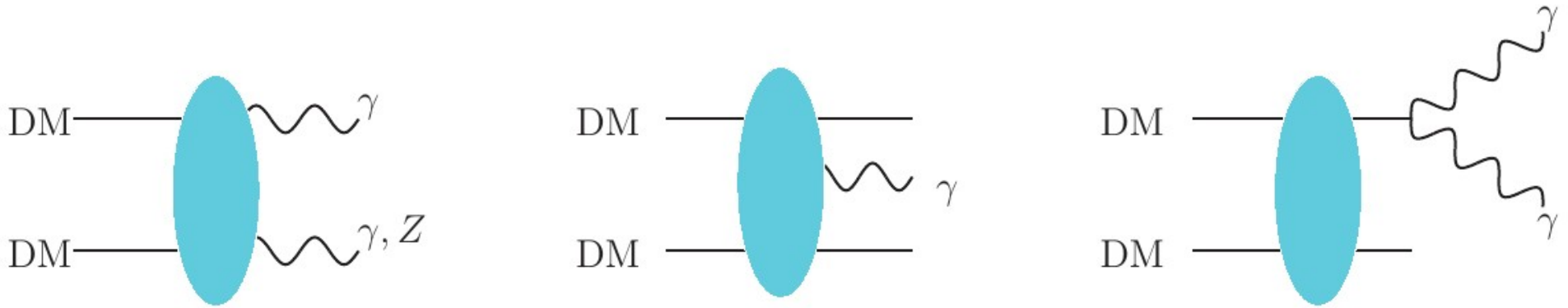
Based on arXiv:1605.08049. In Collaboration with Julian Heeck.

# Outline

- Part I : Motivation  
Box-shaped gamma-ray spectra
- Part II: Another example  
Neutrino features from DM annihilating into SM gauge bosons
- Part III: General case  
Polynomial spectral features
- Part IV: Connection to the diphoton resonance
- Conclusions

# Gamma-ray spectral features

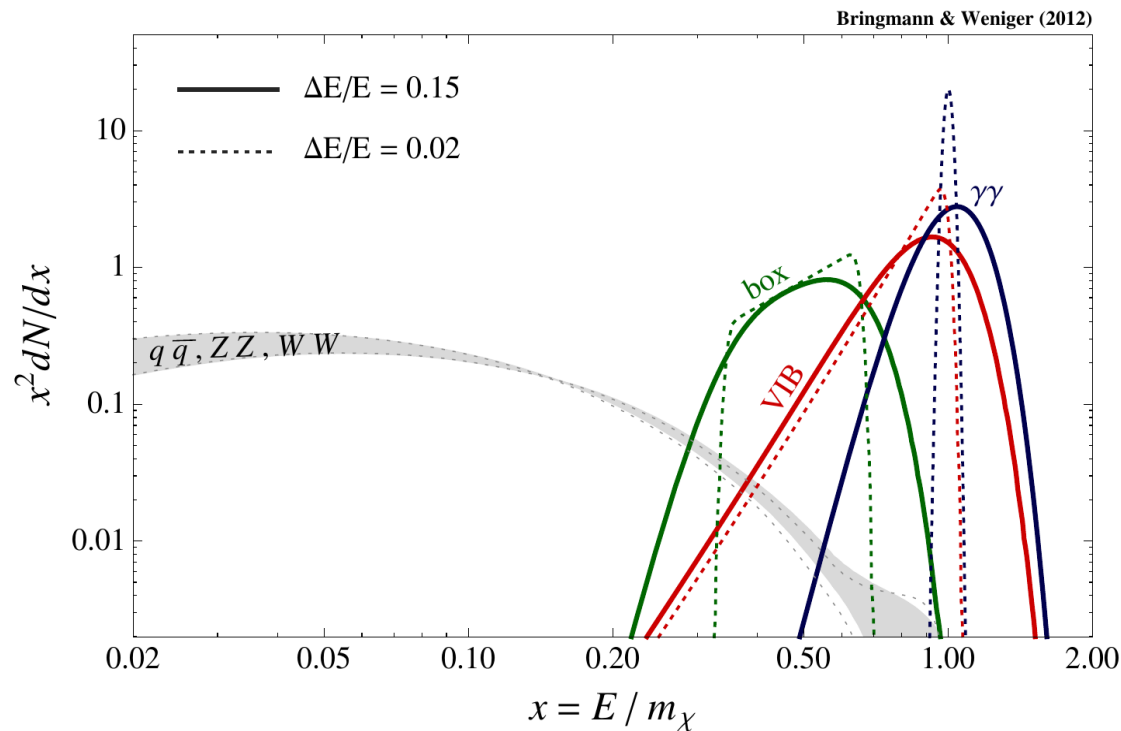
Smoking gun signature for dark matter : no astrophysical process is known to produce a sharp feature in the gamma-ray spectrum



Annihilation into Photons

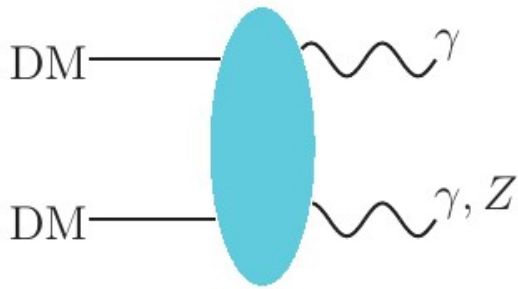
Virtual Internal Bremsstrahlung (VIB)

Box-shaped spectra

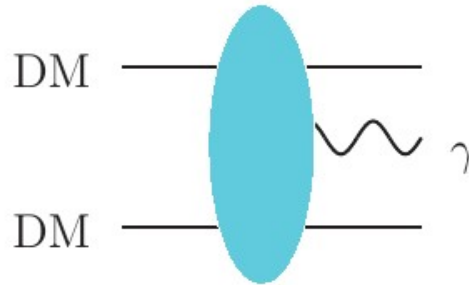


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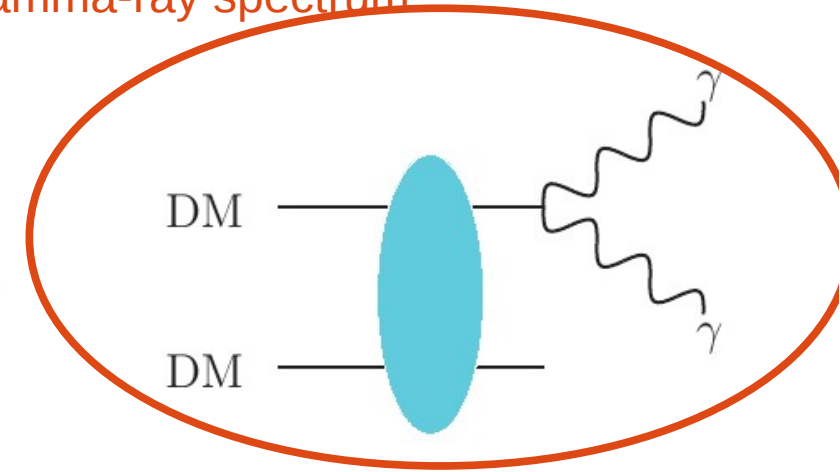
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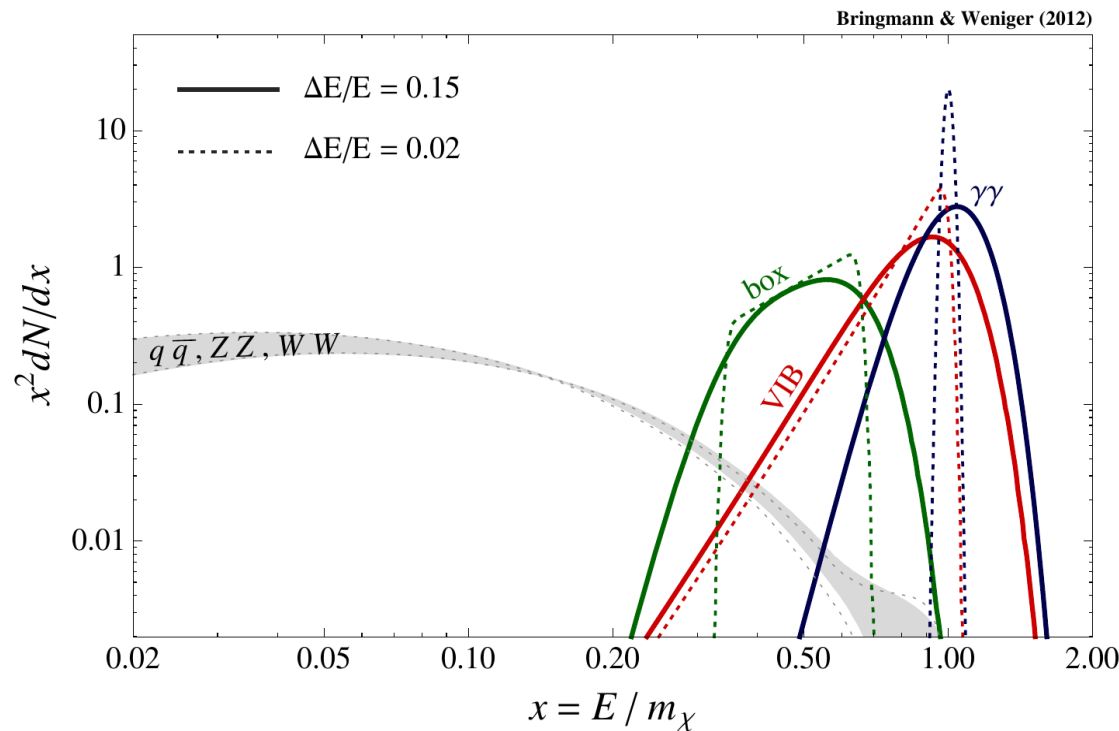
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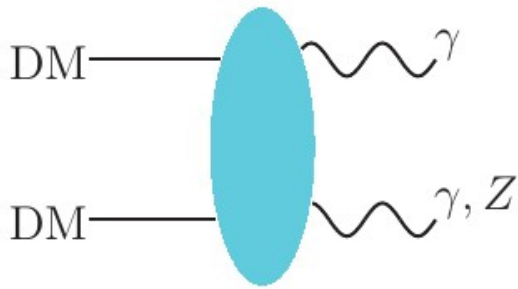
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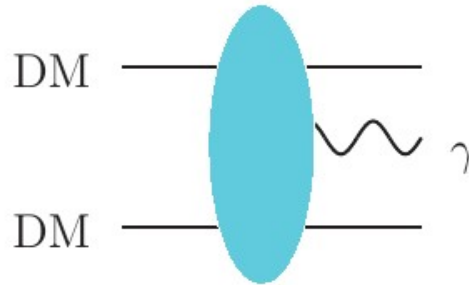
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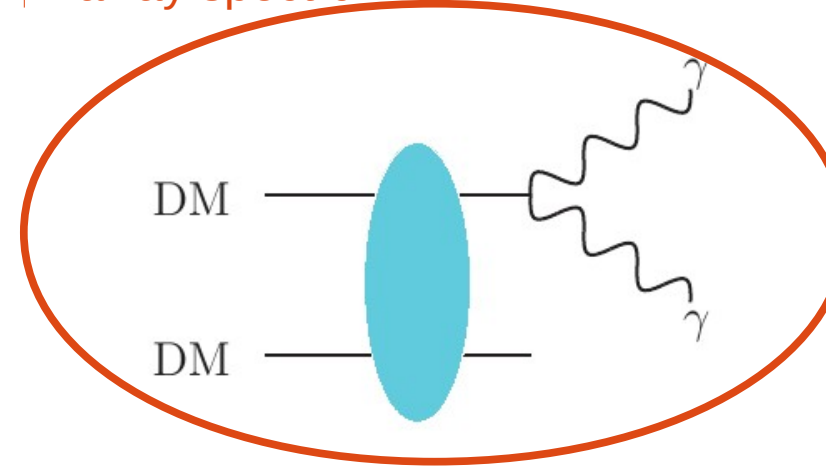
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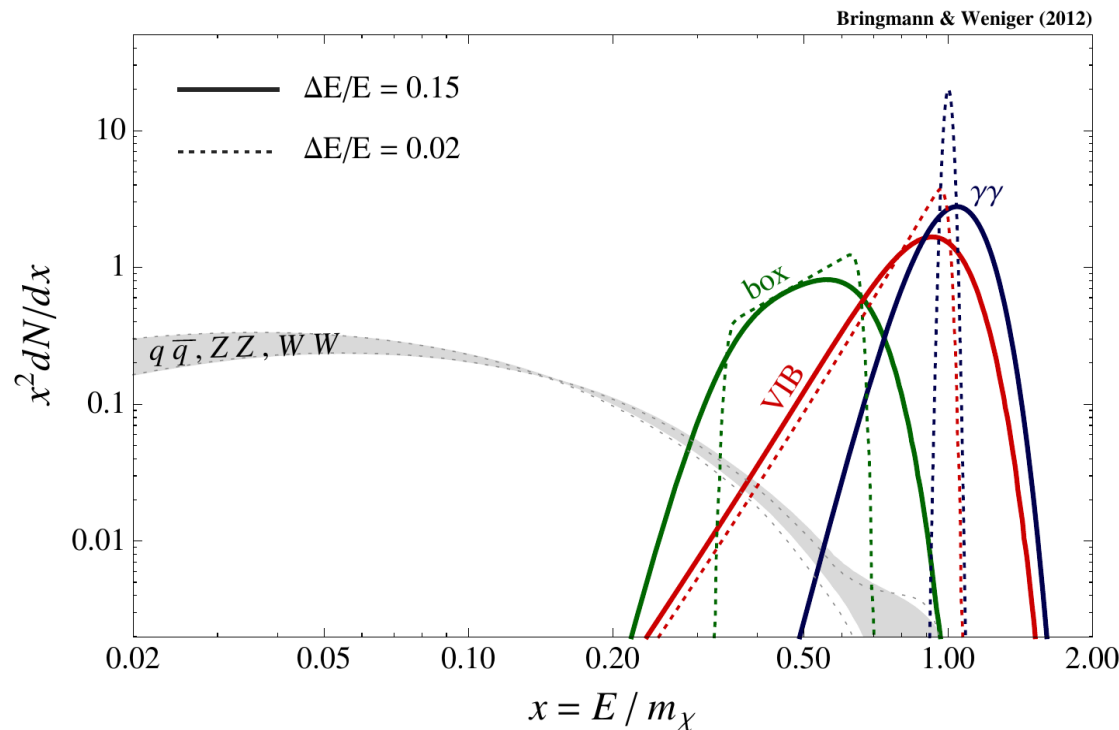
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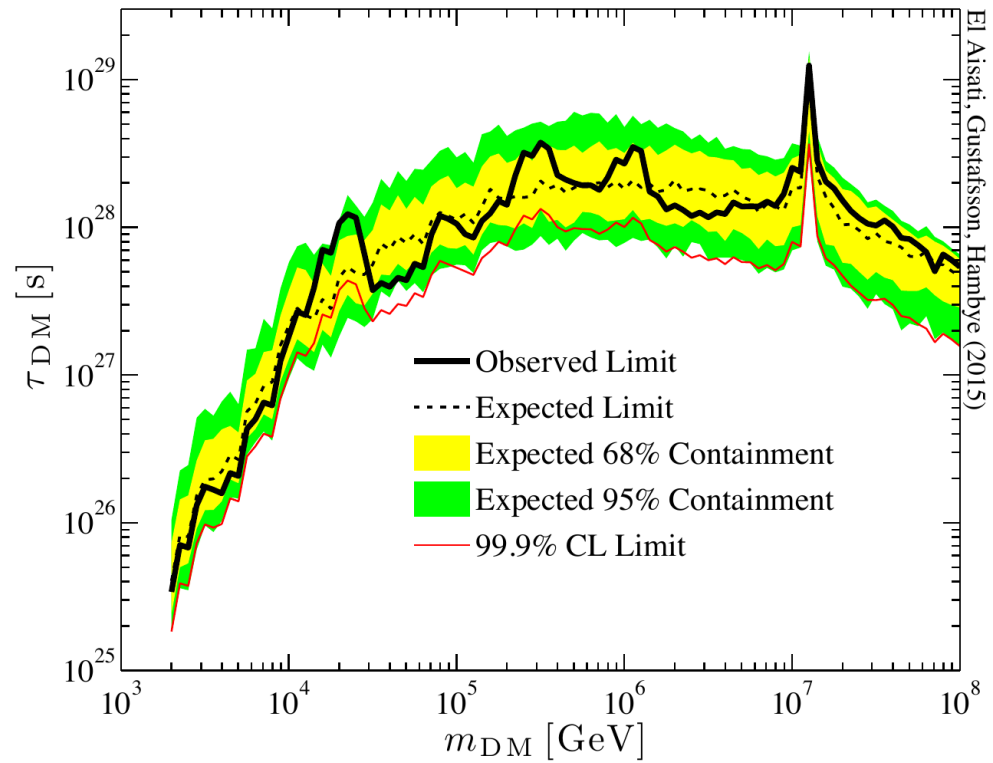
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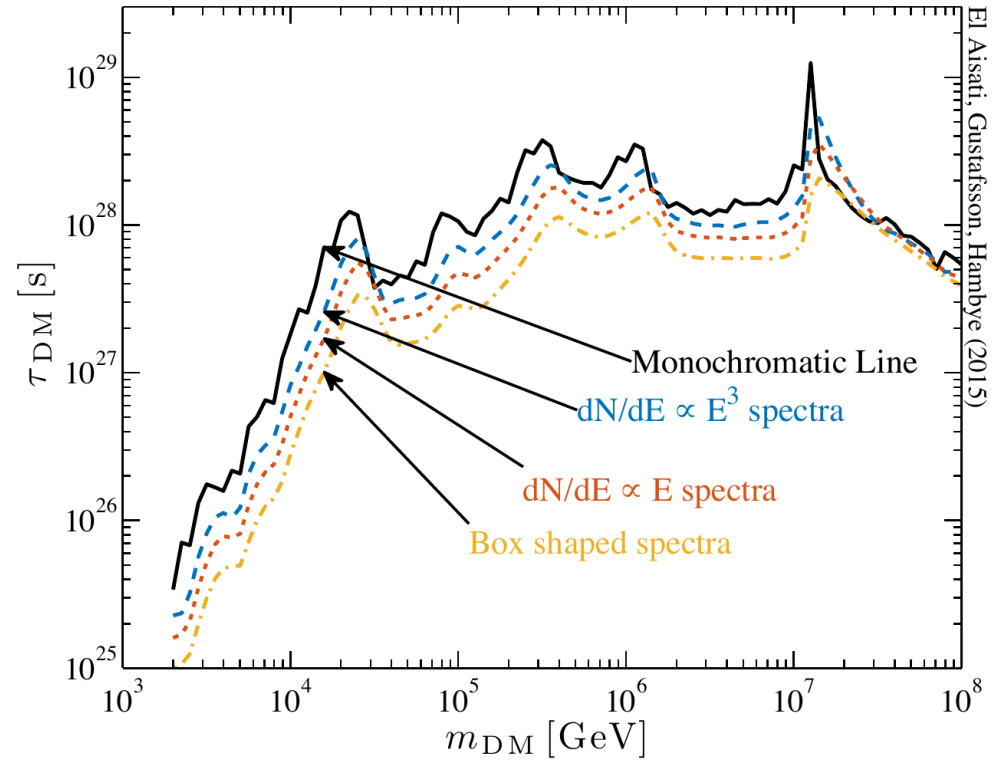
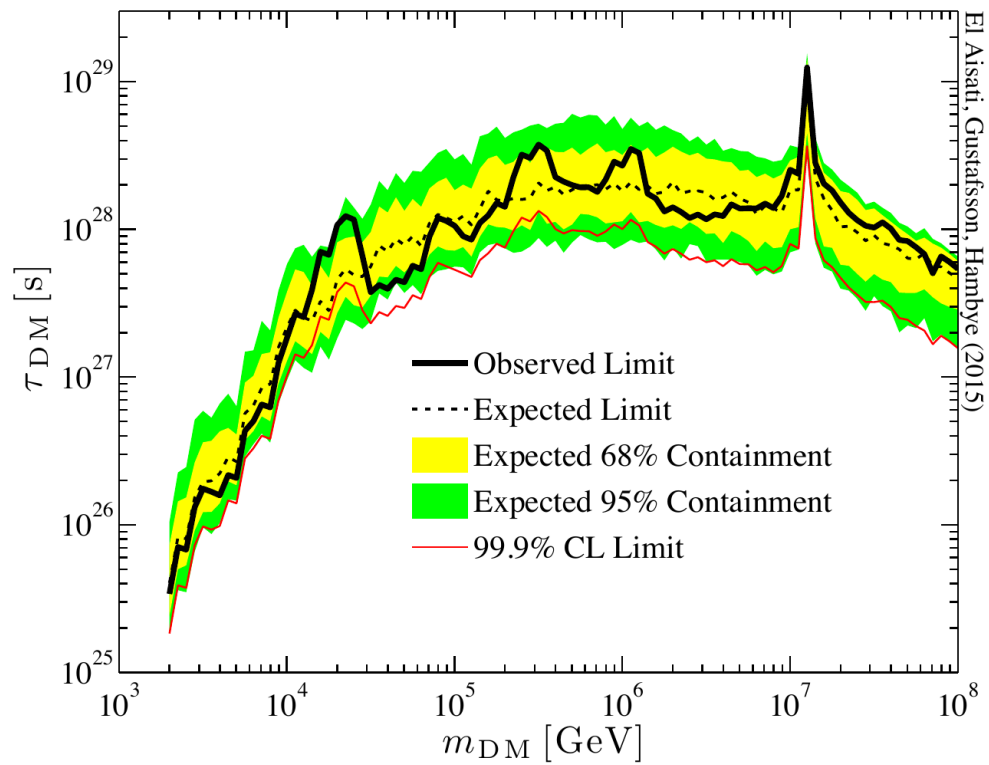
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This talk: Generalize this to an arbitrary intermediate state

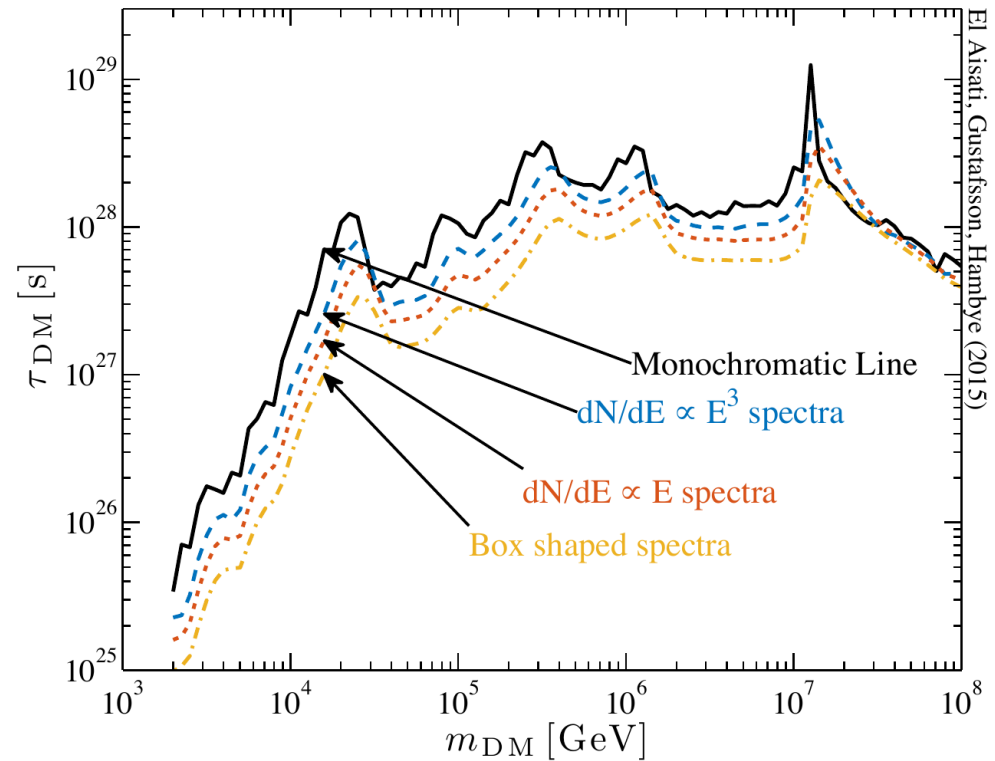
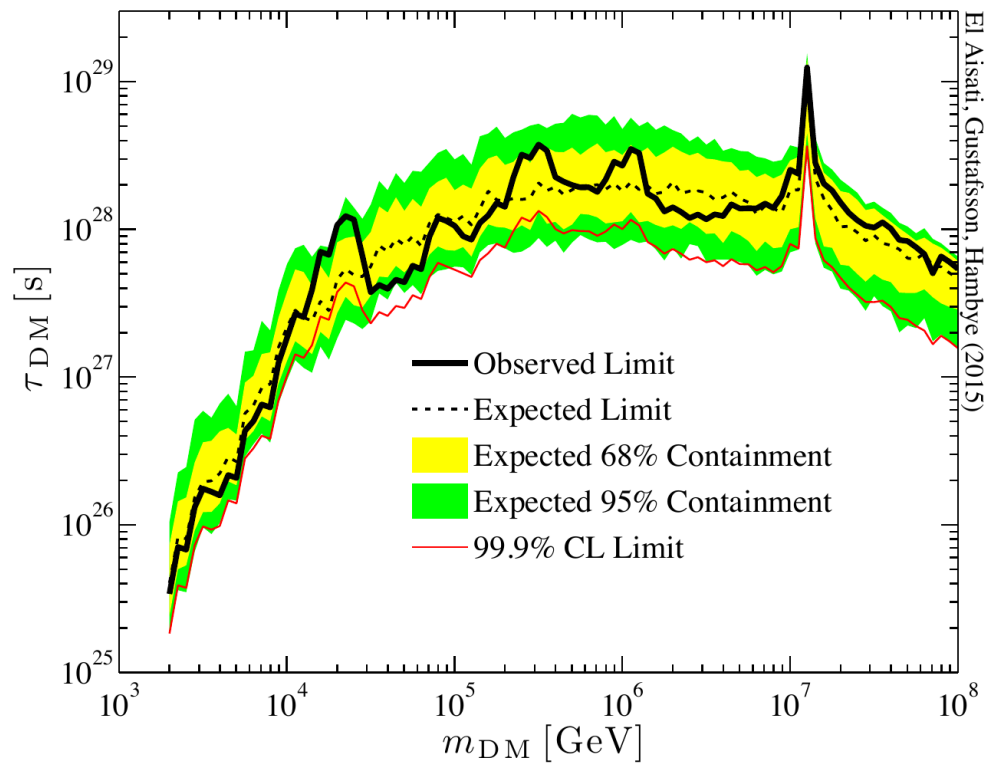
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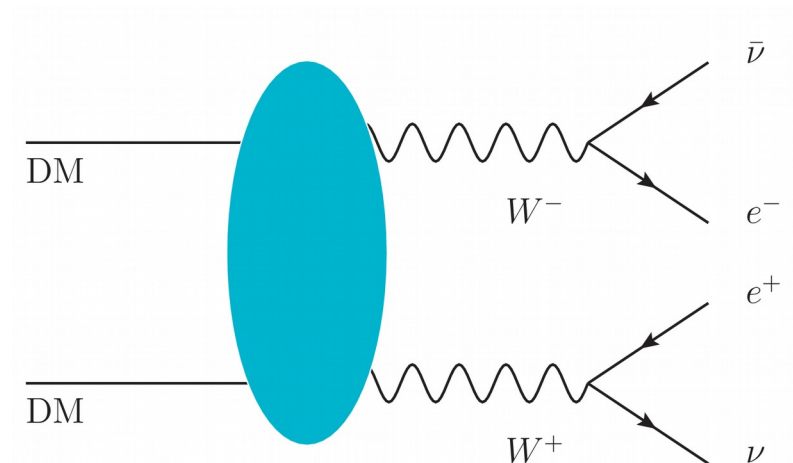
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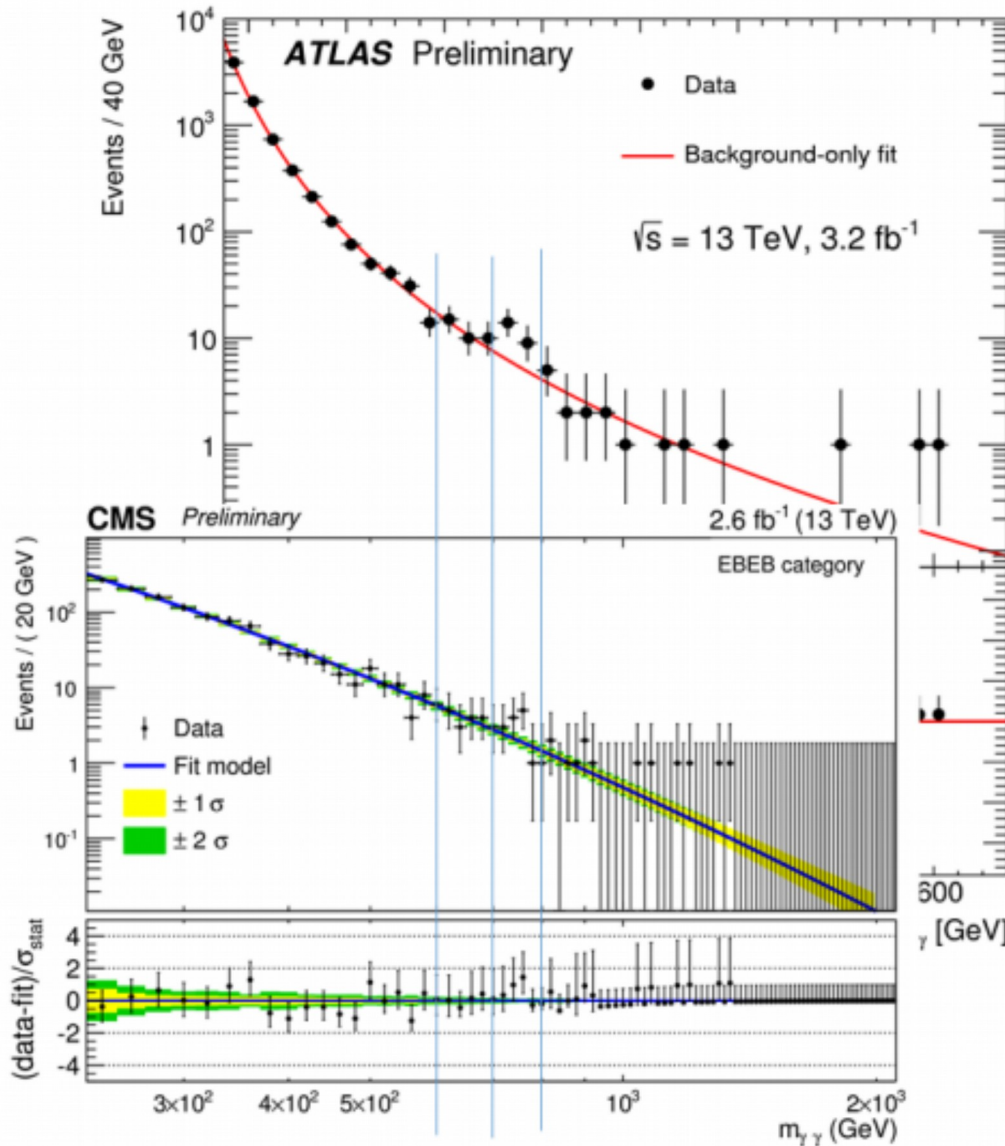


What about?





# Connection to the diphoton resonance

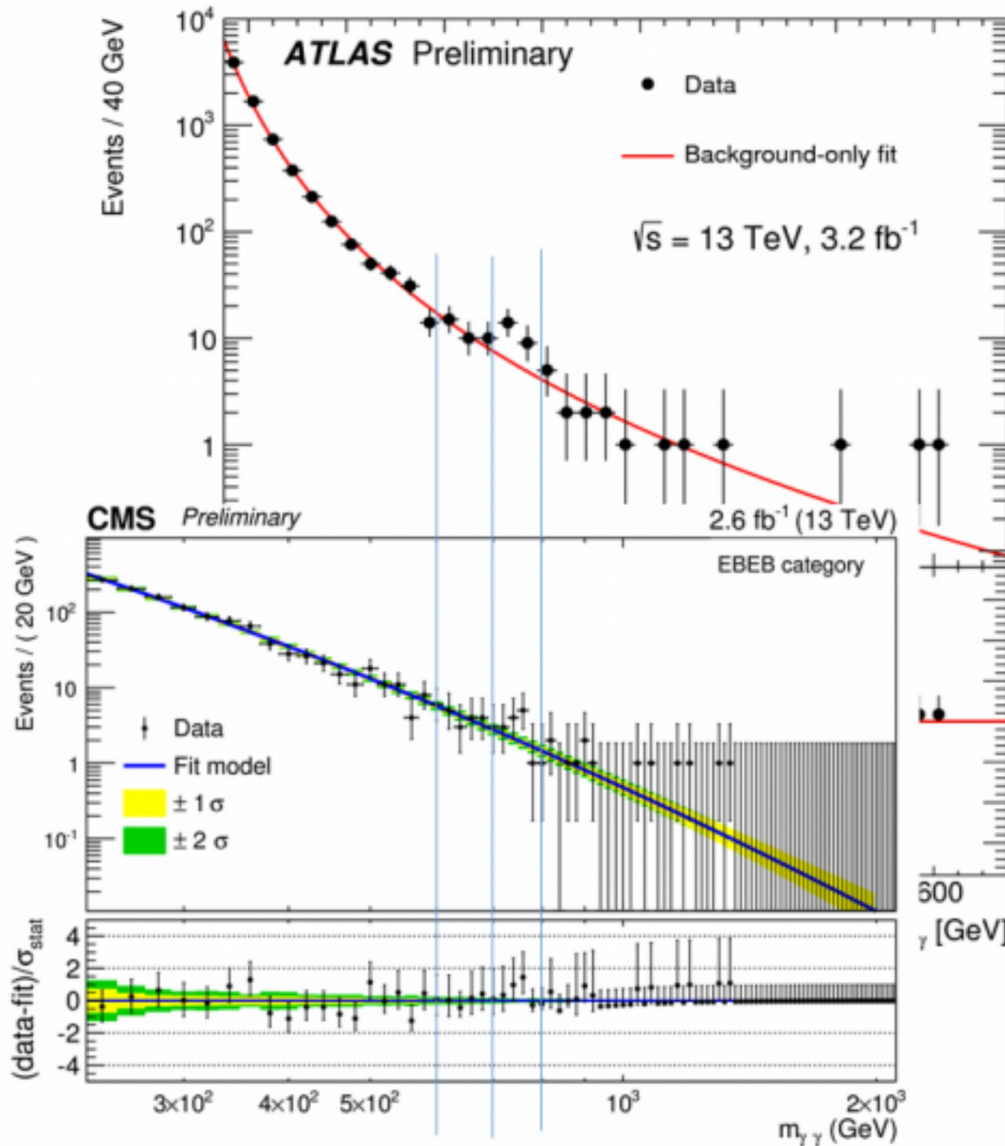


Its simplest explanation assumes a spin-0 or spin-2 particle  $R$  of mass 750 GeV

Landau-Yang theorem

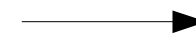
Spin-one is not possible

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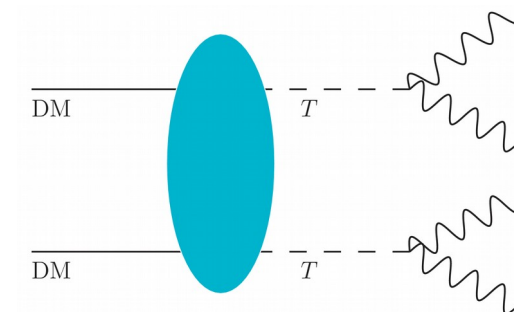
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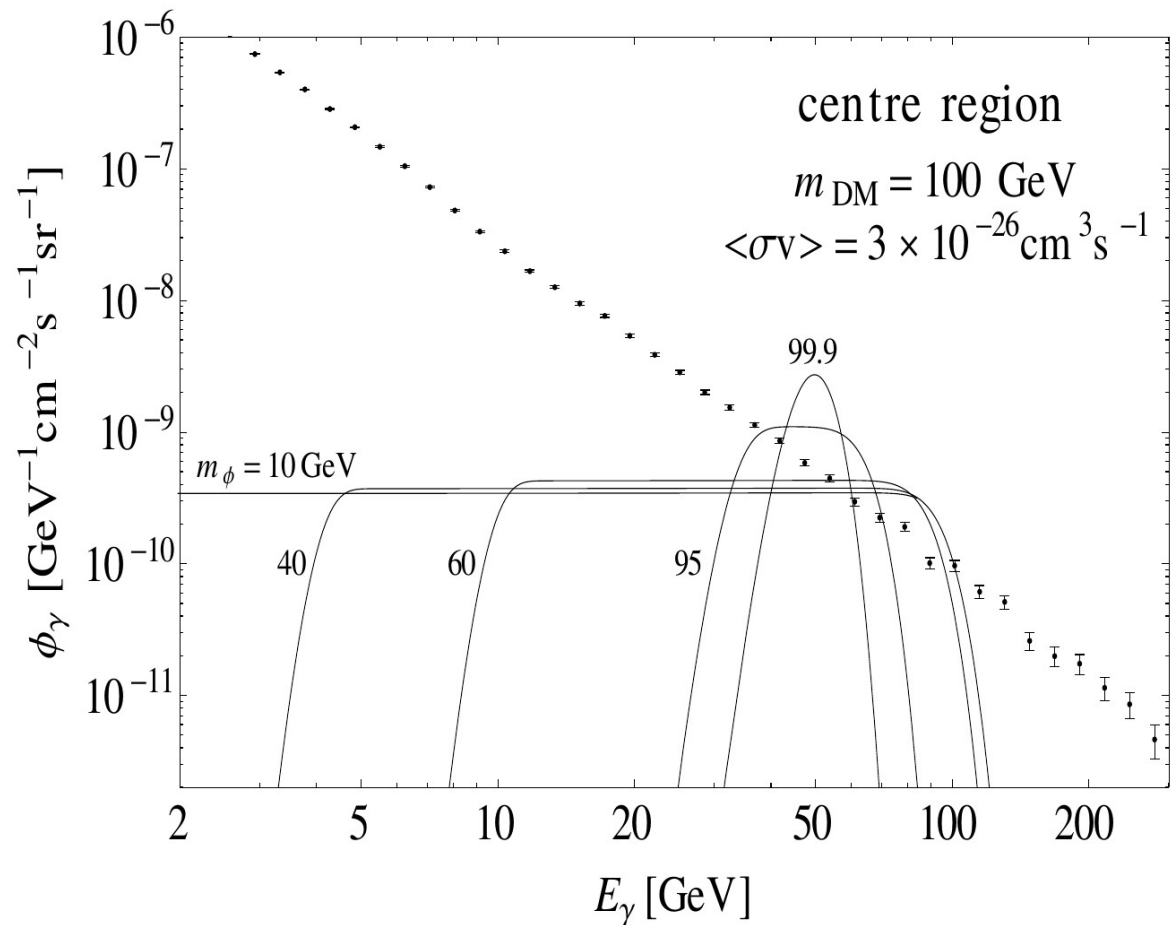
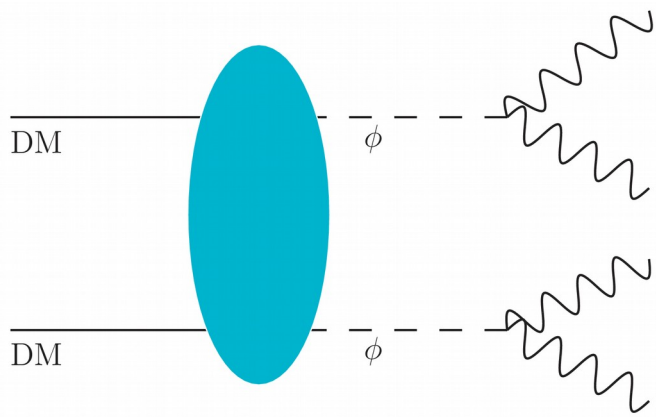


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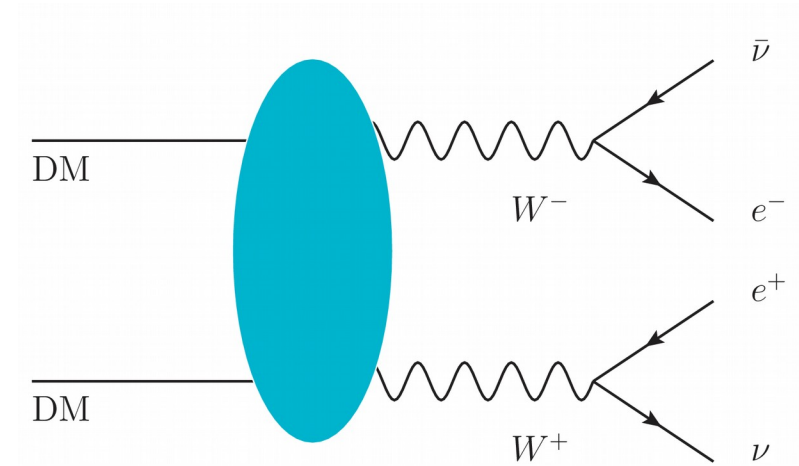
TeV-DM would likely annihilate or decay into the resonance



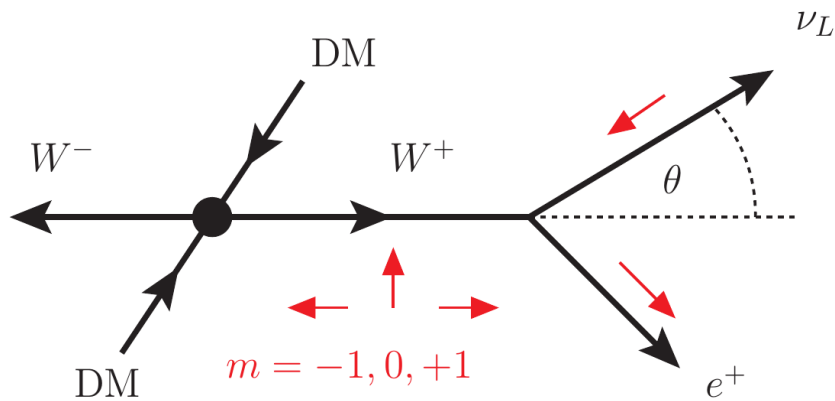
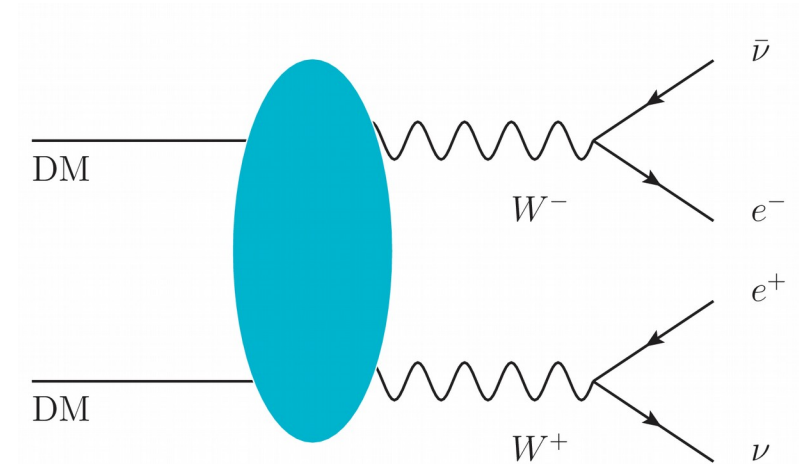
# Box-shaped spectra from intermediary scalars



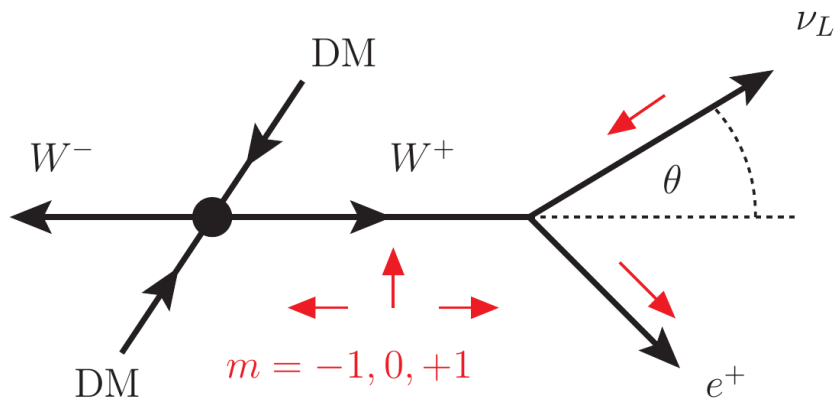
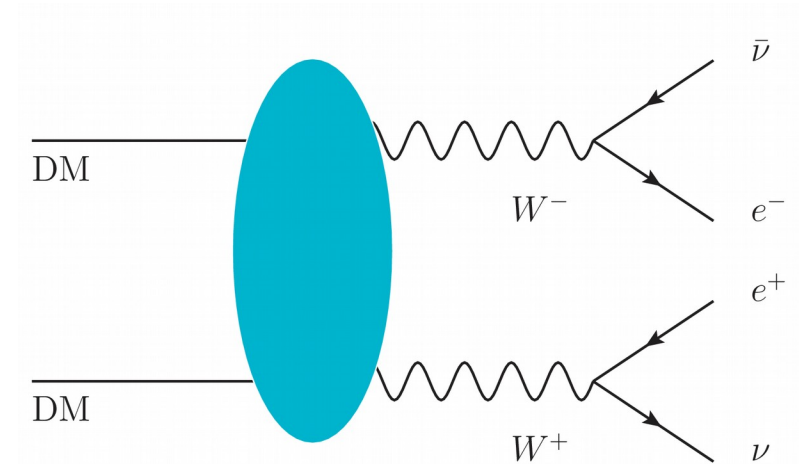
# The case of gauge bosons



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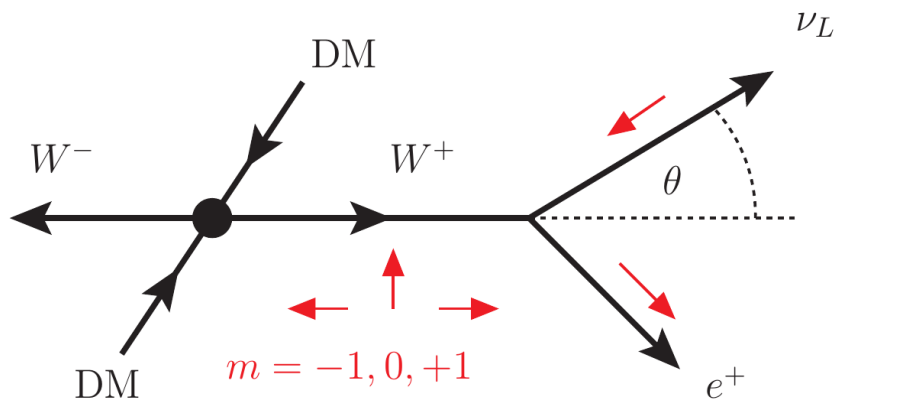
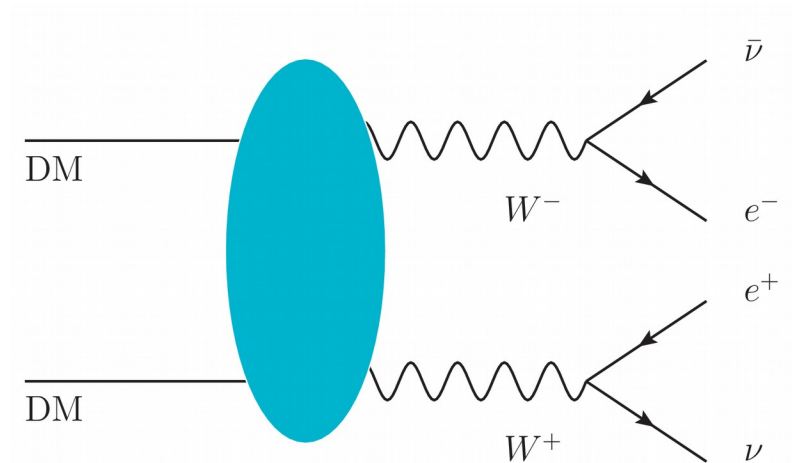


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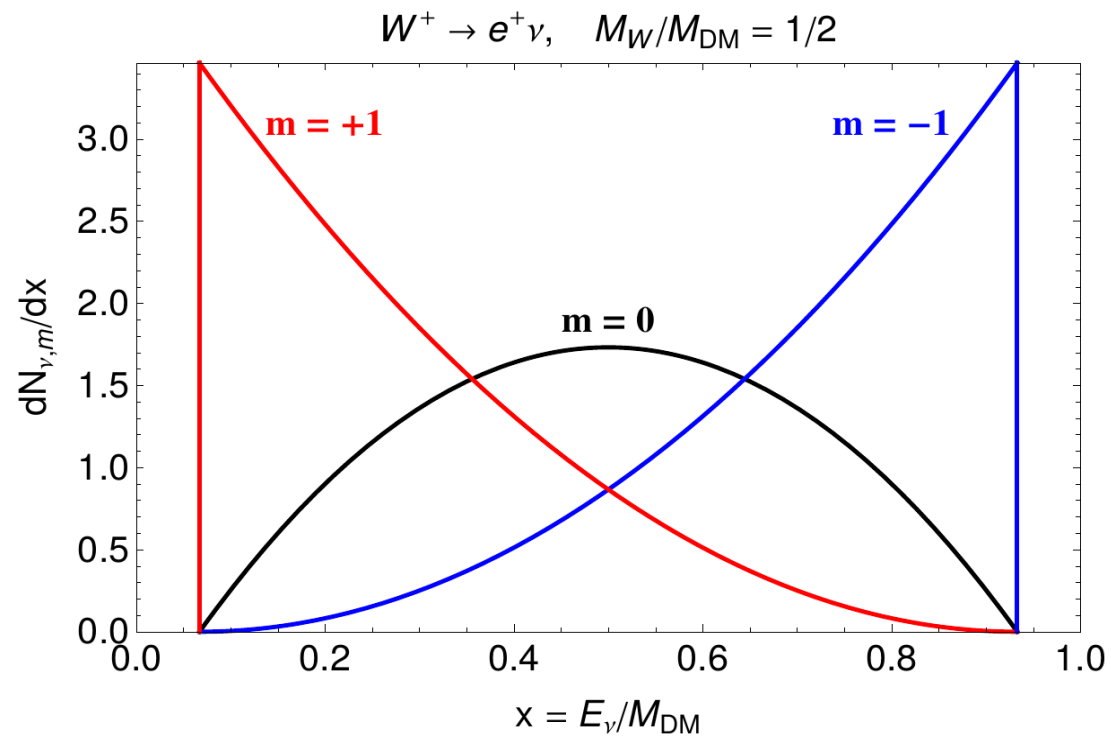


$$\frac{d\Phi_\nu}{dE_\nu} = \Phi_\nu \sum_m \text{Br}_m \frac{dN_{\nu,m}}{dE_\nu}, \quad \Phi_\nu = \frac{(\sigma v)}{8\pi M_{\text{DM}}^2} \bar{J}_{\text{ann}}$$

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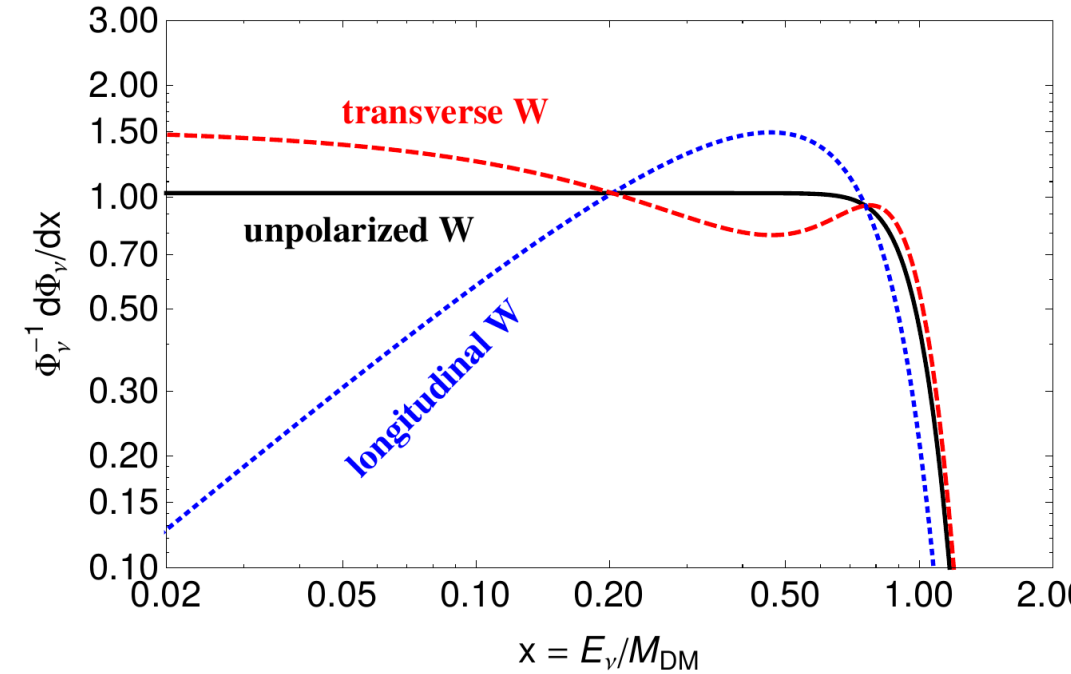


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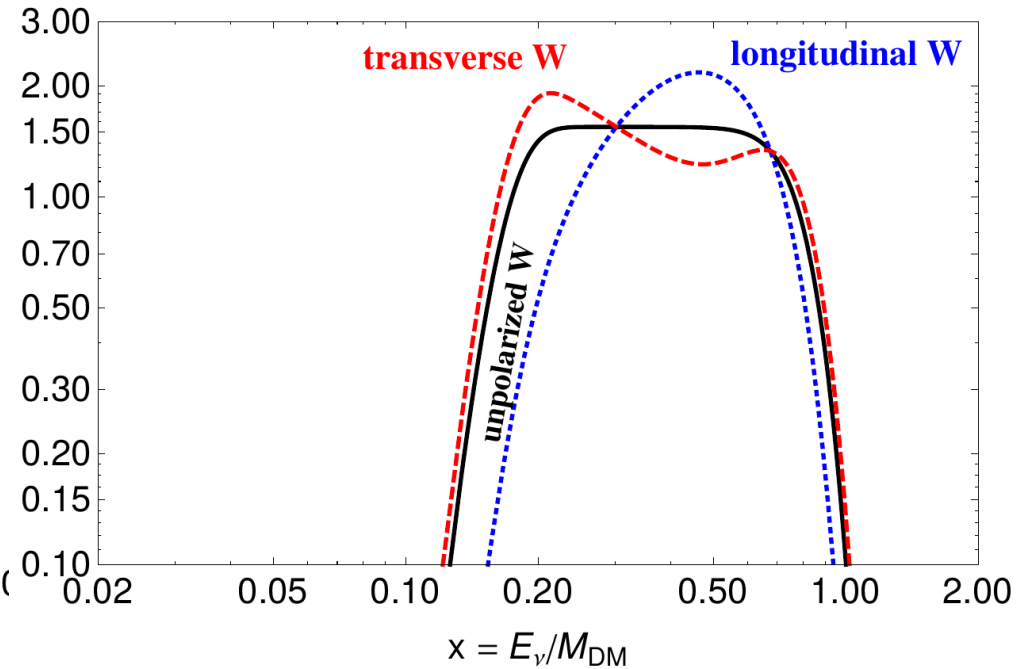


# The case of gauge bosons

DM DM  $\rightarrow W^-W^+$ ,  $W^+ \rightarrow e^+\nu$  with  $M_W/M_{\text{DM}} = 0.01$

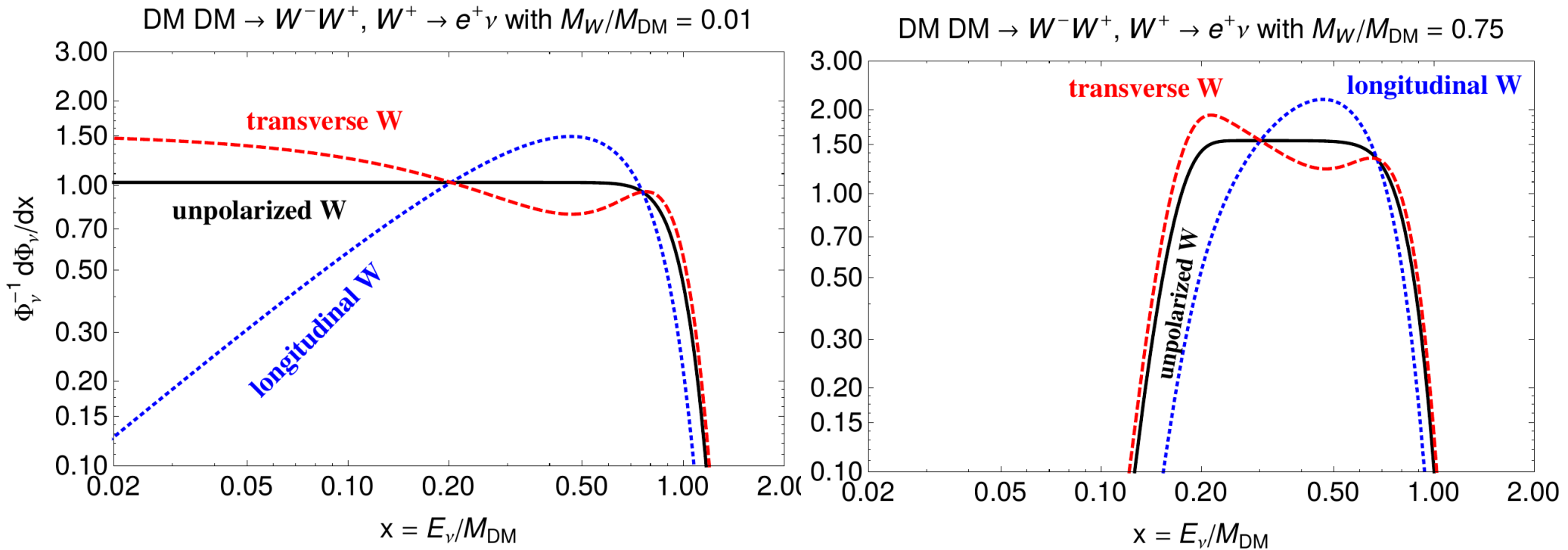


DM DM  $\rightarrow W^-W^+$ ,  $W^+ \rightarrow e^+\nu$  with  $M_W/M_{\text{DM}} = 0.75$





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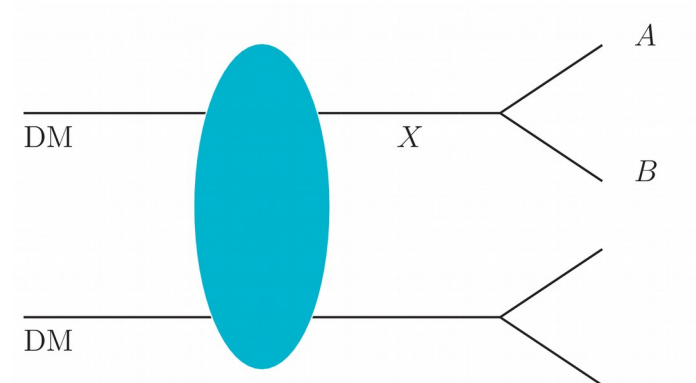


Gauge bosons produced in DM annihilations are typically polarized

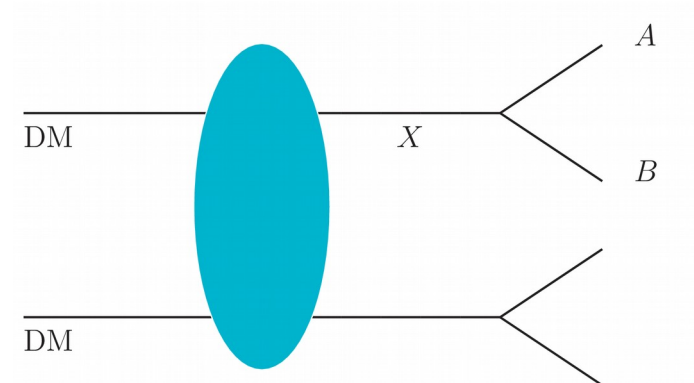
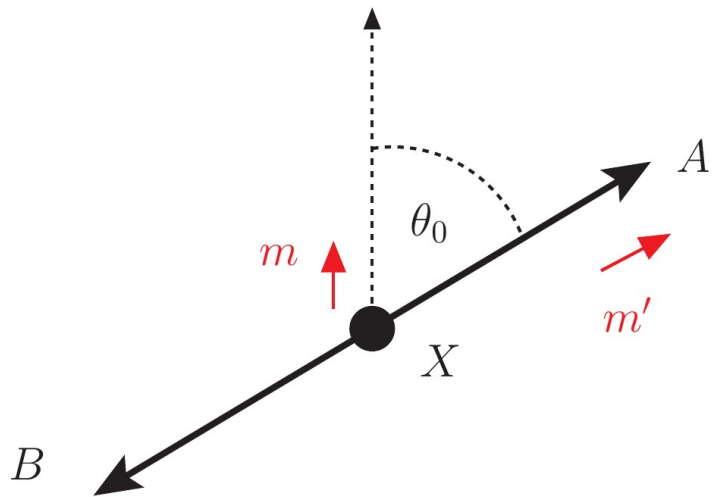
Above the electroweak scale, Majorana DM with  $SU(2)_L$  quantum numbers produce gauge bosons that are mostly transverse.

Scalar DM, also singlet under  $SU(2)_L$ , produces gauge bosons that are mostly longitudinally polarized.

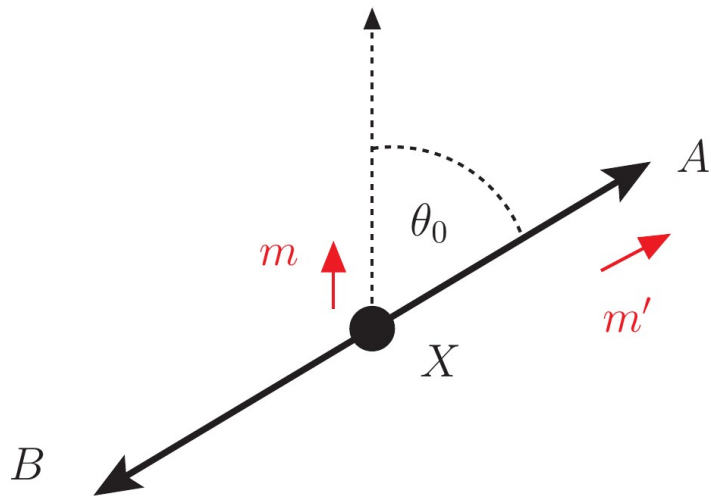
# General case



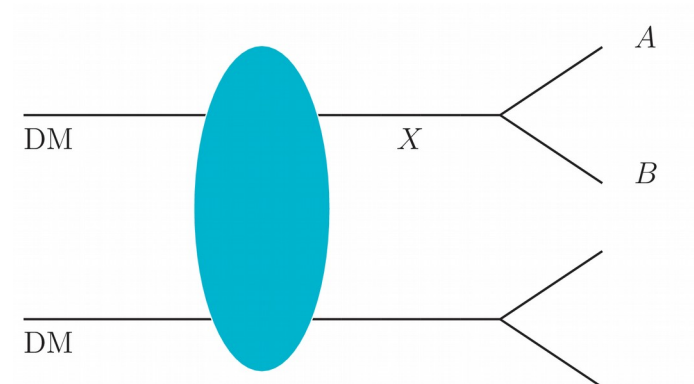
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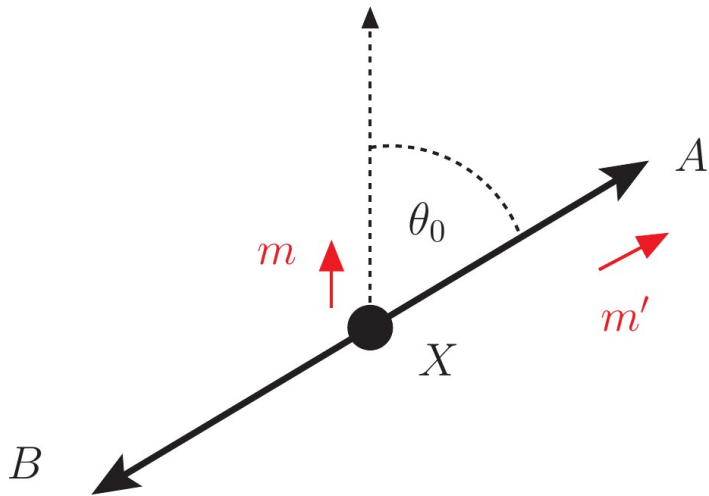
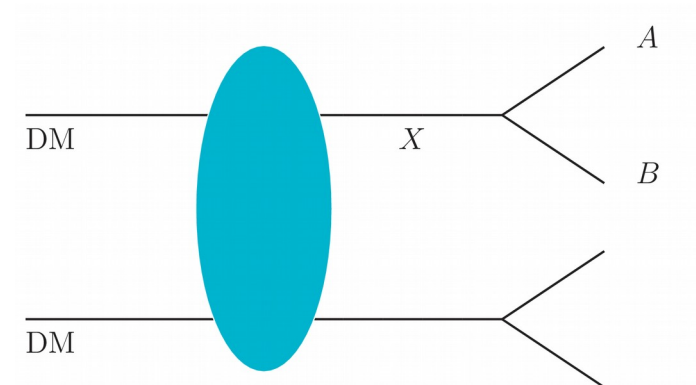
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$${}_{\theta_0} \langle m', S | m, S \rangle = \langle m', S | R(\theta_0) | m, S \rangle \equiv d_{m'm}^S(\theta_0)$$



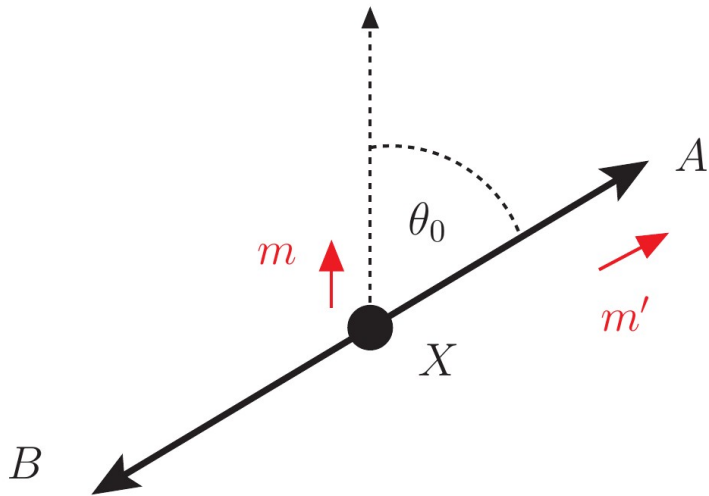
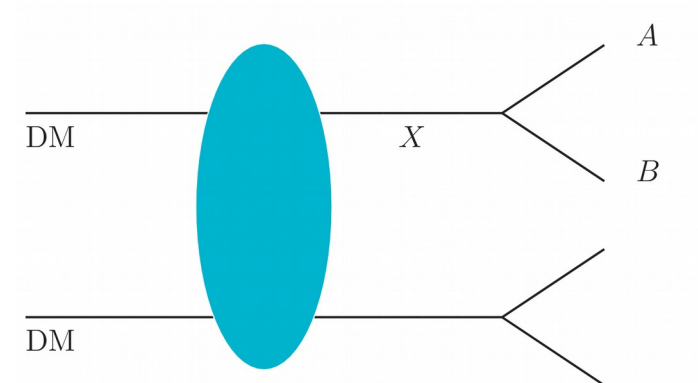
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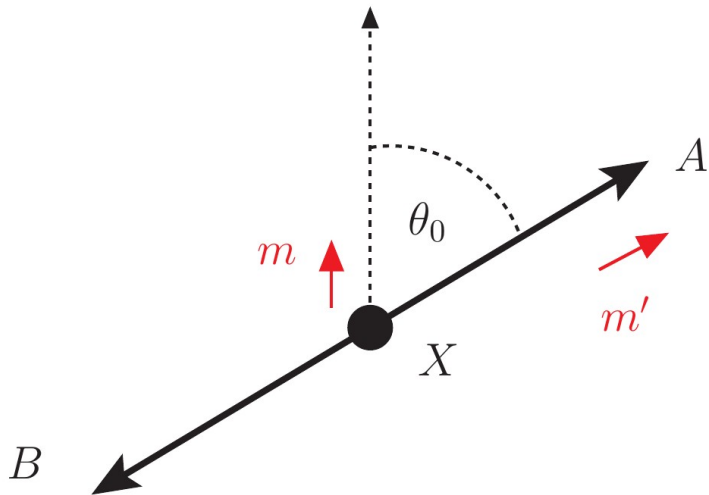


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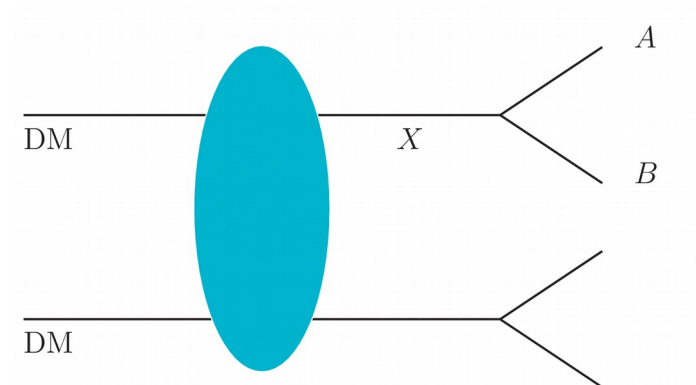
$$\begin{aligned} f_m^S(x, y) &= \frac{(2S+1)}{\sqrt{y^2 - r_X^2}} \Theta(x - x^-(y)) \Theta(x^+(y) - x) \\ &\times \sum_{m'} C_{m'} \left| d_{m'm}^S \left( \arccos \left( \frac{2x - y}{\sqrt{y^2 - r_X^2}} \right) \right) \right|^2 \end{aligned}$$

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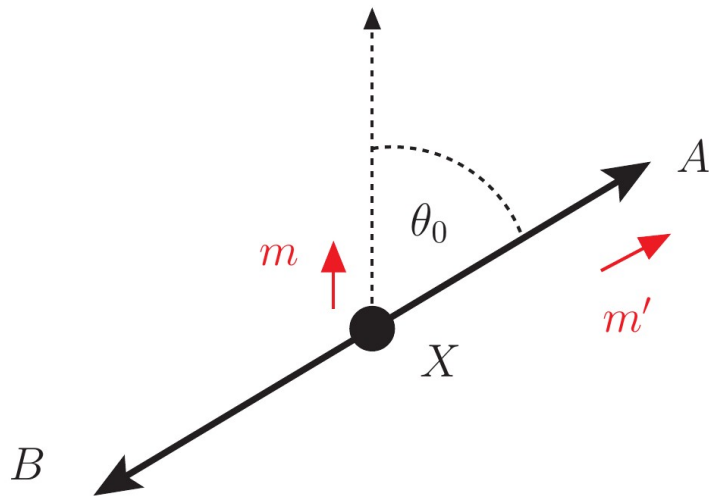


Almost everything fixed by  
angular momentum.

The dependence on the  
DM Model is encoded in two  
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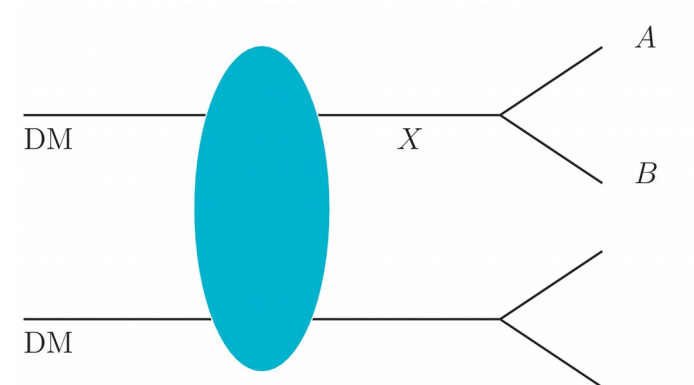
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It determines the  
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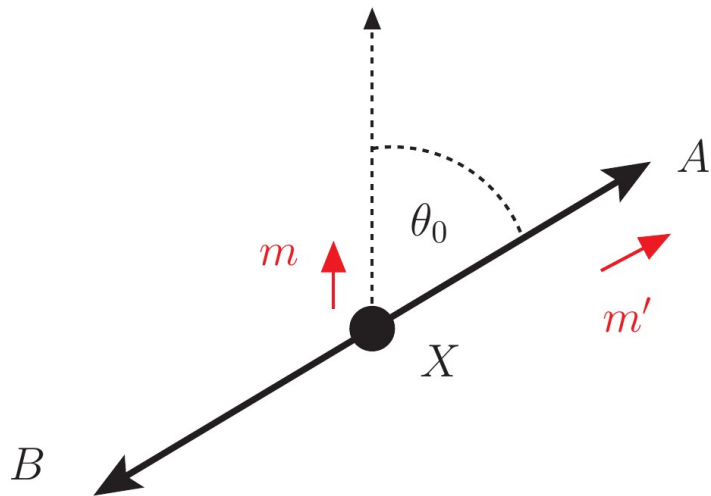
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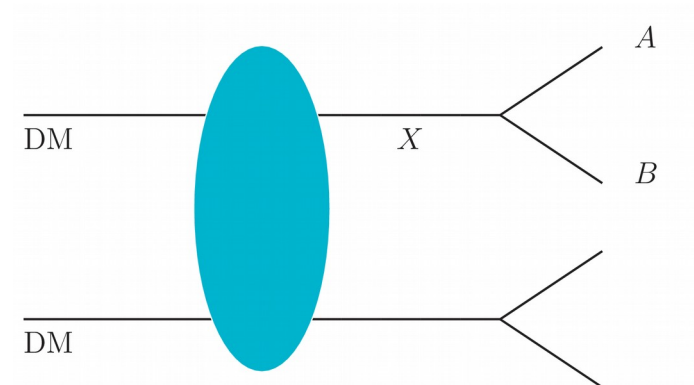
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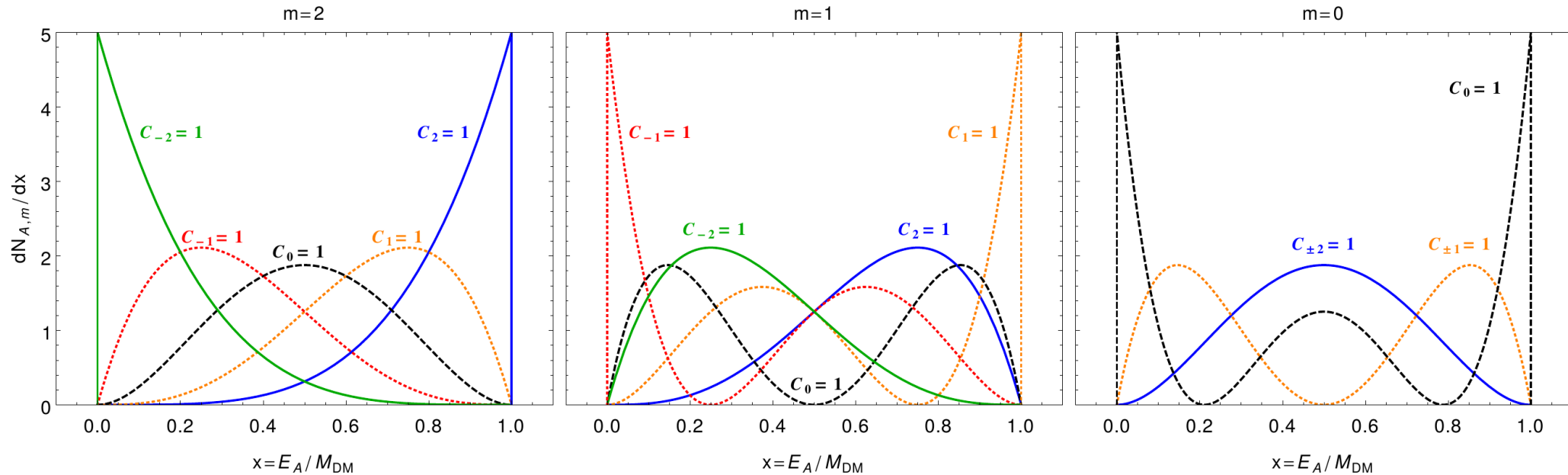
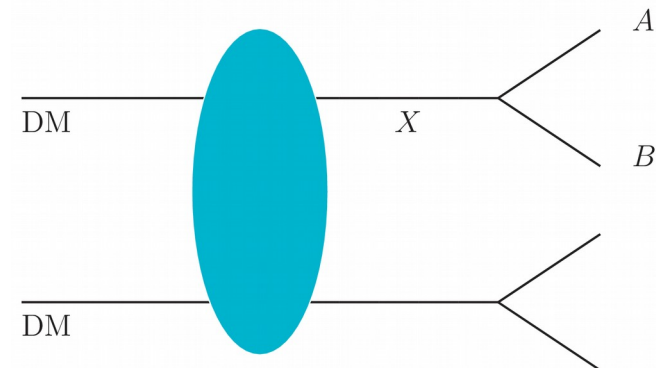
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**Fixed by the properties of the particle X and the final state**

# Example with particles of Spin-2



final state $AB$	$C_{-2}$	$C_{-1}$	$C_0$	$C_1$	$C_2$
$\gamma\gamma$	$\frac{1}{2}$	0	0	0	$\frac{1}{2}$
$ZZ$	$\frac{6}{13}$	0	$\frac{1}{13}$	0	$\frac{6}{13}$
$W^+W^-$	$\frac{6}{13}$	0	$\frac{1}{13}$	0	$\frac{6}{13}$
$hh$	0	0	1	0	0
$\nu_L\bar{\nu}_L$	0	1	0	0	0
$\nu_R\bar{\nu}_R$	0	0	0	1	0
$\nu\bar{\nu}$ (Dirac or Majorana)	0	$\frac{1}{2}$	0	$\frac{1}{2}$	0

$$f_m^S(x, y) = \frac{(2S+1)}{\sqrt{y^2 - r_X^2}} \Theta(x - x^-(y)) \Theta(x^+(y) - x) \times \sum_{m'} \left| C_{m'} \right| d_{m'm}^S \left( \arccos \left( \frac{2x - y}{\sqrt{y^2 - r_X^2}} \right) \right)^2$$

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For spin-2 particles coupled to the energy-momentum tensor

Are spin-2 particles arising in DM annihilations polarized ?

Are they coupled to the  
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# Are spin-2 particles arising in DM annihilations polarized ?

**Boosted regime**  $M_T^2 \ll p^2$

$$\varepsilon^{\mu\nu}(\pm 2) = \varepsilon^\mu(\pm) \varepsilon^\nu(\pm),$$

$$\varepsilon^{\mu\nu}(\pm 1) \simeq \frac{1}{\sqrt{2}M_T} [p^\nu \varepsilon^\mu(\pm) + p^\mu \varepsilon^\nu(\pm)]$$

$$\varepsilon^{\mu\nu}(0) \simeq \frac{\eta^{\mu\nu}}{\sqrt{6}} + \sqrt{\frac{2}{3}} \frac{p^\mu p^\nu}{M_T^2}.$$

**Are they coupled to the  
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**No**



**The spin-2 particles are  
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$$\mathbf{Br}_{0,0} = 1$$

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States with  $m = \pm 1$   
naturally decouple.  
Is there a selection rule  
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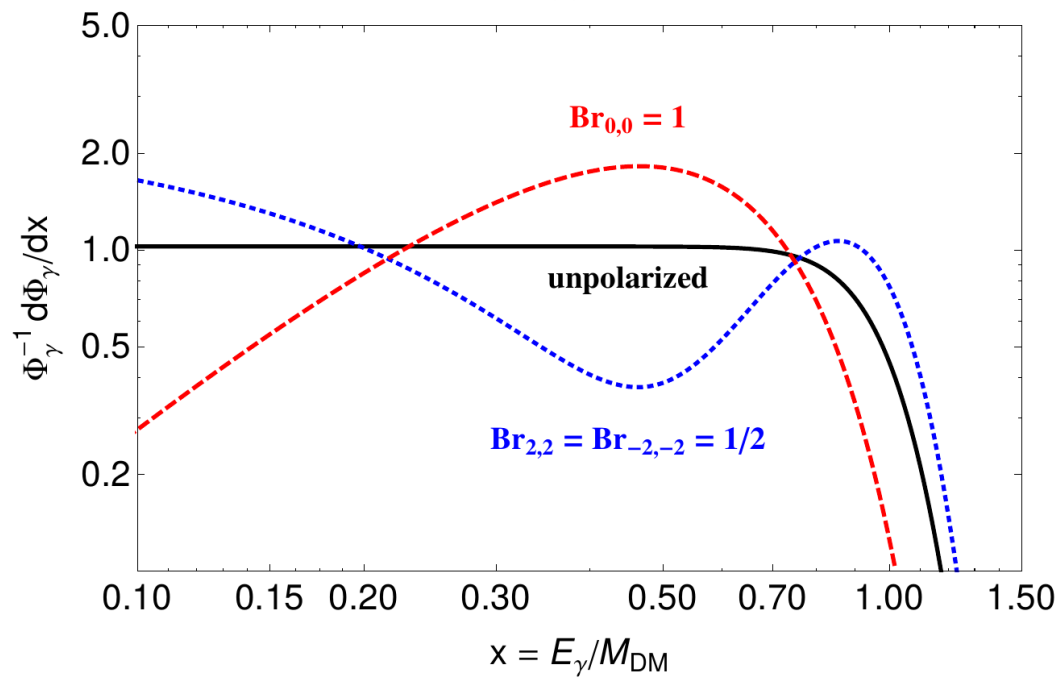
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The spin-2 particles are  
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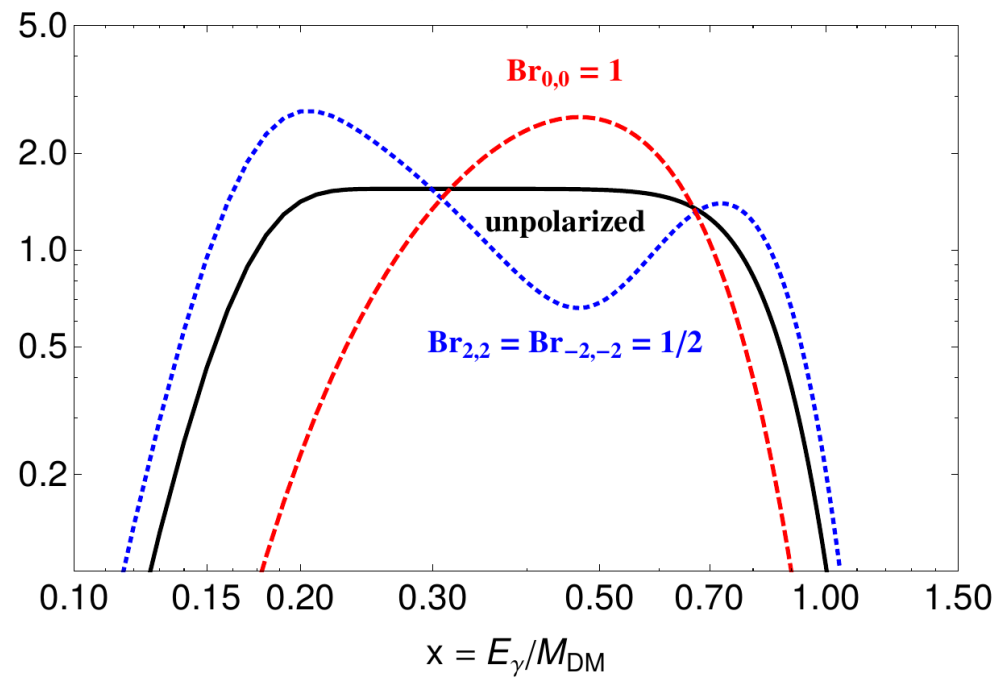
$$\text{Br}_{0,0} = 1$$

$$\text{Br}_{2,2} = \text{Br}_{-2,-2} = 1/2$$

DM DM  $\rightarrow$  TT, T  $\rightarrow \gamma\gamma$  with  $M_T/M_{\text{DM}} = 0.01$

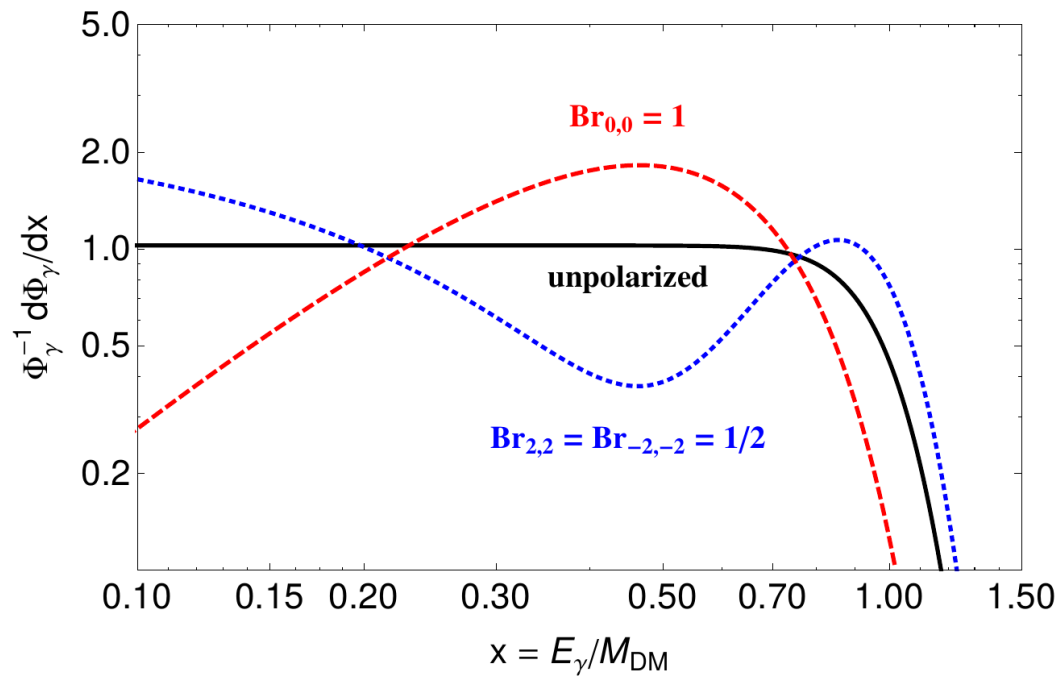


DM DM  $\rightarrow$  TT, T  $\rightarrow \gamma\gamma$  with  $M_T/M_{\text{DM}} = 0.75$

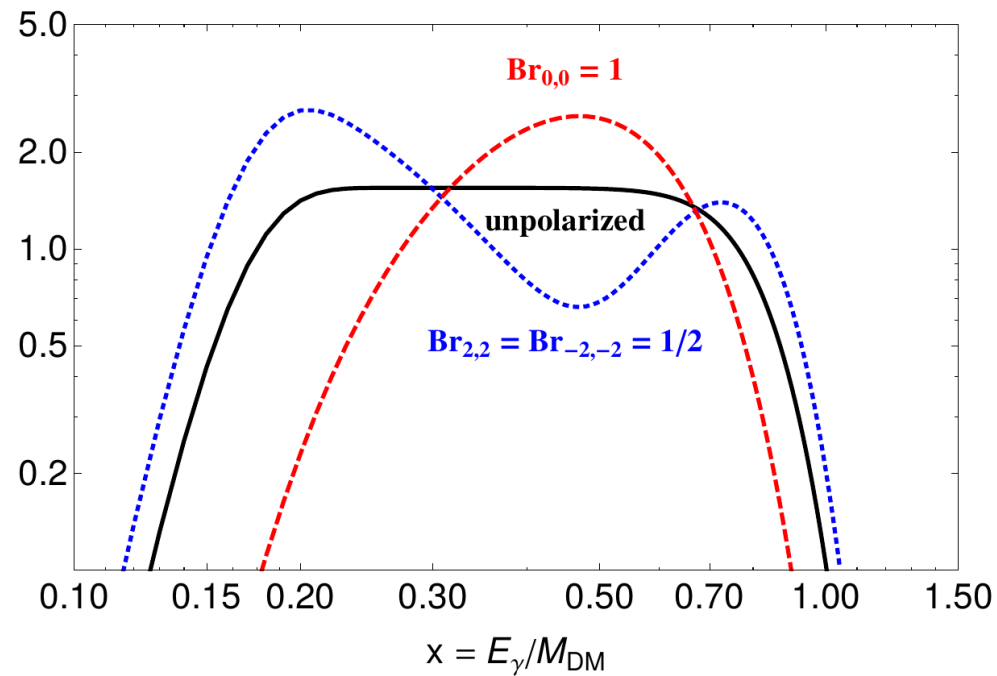




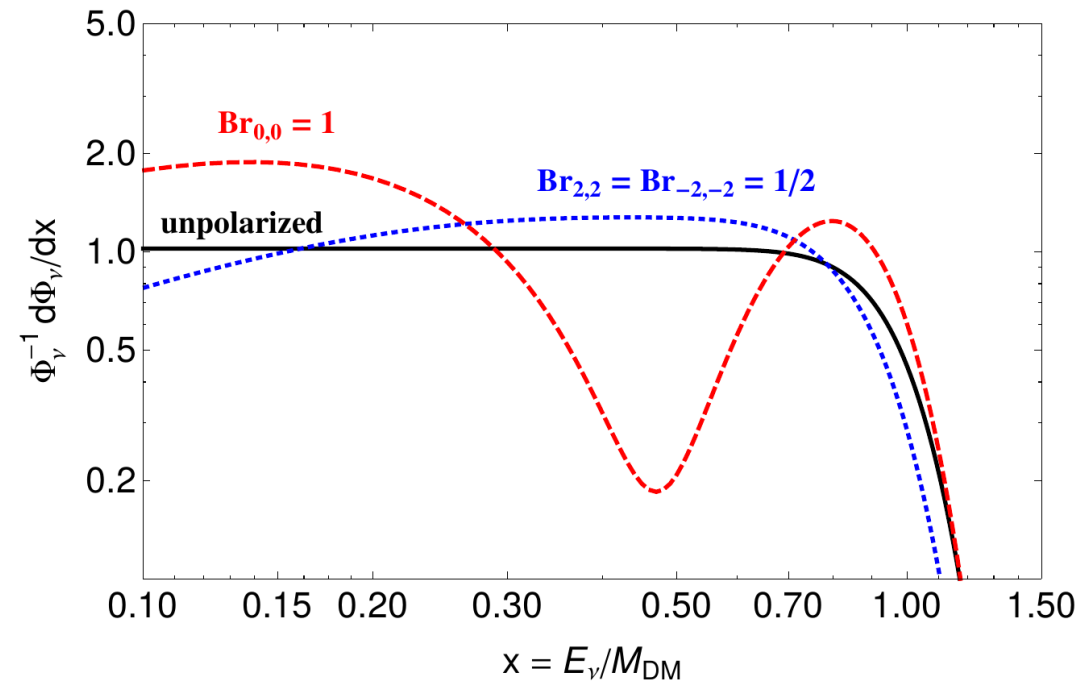
DM DM  $\rightarrow$  TT, T  $\rightarrow \gamma\gamma$  with  $M_T/M_{\text{DM}} = 0.01$



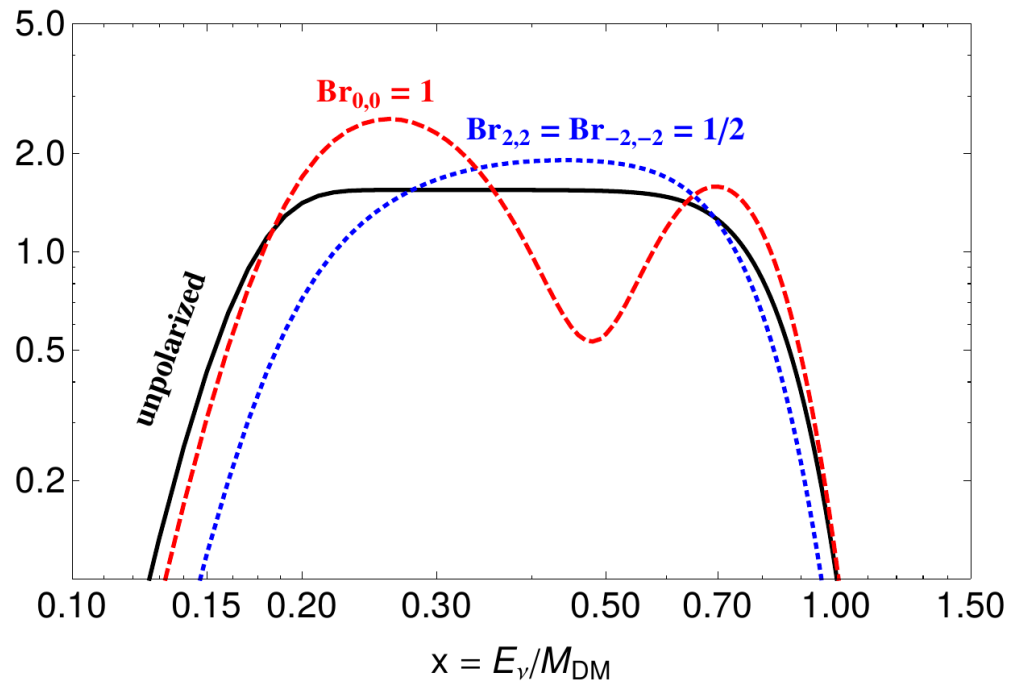
DM DM  $\rightarrow$  TT, T  $\rightarrow \gamma\gamma$  with  $M_T/M_{\text{DM}} = 0.75$



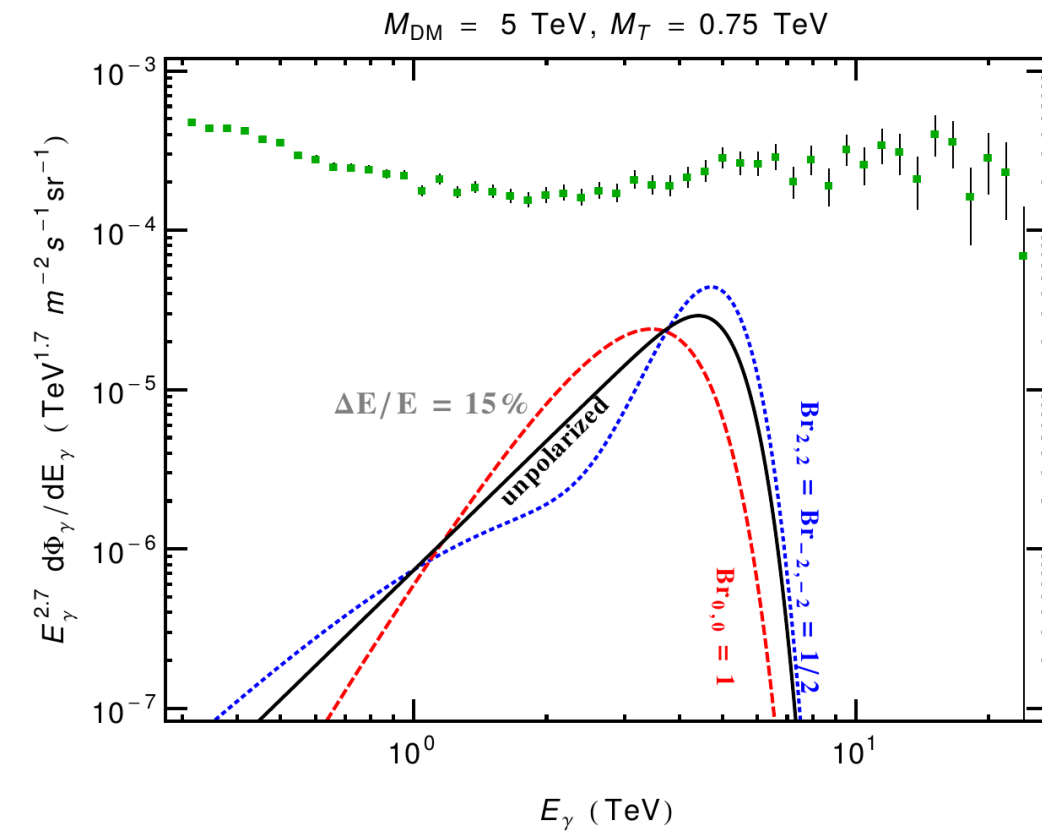
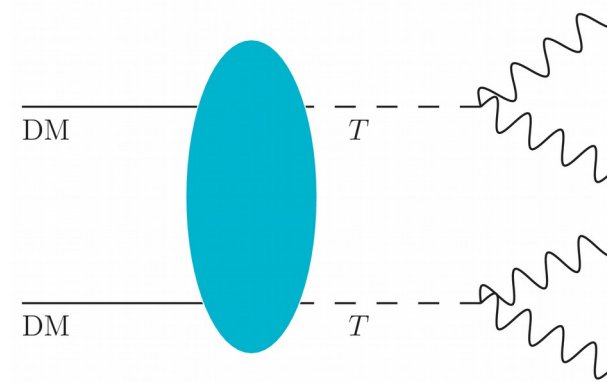
DM DM  $\rightarrow$  TT, T  $\rightarrow \nu\bar{\nu}$  with  $M_T/M_{\text{DM}} = 0.01$



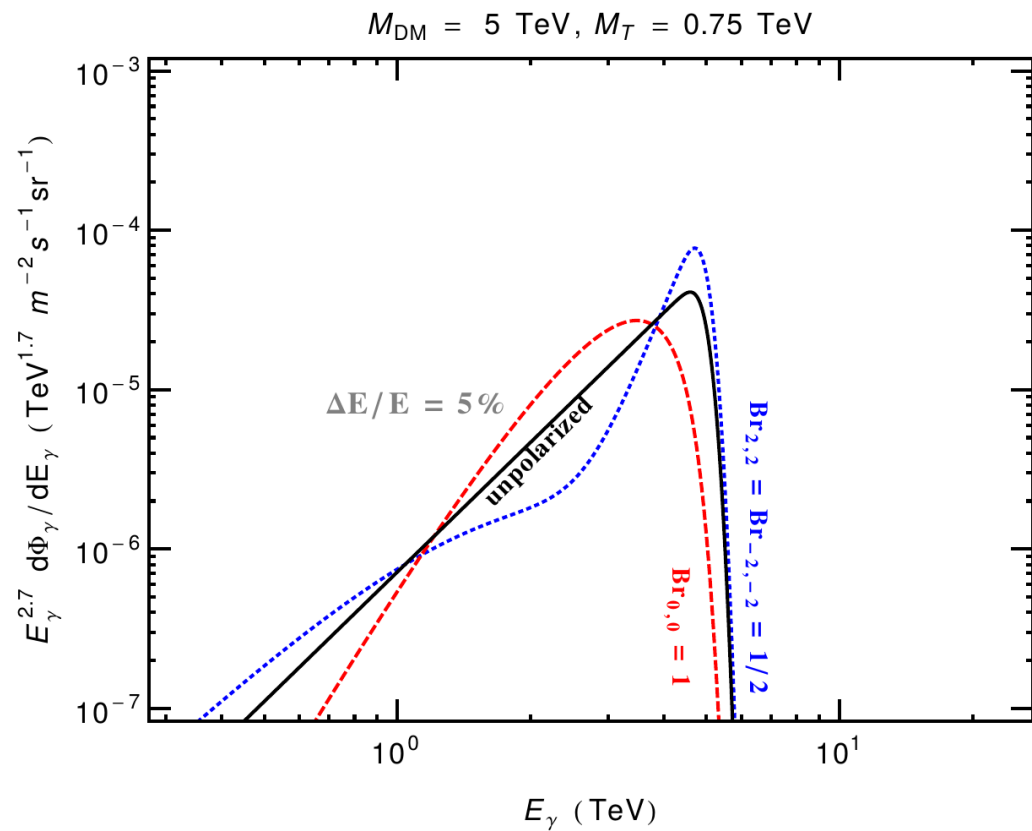
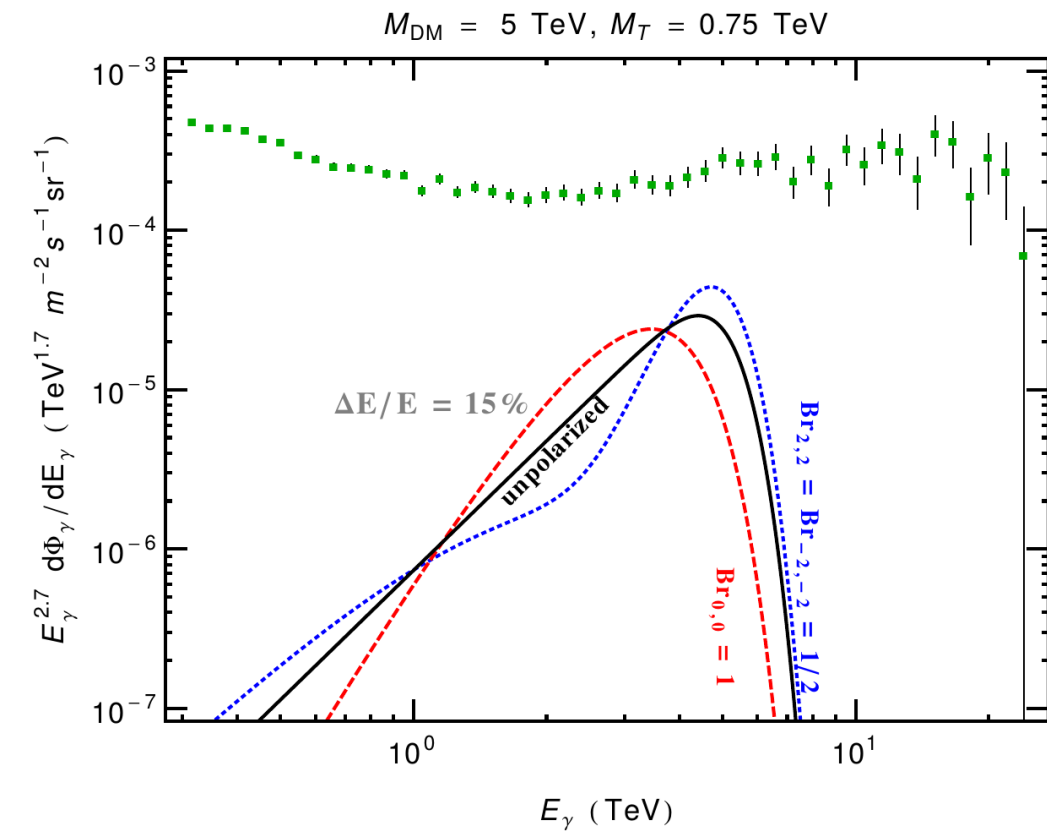
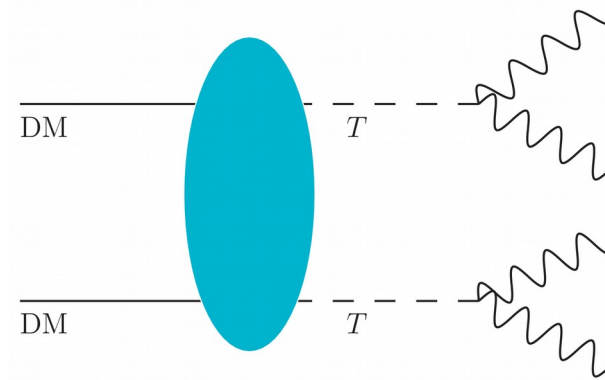
DM DM  $\rightarrow$  TT, T  $\rightarrow \nu\bar{\nu}$  with  $M_T/M_{\text{DM}} = 0.75$



# Connection to the diphoton resonance

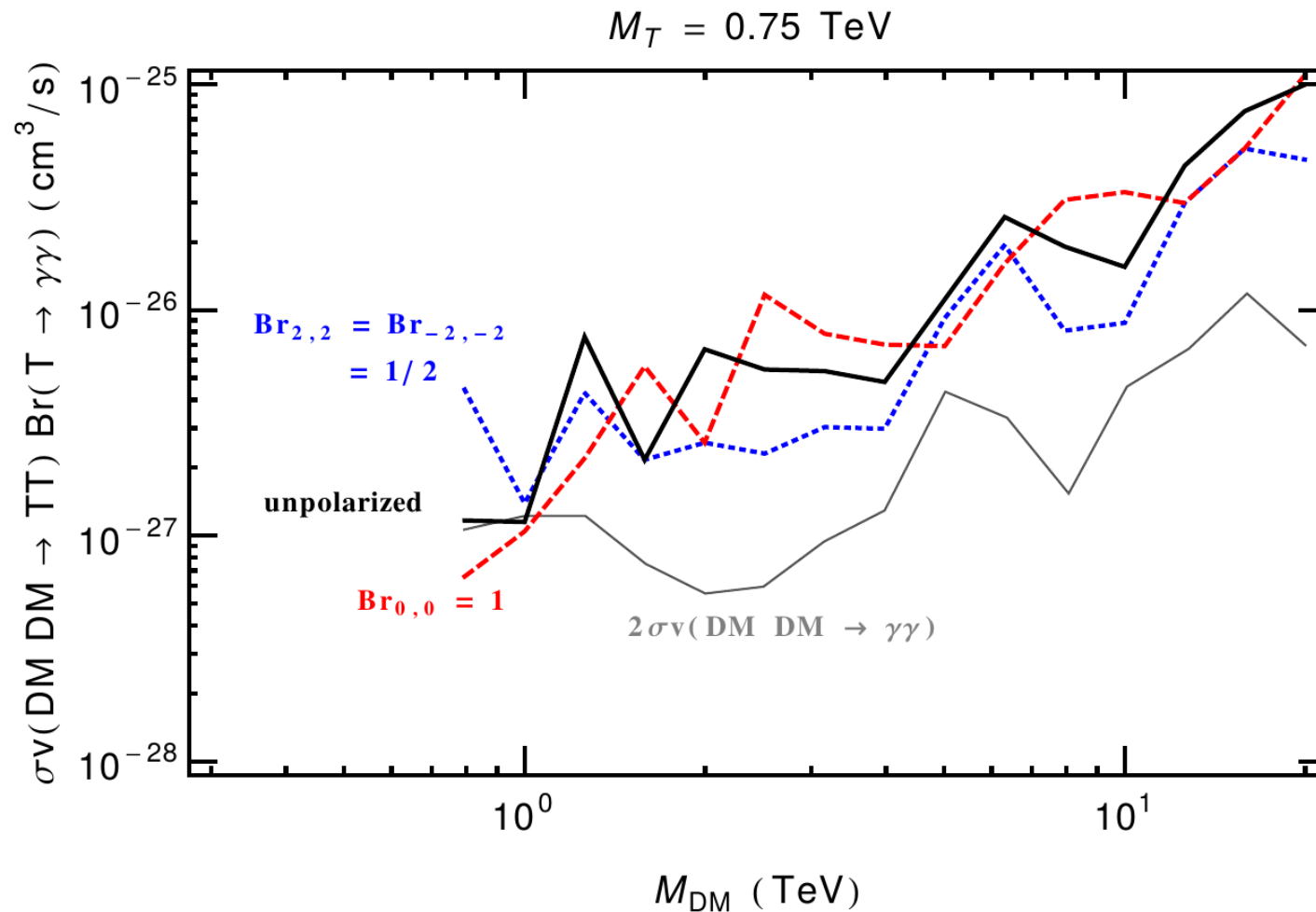


# Connection to the diphoton resonance



# Connection to the diphoton resonance

H.E.S.S. Limits on spectral features



# Conclusions

- DM annihilations into arbitrary particles that subsequently decay into photons or neutrinos lead to **polynomial spectral features**.
- Such features are generic and can be studied using a **model-independent approach**.
- Using this, high resolution of gamma-ray or neutrinos telescopes could **tell the spin of the decaying particle**.
- We calculate the annihilation spectrum that the associated to **750 GeV** resonance if DM annihilates or decays into it.