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will follow recent review S. Pakvasa and JV "Neutrino properties before and after KamLAND" hep-ph/0301061

Atmospheric Neutrinos

are produced in decay cascades initiated by collisions of cosmic rays (p, He, ...) with the Earth's atmosphere



The ν_e flux measured by underground experiments is in agreement with the predictions.

However, these experiments observe a strong deficit of ν_{μ} 's, especially of those coming from "below"



deficit is very well explained by the $u_{\mu}
ightarrow
u_{ au}$ oscillation hypothesis

Atmospheric zenith distribution

Maltoni, Schwetz, Tortola and Valle Phys. Rev. D 67 (2003) 013011 [hep-ph/0207227]



atmospheric neutrinos

1289 vs 1489-day samples

Maltoni et al Phys. Rev. D 67 (2003) 013011 hep-ph/0207227



 10^{-3}

 $\Delta m_{atm}^2 [eV^2]$

0.5

higher sterility rejection

20

15

11

-0.5

0

 $sgn(\theta_{atm}) sin^{2}(\theta_{atm})$

∆χ² atm SK-1489

99% C.L. (1 d.o.f -90% C.L. (1 d.o.f

 10^{-2}

Free [d_=any, d_=any]

Best [d_s=0.92, d_µ=0.04] Active [d_s=1.0, d_µ=0.0]

Mixed [d_=0.5, d_=any]

new parameters in atmospheric analysis

sterile as 4th, not 2nd

neglecting CP phases there are 6 angles in the lepton mixing matrix



sterility rejection depends on new parameters

Solar Neutrinos

are electron neutrinos produced in the core of our Sun by thermonuclear reactions, which generate the solar energy

All reactions result in the overall fusion of protons into helium: $4p \rightarrow {}^{4}\text{He} + 2e^{+} + \gamma + 2\nu_{e}$

- The Standard Solar Model relates the solar parameters (surface luminosity, age, radius, mass) to the total amount of neutrinos produced
- Since 1968 many experiments have measured the flux of electron neutrinos arriving at the Earth, and found they are much less than expected. This has been the Solar Neutrino Problem



early 2002

the SNO experiment showed that ν_e changes flavour due to $\nu_e
ightarrow
u_{\mu/ au}$ oscillation

solar-only regions Maltoni, Schwetz, Tortola and Valle, hep-ph/0207227



Reactor Neutrinos

Neutrinos are also produced in nuclear power plants

reactor neutrino experiments have a **well known & controlled** neutrino source

check of the solar neutrino oscillation hypothesis



KamLAND has just announced

its first results, providing for a confirmation of the solar neutrino oscillation hypothesis.

Implications of first KamLAND reactor results



first terrestrial neutrino experiment probing the solar neutrino anomaly

positive oscillation signal from first 145 days of data

combining with full solar neutrino data sample rules out non-LMA oscillations \Rightarrow oscillations happen inside the sun!

data narrow down allowed $\Delta m_{\rm SOL}^2$ range inside the LMA-MSW region, but have little impact on best fit point

in contrast to atmospheric, solar mixing remains non-maximal

Implications of first KamLAND results-2



solar+KamLAND regions: sterility rejection

Maltoni et al Phys. Rev. D 67 (2003) 013011



Stability of MSW plot



Constraining θ_{13}



Accelerator Neutrinos

Neutrinos are also produced in particle accelerators

Experiments which measure the flux of neutrinos coming from accelerators have the advantage that the **neutrino source is well known/controlled**



- **check** of the atmospheric neutrino oscillation hypothesis
- The K2K accelerator experiment is observing a small deficit in the flux of muon neutrino arriving at the detector, thus confirming atmospheric neutrino oscillations.

K2K agrees with atm results

hints of neutrino conversions also exist from the detection of accelerator-produced neutrinos in the K2K experiment



LSND

hints of neutrino conversions also exist from the detection of acceleratorproduced neutrinos in the LSND experiment



ATM

Peltoniemi, Valle, Nucl. Phys. B 406, 409 (1993); Peltoniemi, Tommasini and Valle, Phys. Lett. B 298 (1993) 383

Caldwell-Mohapatra PRD48 (1993) 325

http://www.to.infn.it/~giunti/neutrino/



can one fit all current nu-data with oscillations ?

sol+atm+reac+sbl/lsnd



stronger sterility rejection by solar & atm data

from SK-1496d-sol + SNO-NC: Maltoni et al Phys. Rev. D 67 (2003) 013011 [hep-ph/0207227]



4-nus do not fit LSND with sol+atm

Maltoni etal NPB643 (2002) 321; upd of PRD65 (2002) 093004

stronger rejection by solar & atm in 2+2 than 3+1



Cosmological Neutrinos

Neutrinos were copiously produced in the early universe, when our Cosmos was hot and dense

there are 336 neutrinos of the three flavours per cm³, a bit more than the number of photons of the Cosmic Microwave Background

Neutrinos could have been important in the production of the relic abundances of light elements

Their tiny mass but huge number might contribute to total mass of the universe and affect its expansion

Although too light to be the main component of the dark matter, neutrinos could have played an important role in the formation of largescale structure





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Dirac or Majorana?

back

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in gauge theories \beta\beta_{0\nu} \leftrightarrow majorana mass
Schechter and Valle, Phys. Rev. D 25 (1982) 2951
Hirsch's talk
like other L violating processes (e.g. nu-
transition magnetic moments) \beta\beta_{0\nu} is sensi-
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tive to Majorana phases Wolfenstein PLB107 (1981) 77; Doi et al; Bilenky et al

no such theorem for flavor violation!





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Supernova neutrinos

Supernovae are massive stars that end their lives in an extremely violent and luminous explosion in which the optical luminosity of the star at maximum can be as great as that of a small galaxy

99% of the total gravitational binding energy is emitted in the form of neutrinos of all flavours. This huge flux can transverse large distances and be detected in underground experiments before the corresponding optical signal, thus providing an early warning for astronomers



neutrinos and SN1987A

In 1987, a few neutrinos were detected from the nearby supernova 1987A that exploded in the Large Magellanic Cloud, a companion to our galaxy about 170,000 light-years away

large angle oscillations may strongly affect $\bar{\nu}_e$ SN-signal Smirnov, Spergel, Bahcall 94; Raffelt et al 96, Kachelriess et al JHEP 0101 (2001) 030, Lunardini & Smirnov

 $E_{\overline{
u}_e}$ =14 MeV, $E_{
m bind} = 3 \times 10^3 \text{ erg}$ $\tau \equiv T_{\nu_h}/T_{\overline{
u}_e}$ =1.4

pre-KamLAND

solar+SN1987A analysis

Kachelriess et al PRD65 (2002) 073016

LMA-MSW may remain best



neutrinos as future Supernova probe

The measurement of a large number of neutrinos from a future galactic supernova will give us important information both on neutrino properties and on the processes that lead to the stellar explosion

assume 10 kpc galactic SN, simulate data with given astro param

see also Barger, Marfatia & Wood



Minakata et al, PLB542 (2002) 239

improved determination of supernova parameters

High-energy astrophysical neutrinos

Neutrino astronomy can provide new information complementary to the more traditional detection of photons and cosmic rays.

Neutrinos can escape from very dense astrophysical sources and travel long distances. Therefore, they are unique candidates for the study the very high-energy Universe.

Today there are several projects of neutrino telescopes, under water or ice, which aim at detecting the Cerenkov light produced by the charged leptons originated from the interaction of upgoing high-energy neutrinos with the surrounding matter.



alternatives to oscillations?

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Oscillation vs Spin Flavor Precession

Spin Flavor Precession

Schechter, Valle PRD24 (1981) 1883 & D25, 283 Akhmedov PLB213 (1988) 64 Lim-Marciano PRD37 (1988) 1368

MHD fixes B-profile Miranda etal NPB595 (2001) 360, PLB521 (2001) 299

Non-standard interactions

consider FC/NU sub-weak strength terms εG_F

can induce oscillations of massless neutrinos in matter, which are E-independent, converting both neutrinos & anti-nu's, can be resonant in SNovae Valle PLB199 (1987) 432, Roulet 91; Guzzo etal 91; Barger etal 91

global analysis in Guzzo et al NPB629 (2002) 479 MORE





Oscillation vs Spin Flavor Precession

Barranco et al PRD66 (2002) 093009 [hep-ph/0207326]

current solar data still do not allow the reconstruction of the profile of ν_e -conversion probability

LMA-MSW, RSFP, NRSFP equivalent



KamLAND lifts degeneracy

ruling out SFP as main solution

testing SFP as sub-leading

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probing neutrino magnetic moments at LMA-MSW



a

tan²0

0.3

tan²θ

H

D

How robust are atmospheric oscillations?

very good contained atm-fit, Gonzalez-Garcia etal, Phys. Rev. Lett. 82 (1999) 3202 [hep-ph/9809531]



probing NSI with atmospheric data

Fornengo et al, Phys. Rev. D 65 (2002) 013010 [hep-ph/0108043].



atm bounds on **FC** and **NU** nu-interactions

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Neutrino Factories

apart from probing s_{13} and δ ... Cervera et al, De Rujula, Gavela, Hernandez, Freund, Huber, Lindner, Albright et al, Barger et al...

they can probe NSI



Improved FC-tests at NuFact back

Huber & JV Phys. Lett. B 523 (2001) 151 [hep-ph/0108193]

10 kt detector, .33 ν_{τ} detection eff above 4 GeV; need no tau charge id



FCI-oscillation confusion theorem back a neutrino factory is less sensitive to θ_{13} 10^{-1} because non-standard neutrino interactions are confused with oscillations Huber, Schwetz & JV Phys. Rev. Lett. 88 (2002) d 101804 [hep-ph/0111224] b $J_{\Psi}^{5} 10^{-2}$ 10^{-3} near-site programme essential 10^{-4} 10^{-3} 10^{-2} 10^{-1} $\sin^2 2\theta_{13}$

 2×10^{20} mu/yr/polarity \times 5 yr, 40 kt magn iron calorim, 10% muon E-resoln above 4 GeV

FCI-oscillation confusion theorem-2 back

Huber, Schwetz and J. V. Phys. Rev. D 66, 013006 (2002) [hep-ph/0202048]



 2×10^{20} mu/yr/polarity \times 5 yr, 40 kt magn iron calorim, 10% muon E-resoln above 4 GeV

90% CL reach on $\sin^2 2\theta_{13}$ vs NSI bounds

The dotted line is for $700~{\rm km}$, dash-dotted for $3\,000~{\rm km}$ and dashed is for $7\,000~{\rm km}$ baseline

horizontal black line is the current NSI limit

vertical grey band is the sensitivity without NSI

diagonal solid line is the theoretical bound derived from our confusion theorem

Neutrinos and new physics

massive neutrinos of different flavour $(\nu_e, \nu_\mu, \nu_\tau)$ can mix by quantum mechanics

Since 1998, measurements of the flux of atmospheric neutrinos in the Super-Kamiokande detector have shown evidence for neutrino $\nu_{\mu} \rightarrow \nu_{\tau}$ oscillations

In April 2002, the SNO collaboration demonstrated that the solar neutrinos were also converted into other flavours before reaching the Earth

study of neutrino properties is the best way

to probe new physics at scales





nu-oscillations first/only evidence for physics beyond the Standard Model

Theory of neutrino properties

how to reconstruct the parameters

how to reconstruct the underlying Theory

simplest gauge theory mixing matrix 12=sol 13=reac 23=atm 3 angles θ_{ij} 1 KM-like \mathcal{O} 2 Majorana phases $\beta \beta_{0\nu}$ $\phi_1\,,\phi_2$ Schechter and Valle, Phys. Rev. D 22 (1980) 2227 hierarchical splittings max θ_{23} , large θ_{12} & small θ_{13} **NORMAL INVERSE**

quasi-degenerate may lead to $\beta\beta_{0\nu}$ rate similar to present hint Ioannisian & J. V. PL B332 (1994) 93; Caldwell & Mohapatra; Joshipura; Bamert & Burgess; Balaji, Mohapatra, Parida & Paschos, Babu, Ma & Valle, ... Ellis & Lola, Ma, Casas et al, Haba et al, ... back

leptonic CP violation

will be a challenge !

"Dirac" CPV suppressed, since ϕ disappears when $\Delta_{12} \rightarrow 0$ Schechter and Valle, Phys. Rev. D **21** (1980) 309

- "Majorana" CPV absent from conventional $\Delta L = 0$ oscillations
- require L violation and effect is V-A suppressed (Dirac-Majorana confusion theorem) e.g. $\beta\beta_{0\nu}$ (Doi et al 1981) & $\Delta L = 2$ oscillations

Schechter and Valle, Phys. Rev. D 23 (1981) 1666

- must look for chirality violating processes, such as neutrino electromagnetic form factors
- or processes involving large "Majorana"-masses, such as | leptogenesis
- the seesaw connection opens the remarkable possibility that the Majorana phases of the light neutrrinos affect the baryon asymmetry generated in leptogenesis models in an unsuppressed way



basic dim-5 operator back



from Gravity

Weinberg; Barbieri, Ellis, Gaillard; Akhmedov et al

from seesaw schemes

Gell-Mann, Ramond, Slansky; Yanagida; Mohapatra, Senjanovic Phys. Rev. Lett. **44** (1980) 91 Schechter, Valle Phys. Rev. D **22** (1980) 2227

here I consider at an effective level

neutrino unification

back

- Babu, Ma and Valle, Phys. Lett. B 552 (2003) 207 [hep-ph/0206292]
- due to A_4 symmetry neutrino masses unify as they run up



Chankowski et al, Phys. Rev. Lett. 86 (2001) 3488 [hep-ph/0011150]

solar & atm splittings arise from RGE + threshold effects

common origin for neutrino and KM mixing

maximal θ_{23} ; large θ_{12} & small θ_{13}

observable neutrino mass eg in cosmology, β and $\beta\beta_{0\nu}$ decays

observable LFV
$$B(\tau \to \mu \gamma) \sim 10^{-6}$$

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The seesaw

neutrino masses follow from the exchange of heavy isosinglet neutral heavy leptons with mass $M_R = M_R^T \propto \Delta_R$ (126 of SO(10), type I) or from the exchange of heavy scalar bosons (type II)

$$\left(\begin{array}{cc} M_L & D \\ D^T & M_R \end{array}
ight)$$



the first gives $M_{\nu \text{ eff}} = M_L - DM_R^{-1}D^T$ where D is the standard $SU(2) \otimes U(1)$ breaking Dirac mass term

the M_L term is propto an effective iso-triplet vev Schechter, Valle Phys. Rev. D 22 (1980) 2227 also suppressed by the left-right breaking scale, $M_L \propto 1/M_R$ Mohapatra, Senjanovic

hardly any predictivity, unless specific symmetries are assumed

global realization of seesaw mechanism



seesaw charged and neutral currents

Schechter, Valle, Phys. Rev. D 22, 2227 (1980) & D 25, 774 (1982)

- arbitrary number of $SU(2) \otimes U(1)$ singlets implies that the mixing matrix describing the charged leptonic weak interaction is a rectangular matrix K which may be decomposed as $K = (K_L, K_H)$ where K_L and K_H are 3×3 matrices.
- far more mixing angles θ_{ij} and CP violating phases ϕ_{ij} than needed to describe the charged current weak interaction of quarks, since (i) neutrinos are Majorana particles so that their mass terms are not invariant under rephasings, and (ii) the isodoublet neutrinos mix with the isosinglets.
- The NC weak interactions are described by a non-trivial matrix $P = K^{\dagger}K$ implies $\nu_h \to 3\nu$
- strength of new couplings maybe sizable in variants of the seesaw: NSI
- The (3, 1) model has 2 massless neutrinos. The other two form a light-heavy Majorana pair or a Dirac pair if lepton number is a good symmetry. If not the massless degeneracy is lifted by radiative corrections.
- forms the basis for hybrid model where atm scale comes from tree, while solar arises from loops

minimalistic neutrino masses?

Gouvea and Valle, Phys. Lett. B 501 (2001) 115 [hep-ph/0010299]

- based on 3 + 1 scheme
- seesaw generates atm scale, leaving 2 neutrinos massless
- degeneracy lifted by gravity-induced dim-5: solar scale



• KamLAND implies need for a non-gravitational mechanism to generate the solar scale

SUSY origin for neutrino mass

back

spontaneous RPV

Aulakh, Mohapatra 83; Hall, Suzuki 84 Ross, Valle 85; Ellis et al 85; Santamaria, Valle 87 ...

singlet sneutrino vev

Masiero and Valle, Phys. Lett. B 251 (1990) 273

attractive dynamics and systematic parametrization of RPV

spontaneous RPV \rightarrow effective bilinear RPV

hybrid neutrino masses





U(1) family symmetries

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Mira, Nardi, Restrepo and Valle, Phys. Lett. B 492 (2000) [hep-ph/0007266] back

- quark and lepton mixing from textures
- U(1) symmetry \rightarrow simplest bilinear RPV SUSY model: $W = W_{MSSM} + \mu_{\alpha} \ell_{\alpha} H_u$

common origin for μ -problem & nu-anomalies

• $\mu_0 \sim m_{3/2} \theta$ Giudice-Masiero $\mu_i \sim m_{3/2} \theta^{7+x}$ Nilles-Polonsky

- RPV-seesaw gives atm scale, leaving 2 nu's massless
- degeneracy lifted by loops, to give solar scale. Need to change U(1) assignments so as to get hugher Δm_{SOL}^2 in agreement with Kamland



bilinear RPV solution to neutrino anomalies

back Diaz, Hirsch, Porod, Romao and Valle, hep-ph/0302021; Phys. Rev. D 62 (2000) 113008 [Err-ibid. D 65 (2002) 119901]; Phys. Rev. D 61 (2000) 071703



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LSP decay length [cm]: BRPV

back from Bartl et al NPB 600 (2001) 39



Mukhopadhyaya, Roy & Vissani; Chun & Lee; Choi et al; Datta et al

no neutralino dark matter

any charged SUSY particles can be the LSP



light sterile-nus from extra dimensions

back

Ioannisian, JV PRD63 (2001) 073002

Antoniadis, Arkani-Hamed, Dimopoulos, Dvali... Mohapatra, Perez-Lorenzana...

sterile-nu as zero-th mode of the Kaluza-Klein tower



light-nu's without new scale

tree

Gonzalez-Garcia, JV, Phys. Lett. B 216, 360 (1989); Romao and JV, Nucl. Phys. B 381 (1992) 87...

vs radiative

Zee or Babu-type

In contrast to seesaw

m-nu $\rightarrow 0$ as the LNV scale \rightarrow zero

Higgs \rightarrow 2-majoron "invisible" decays

Joshipura, JV NPB397 (1993) 105; Campos et al PRD55 (1997) 1316 back



predicting nu-mass and mixing?

back

- top-bottom vs bottom-up
- hierarchical vs quasi-degenerate, sterile-nus?
- what is the scale ?
 - Planck scale: Strings?
 - GUT scale E(6), SO(10),...
 - Intermediate scale: P-Q, L-R ...
 - Weak $SU(3)\otimes SU(2)\otimes U(1)$ scale
- what is the mechanism?
 - tree vs radiative
 - B-L gauged vs ungauged...
 - no theory of flavour

In short

- oscillations fit well sol+atm, but **not LSND**
- LMA-MSW as astro-probe: SuperNovae and Sun probing the Sun beyond helioseismology
- non-standard properties can only play a sub-leading role in solar and atm robustness
- NSI test @ NuFact e-tau NSI-OSC confusion...
- no hard theory predictions

but suggests Majorana

• if neutrino masses arise from low energy supersymmetry, neutrino properties may be testable at high energy accelerators