

Relic density of neutrinos with primordial asymmetries

SP, T.Pinto & G.Raffelt,
PRL 102 (2009) 241302 [arXiv:0808.3137]

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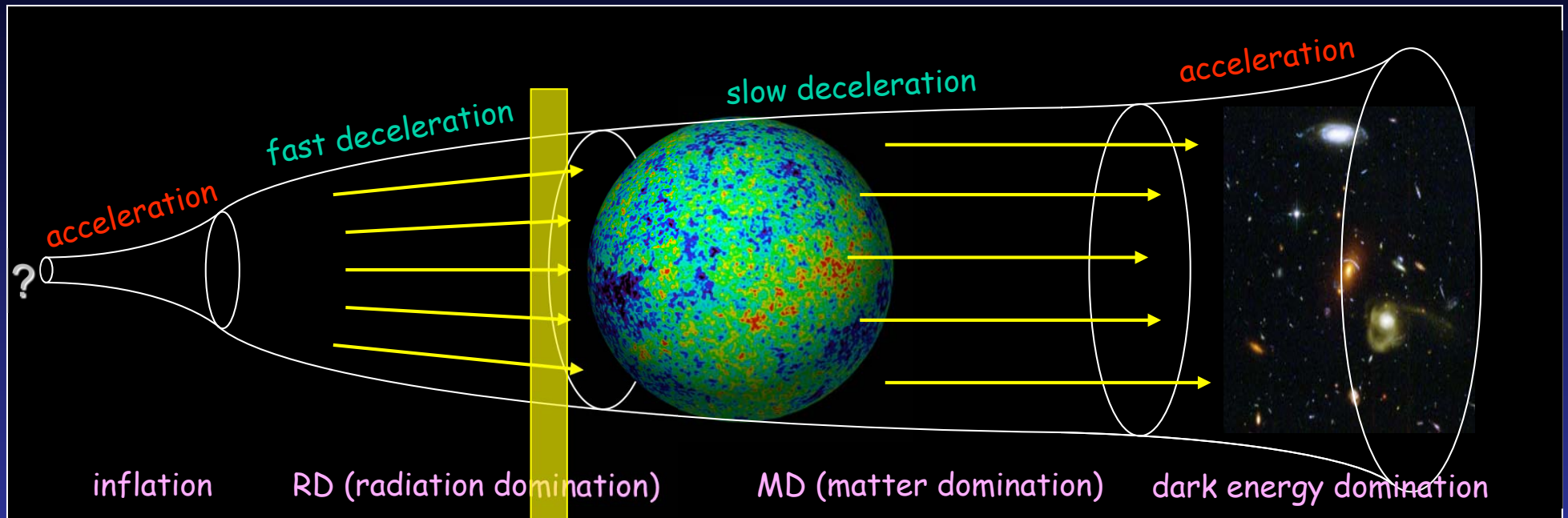
Outline

Introduction: the radiation component of the Universe

Relic neutrino asymmetries and flavour oscillations

Relic density of neutrinos with primordial asymmetries

Evolution of the Universe



a_{eq} : Equality $\rho_r = \rho_m$

Relativistic energy density

At $T \gg m_e$, the radiation content of the Universe is

$$\rho_r = \rho_\gamma + \rho_\nu = \frac{\pi^2}{15} T^4 + 3 \times \frac{7}{8} \times \frac{\pi^2}{15} T^4 = \left[1 + \frac{7}{8} \times 3 \right] \rho_\gamma$$

At $T < m_e$, the radiation content of the Universe is

$$\rho_r = \rho_\gamma + \rho_\nu = \frac{\pi^2}{15} T_\gamma^4 + 3 \times \frac{7}{8} \times \frac{\pi^2}{15} T_\nu^4 = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} 3 \right] \rho_\gamma$$

$$\rho_r = \rho_\gamma + \rho_\nu = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right] \rho_\gamma$$

$\frac{T_\nu^4}{T_\gamma^4}$

Relativistic energy density

At $T < m_e$, the radiation content of the Universe is

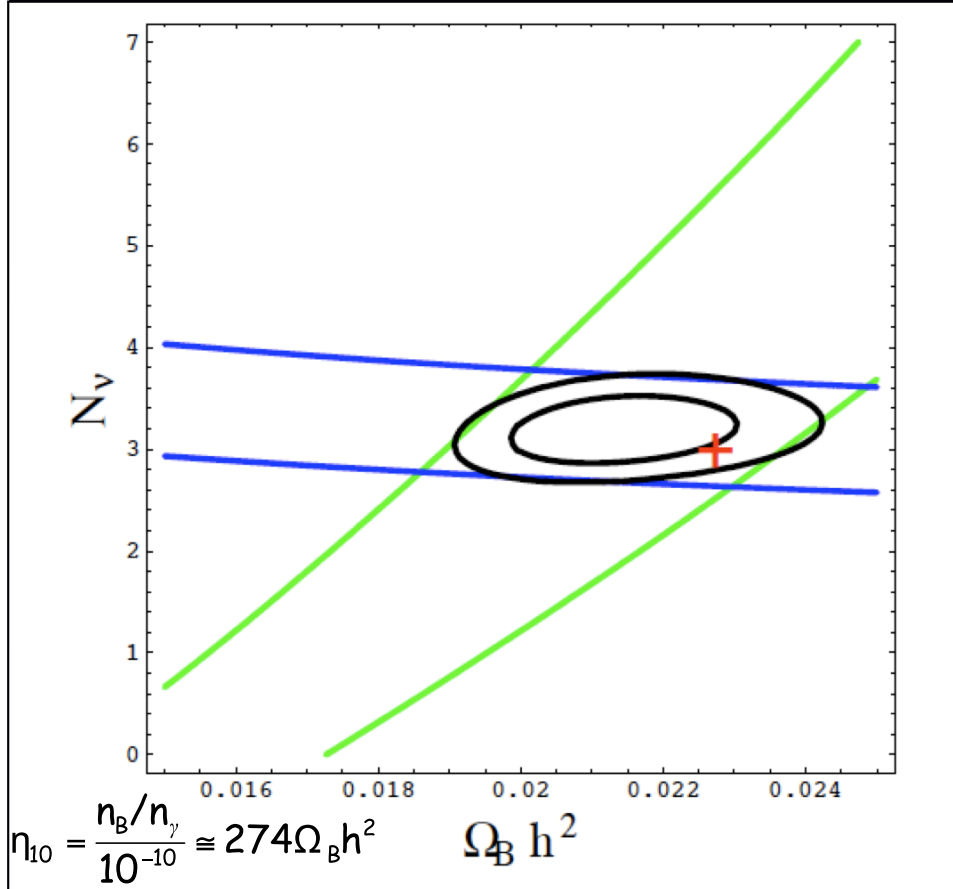
$$\rho_r = \rho_\gamma + \rho_\nu + \rho_x = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right] \rho_\gamma$$

Effective number of relativistic neutrino species

Traditional parametrization of the energy density stored in relativistic particles

of flavour neutrinos: $N_\nu = 2.984 \pm 0.008$ (LEP data)

allowed ranges for N_{eff}

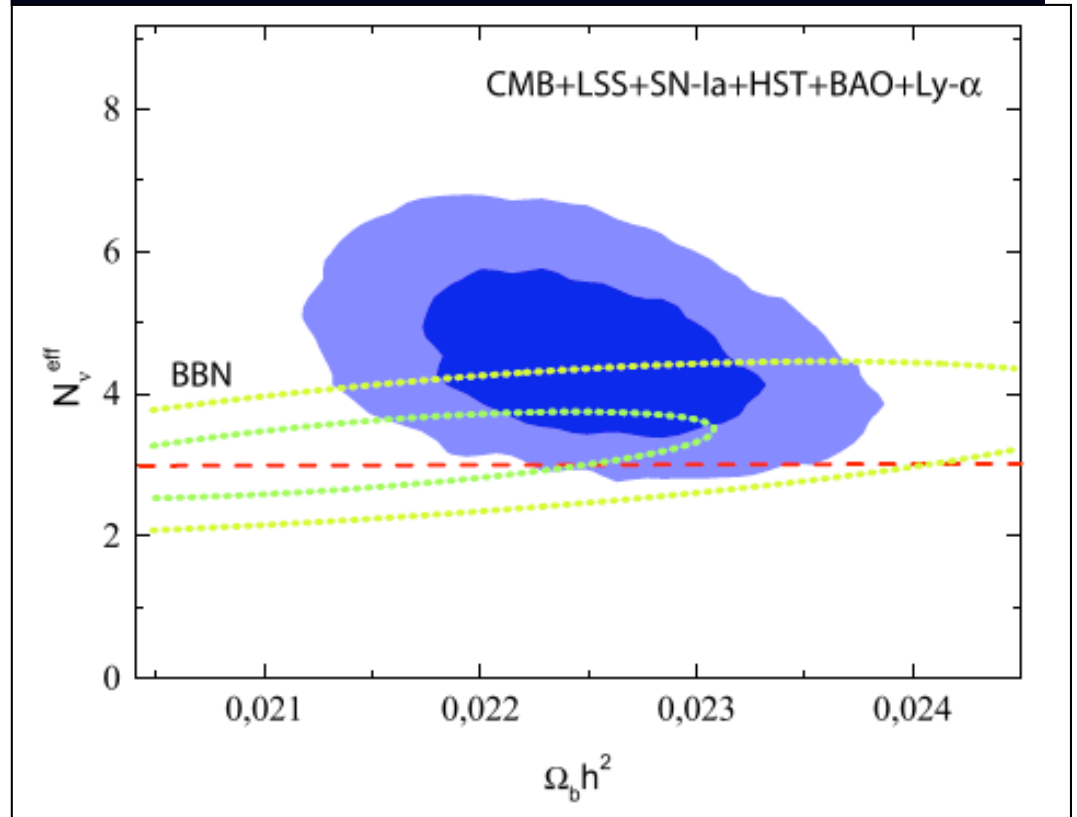


Bounds from BBN

F. Iocco et al, Phys. Rep. 472 (2009) 1

$$N_{\text{eff}} = 3.0 \pm 0.3_{\text{stat}} \pm 0.3_{\text{syst}}$$

(95% CL)



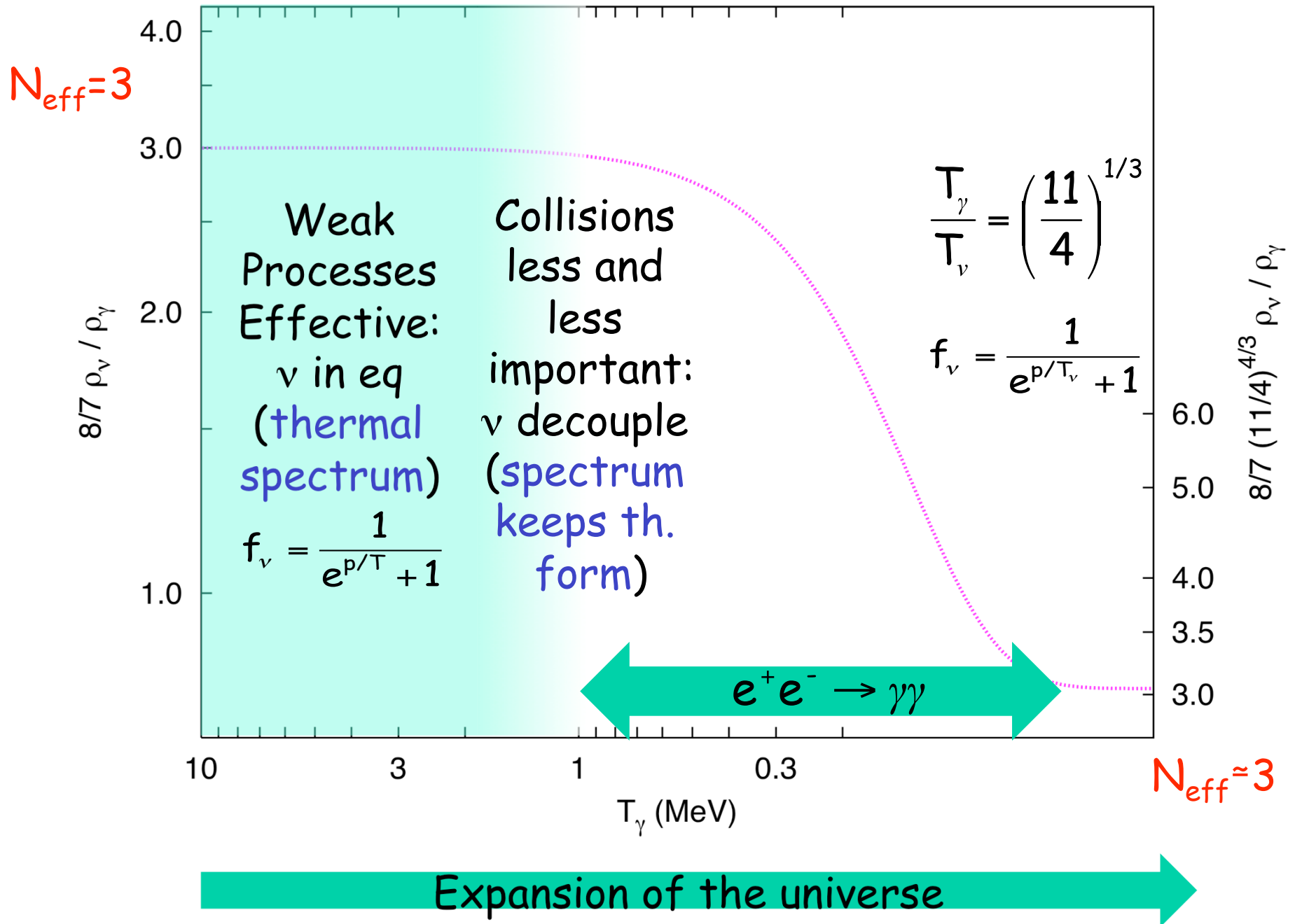
Bounds from non-BBN data

Mangano et al, JCAP 0703 (2007) 006

$$3.0 < N_{\text{eff}} < 7.9 \quad (\text{CMB} + \text{LSS data})$$

$$3.1 < N_{\text{eff}} < 6.2 \quad (+\text{BAO and Ly} - \alpha)$$

Evolution of ρ_ν/ρ_γ before and after ν decoupling

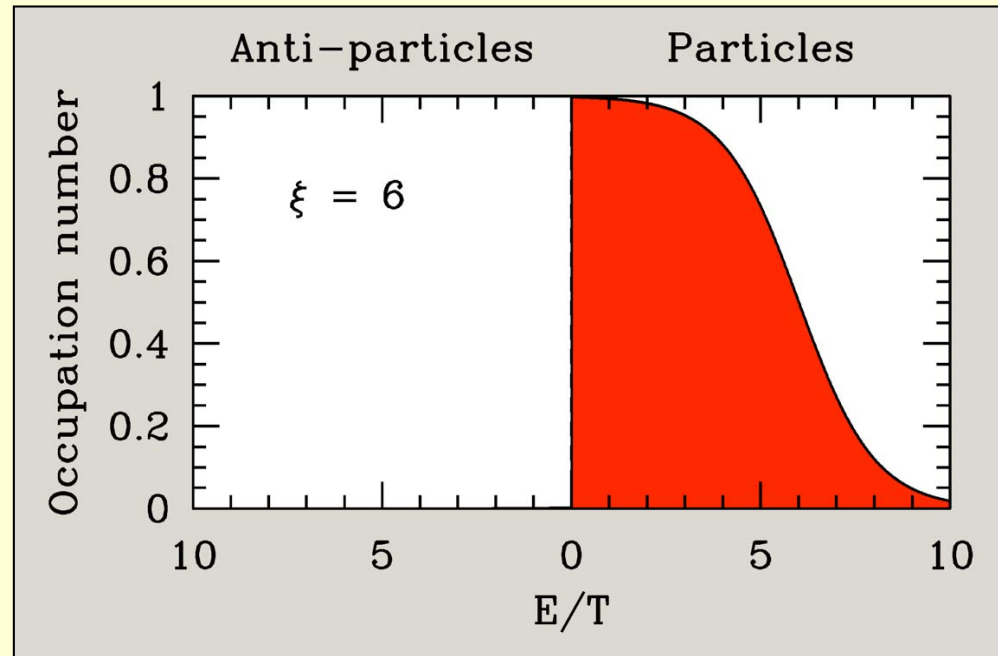


Primordial Neutrino asymmetries

Fermi-Dirac distribution

- Temperature T
- Chemical potential μ
 - + μ Particles
 - μ Anti-particles

$$f_p = \frac{1}{\exp\left(\frac{E - \mu}{T}\right) + 1}$$



Degeneracy parameter

$$\xi = \frac{\mu}{T}$$

Invariant under cosmic expansion

Number density

$$n_{\nu\bar{\nu}} = \int dE \frac{4\pi}{(2\pi)^3} \left(\frac{E^2}{1 + \exp(E/T - \xi)} + \frac{E^2}{1 + \exp(E/T + \xi)} \right)$$

$$= \frac{3\zeta_3}{2\pi^2} T_v^3 \left[1 + \frac{2\ln(2)}{3\zeta_3} \xi^2 + \frac{1}{72\zeta_3} \xi^4 + \dots \right]$$

BBN and Neutrino Chemical Potentials

Expansion Rate
Effect
(all flavors)

Energy density in one neutrino flavor with degeneracy parameter $\xi = \mu/T$

$$\rho_{\nu\bar{\nu}} = \frac{7\pi^2}{120} T_\nu^4 \left[1 + \underbrace{\frac{30}{7} \left(\frac{\xi}{\pi}\right)^2 + \frac{15}{7} \left(\frac{\xi}{\pi}\right)^4}_{\Delta N_{\text{eff}}} \right]$$

Beta equilibrium
effect for
electron flavor



Helium abundance essentially fixed by n/p ratio at beta freeze-out

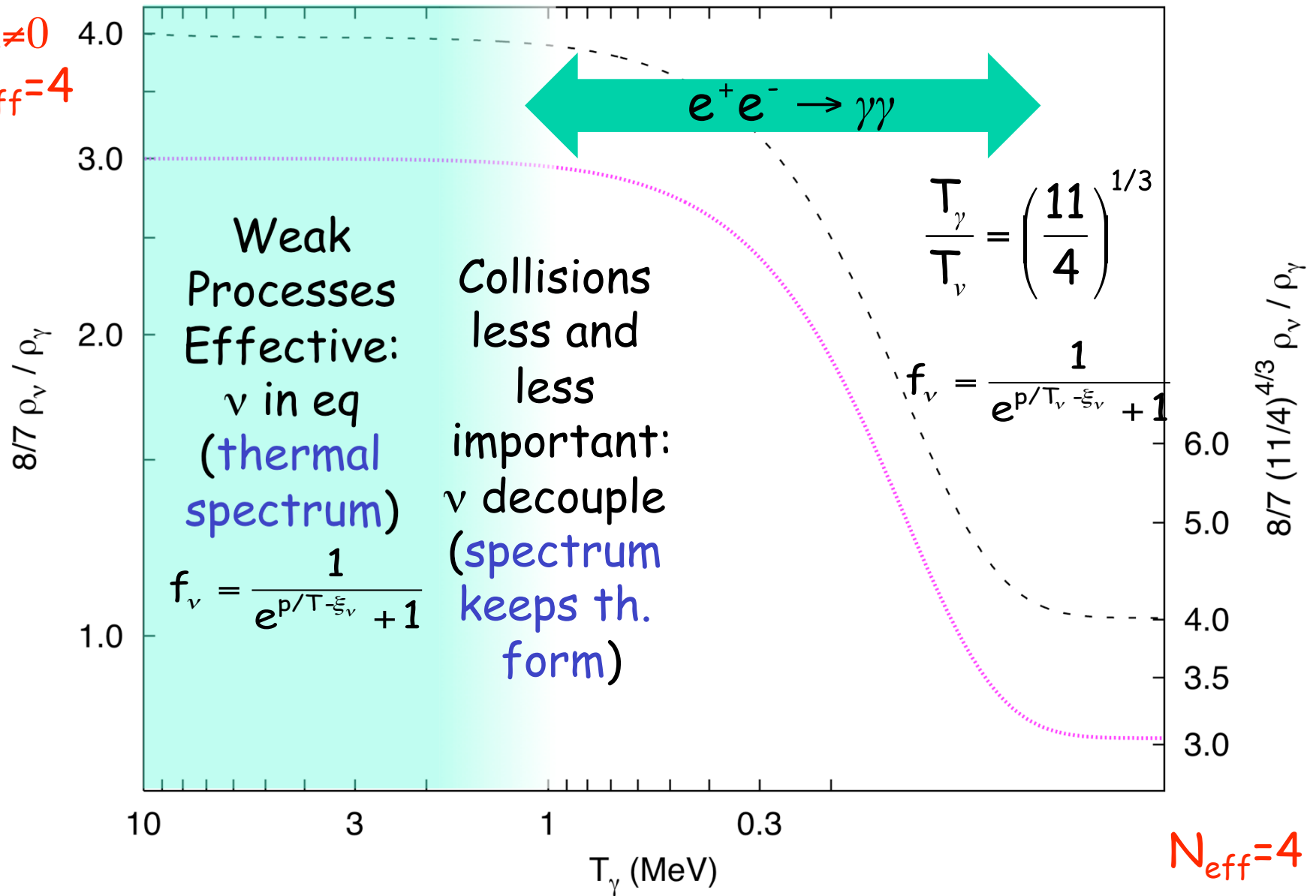
$$\frac{n}{p} = e^{-(m_n - m_p)/T - \xi_{\nu_e}} \quad |\xi_{\nu_e}| \lesssim 0.07$$

Effect on ${}^4\text{He}$ equivalent to $\Delta N_{\text{eff}} \sim -18 \xi_{\nu_e}$

- ν_e beta effect can compensate expansion-rate effect of $\nu_{\mu,\tau}$
- No significant BBN limit on neutrino number density

Evolution of ρ_ν/ρ_γ with asymmetries ($L_\nu^{\text{tot}}=0$)

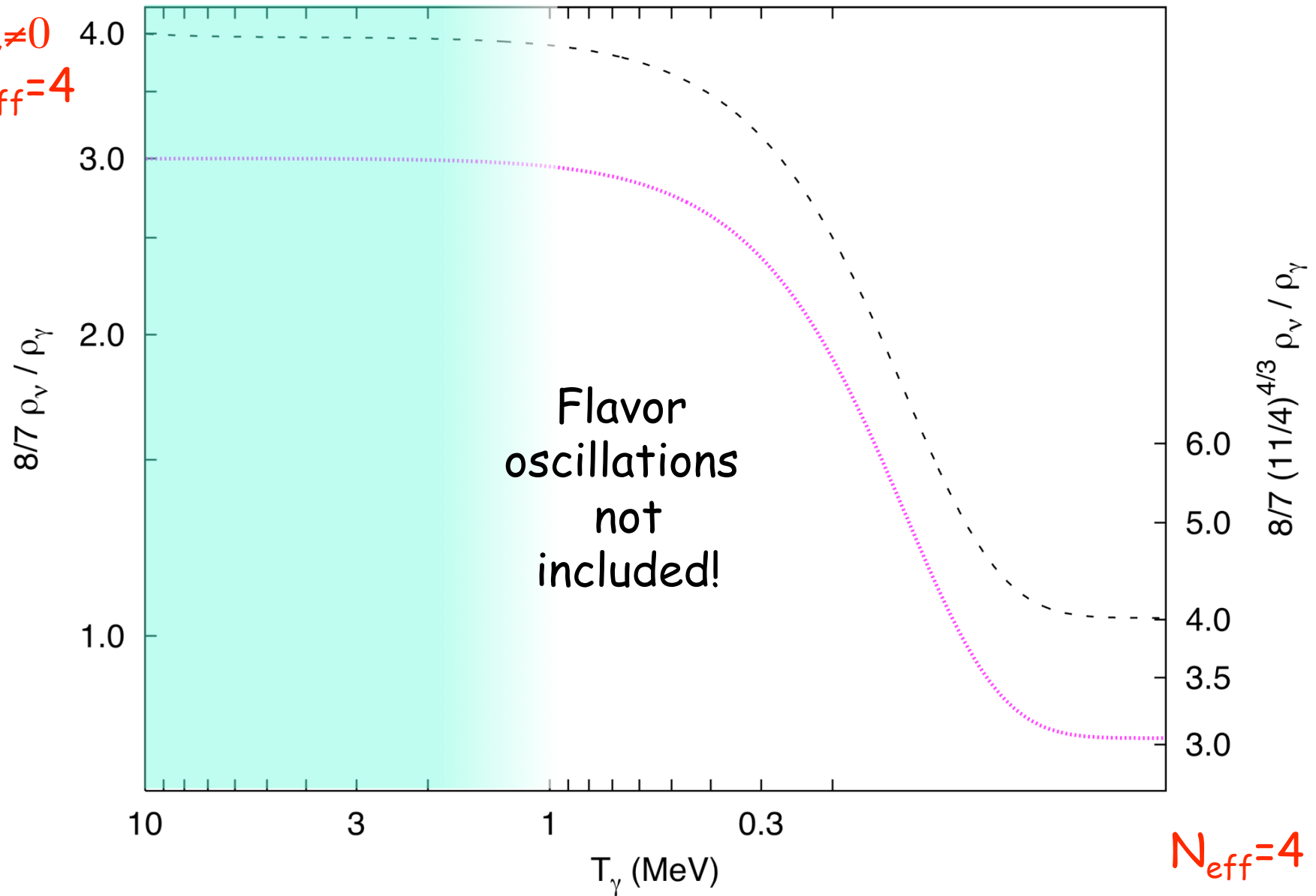
$\xi_\nu^0 \neq 0$
 $N_{\text{eff}}^0 = 4$



Expansion of the universe

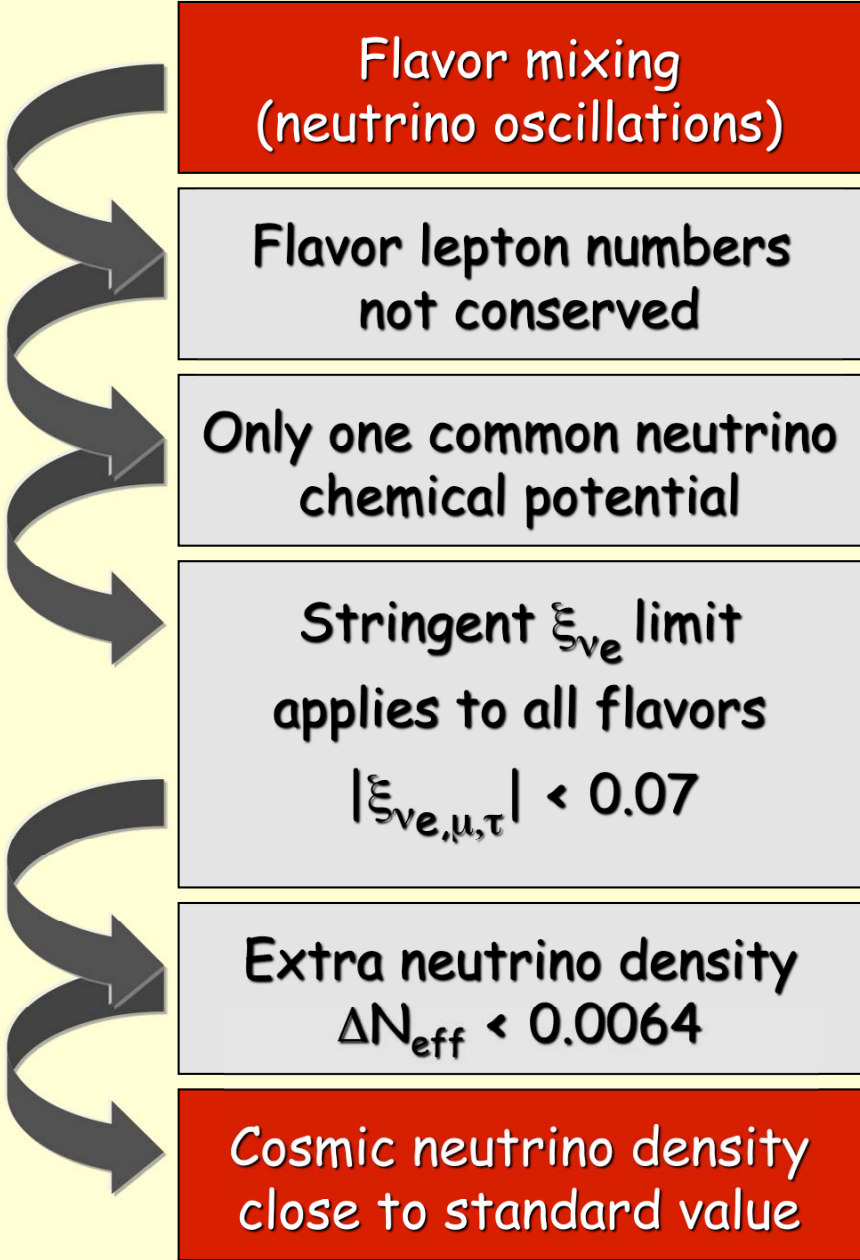
Evolution of ρ_ν/ρ_γ with asymmetries ($L_\nu^{\text{tot}}=0$)

$\xi_\nu^0 \neq 0$
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Expansion of the universe

Chemical Potentials and Flavor Oscillations



Flavor mixing
(neutrino oscillations)

Flavor lepton numbers
not conserved

Only one common neutrino
chemical potential

Stringent ξ_{ν_e} limit
applies to all flavors

$$|\xi_{\nu_{e,\mu,\tau}}| < 0.07$$

Extra neutrino density
 $\Delta N_{\text{eff}} < 0.0064$

Cosmic neutrino density
close to standard value

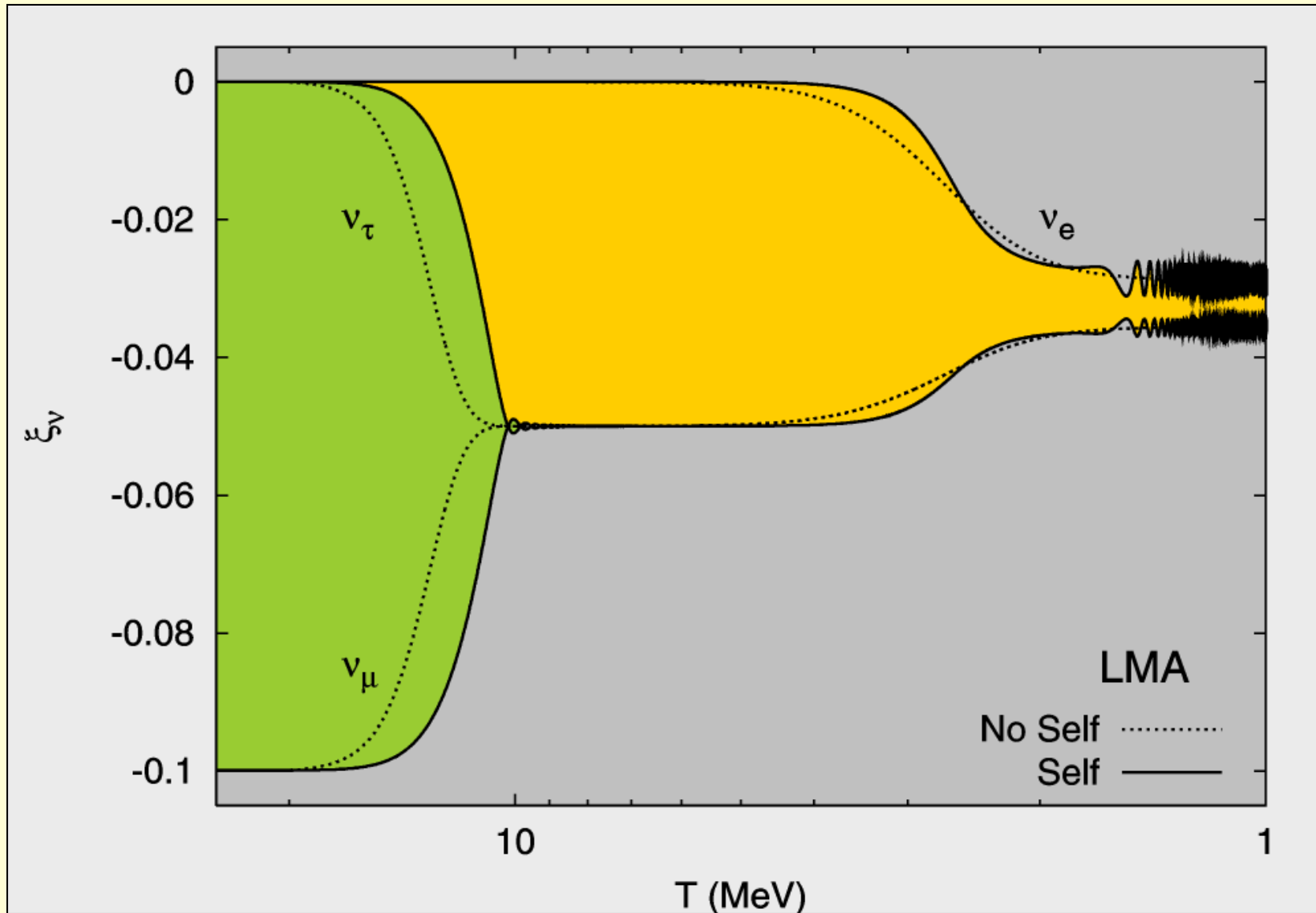
- Oscillations effective before n/p freeze out ?
- YES for solar LMA solution
- Our knowledge of the cosmic neutrino density depends on measured oscillation parameters

- Lunardini & Smirnov, hep-ph/0012056
- Dolgov, Hansen, SP, Petcov, Raffelt & Semikoz, hep-ph/0201287
- Abazajian, Beacom & Bell, astro-ph/0203442
- Wong, hep-ph/0203180

However, important caveats!

SP, Pinto & Raffelt
PRL 102 (2009) 241302 [arXiv:0808.3137]

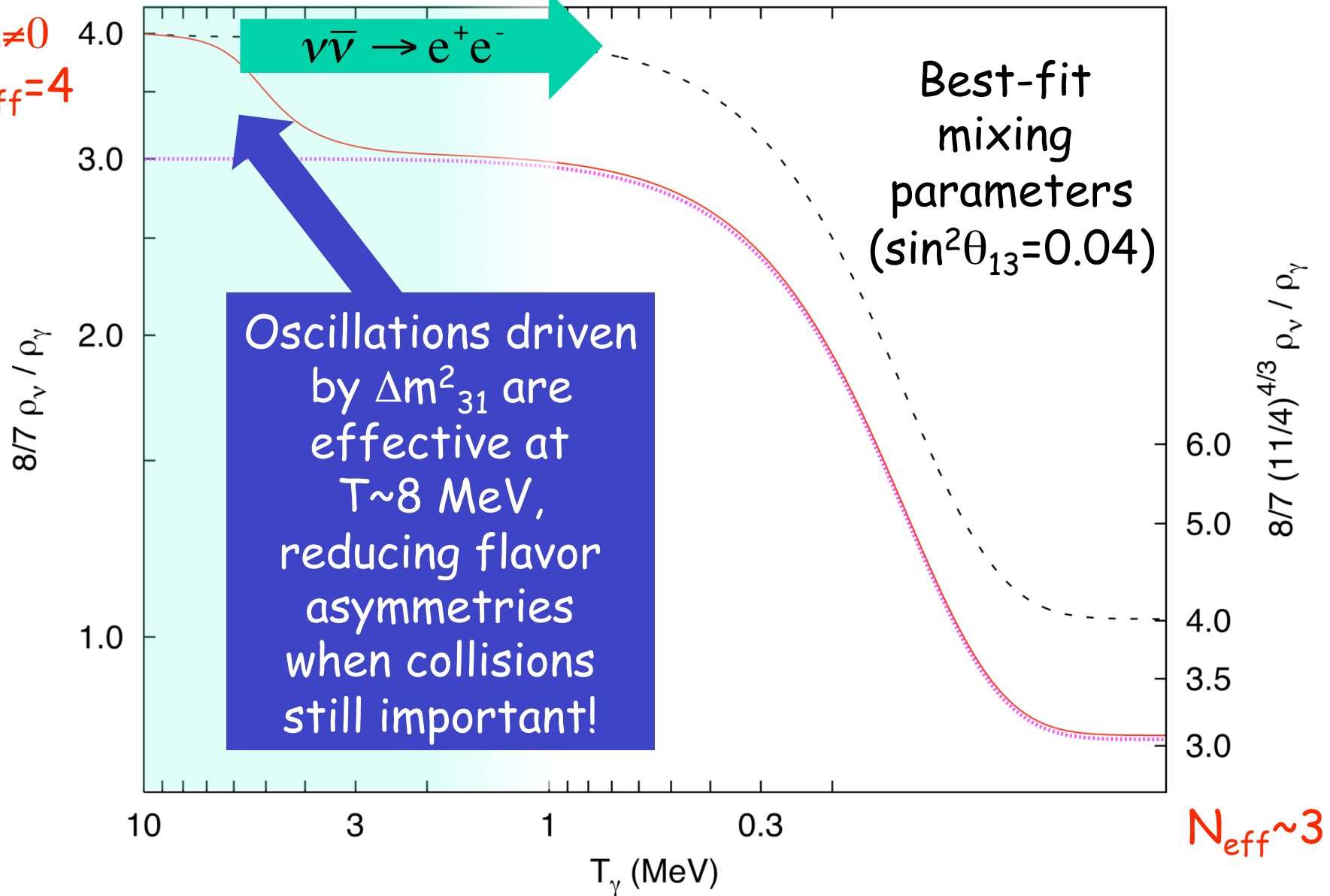
Flavor Transformation for LMA Solution



Dolgov, Hansen, SP, Petcov, Raffelt & Semikoz, hep-ph/0201287

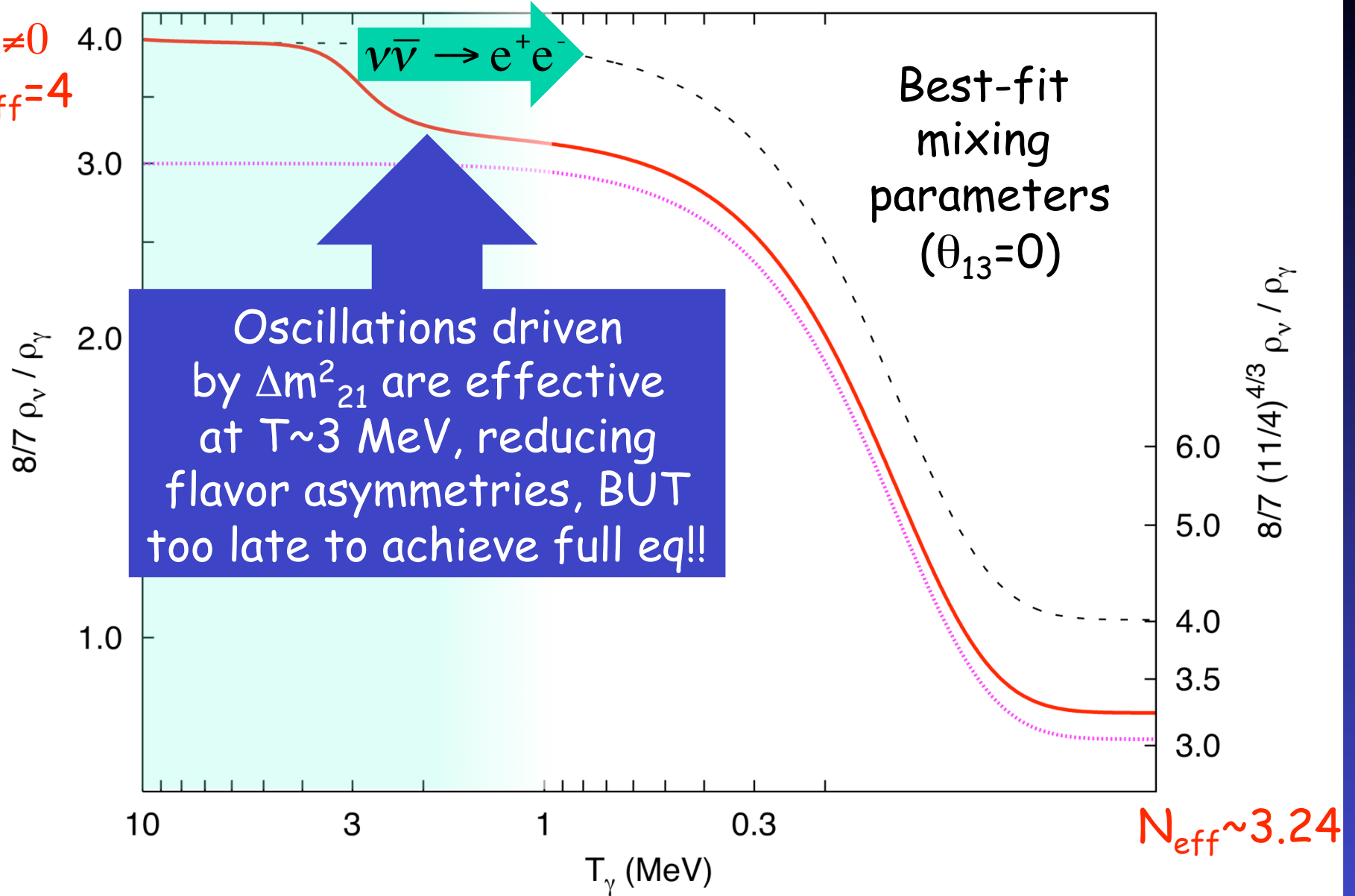
Evolution of ρ_ν/ρ_γ with asymmetries ($L_\nu^{\text{tot}}=0$)

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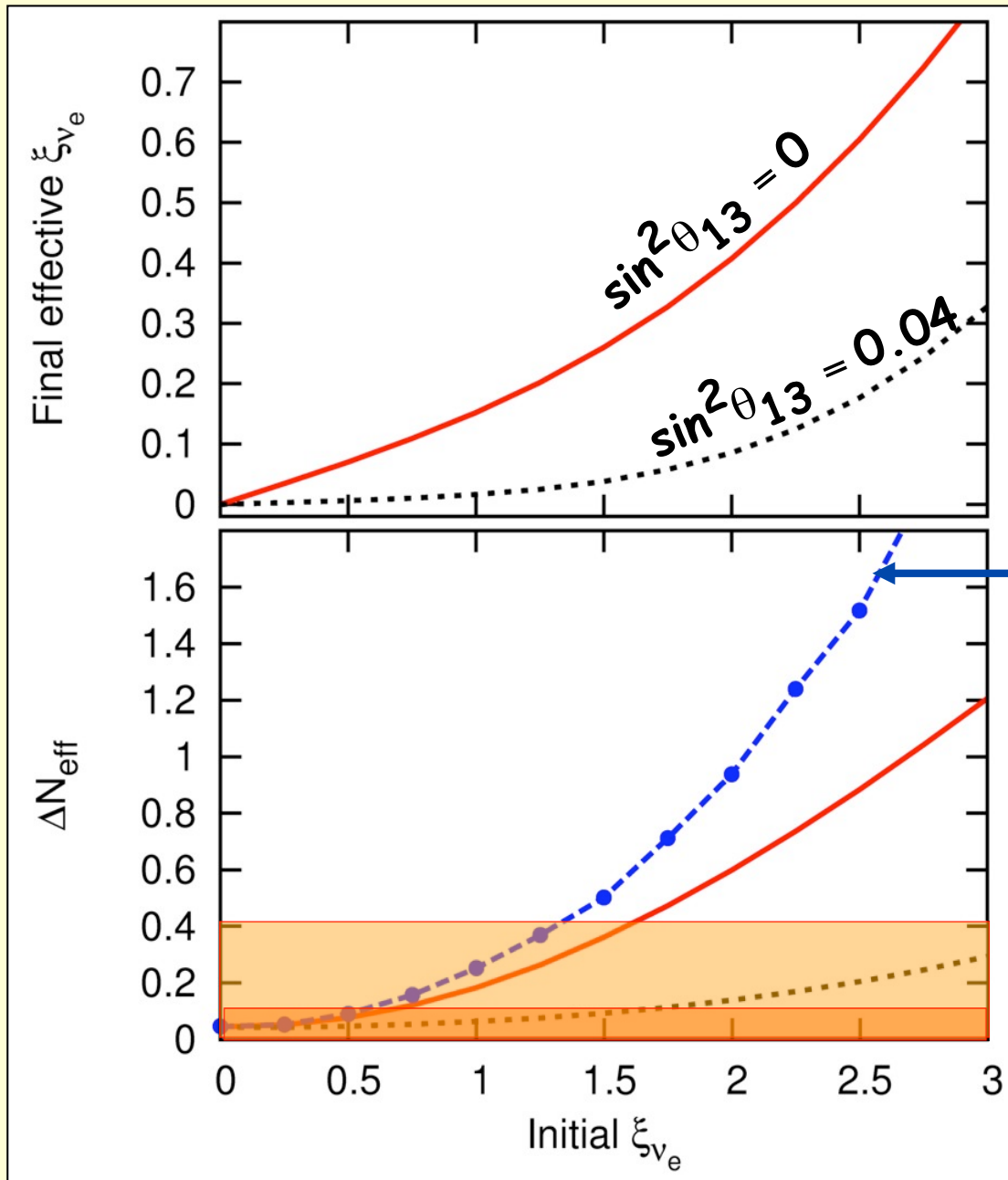


Evolution of ρ_ν/ρ_γ with asymmetries ($L_\nu^{\text{tot}}=0$)

$\xi_\nu^0 \neq 0$
 $N_{\text{eff}}^0 = 4$



Final radiation density



SP, Pinto & Raffelt
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- No global lepton asymmetry
 $n(\nu_e) = n(\bar{\nu}_\mu) + n(\bar{\nu}_\tau)$
- Equilibrium between ν_μ and ν_τ
 $n(\bar{\nu}_\mu) = n(\bar{\nu}_\tau)$

Initial ν_e and ν_μ/ν_τ asymmetries not exactly opposite, but adjusted such that $\xi(\nu_e)$ falls into the allowed BBN range

PLANCK sensitivity: $N_{\text{eff}} \sim 0.4$

Best CMB sensitivity: $N_{\text{eff}} \sim 0.1$

Summary

- **Flavor neutrino oscillations with the measured mixing parameters do not lead to full equilibrium before BBN, in contrast to what is frequently stated**
- **A 13-mixing angle close to the Chooz limit strongly helps to establish equilibrium**
- **The initial densities of ν_e and $\bar{\nu}_\mu$ plus $\bar{\nu}_\tau$ can be adjusted such that at BBN the ν_e asymmetry falls into the allowed range, yet there is a significant excess radiation density**