

Computer tools in particle physics

- Lecture 1 : SARAH and SPheno -

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SARAH

SARAH



[Staub]

- **Name of the tool:** SARAH
- **Author:** Florian Staub (florian.staub@cern.ch)
- **Type of code:** Mathematica package
- **Website:** <http://sarah.hepforge.org/>

SARAH

SARAH



[Staub]

- Lagrangian derivation: SUSY and non-SUSY models
- Mass matrices
- All vertices
- Tadpole equations
- 1-loop corrections for tadpoles and self-energies
- 2-loop renormalization group equations
- 1-loop Wilson coefficients for flavor observables
- Input files for other codes



Crucial for this course!

Models already in SARAH

Supersymmetric Models

- MSSM [in several versions]
- NMSSM
- Near-to-minimal SSM (near-MSSM)
- General singlet extended SSM (SMSSM)
- DiracNMSSM
- Triplet extended MSSM/NMSSM
- Several models with R-parity violation
- U(1)-extended MSSM (UMSSM)
- Secluded MSSM
- Several B-L extended models
- Inverse and linear seesaws [several embeddings]
- MSSM/NMSSM with Dirac Gauginos
- Minimal R-Symmetric SSM
- Minimal Dirac Gaugino SSM
- Seesaws I-II-III [SU(5) versions]
- Left-right symmetric model
- Quiver model

Non-Supersymmetric Models

- Standard Model
- Inert Higgs doublet model
- B-L extended SM
- B-L extended SM with inverse seesaw
- SM extended by a scalar color octet
- Two Higgs doublet model
- Singlet extended SM
- Singlet Scalar DM

SARAH



<http://sarah.hepforge.org/>

The scotogenic model

Also known as...

The inert doublet model

The radiative seesaw

Ma's model

The scotogenic model

[Ernest Ma, 2006]

Field	$SU(2)_L \times U(1)_Y$	Z_2
L_i	$(2, -1/2)$	+
e_i	$(1, 1)$	+
ϕ	$(2, -1/2)$	+
N_i	$(1, 0)$	−
η	$(2, -1/2)$	−

σΚÓΤΟΣ
skotos = darkness



← Inert (or dark) doublet

**Dark
Matter!**

$$\mathcal{L}_N = \overline{N_i} \not{\partial} N_i - \frac{m_{N_i}}{2} \overline{N_i^c} N_i + y_{i\alpha} \eta \overline{N_i} \ell_\alpha + \text{h.c.}$$

$$\begin{aligned} \mathcal{V} = & m_\phi^2 \phi^\dagger \phi + m_\eta^2 \eta^\dagger \eta + \frac{\lambda_1}{2} (\phi^\dagger \phi)^2 + \frac{\lambda_2}{2} (\eta^\dagger \eta)^2 + \lambda_3 (\phi^\dagger \phi) (\eta^\dagger \eta) \\ & + \lambda_4 (\phi^\dagger \eta) (\eta^\dagger \phi) + \frac{\lambda_5}{2} \left[(\phi^\dagger \eta)^2 + (\eta^\dagger \phi)^2 \right] \end{aligned}$$

The scotogenic model

[Ernest Ma, 2006]

$$\mathcal{V} = m_\phi^2 \phi^\dagger \phi + m_\eta^2 \eta^\dagger \eta + \frac{\lambda_1}{2} (\phi^\dagger \phi)^2 + \frac{\lambda_2}{2} (\eta^\dagger \eta)^2 + \lambda_3 (\phi^\dagger \phi) (\eta^\dagger \eta) \\ + \lambda_4 (\phi^\dagger \eta) (\eta^\dagger \phi) + \frac{\lambda_5}{2} [(\phi^\dagger \eta)^2 + (\eta^\dagger \phi)^2]$$

Inert scalar sector: $\eta^\pm$ $\eta^0 = (\eta_R + i\eta_I)/\sqrt{2}$

$$\begin{aligned} m_{\eta^+}^2 &= m_\eta^2 + \lambda_3 \langle \phi^0 \rangle^2 \\ m_R^2 &= m_\eta^2 + (\lambda_3 + \lambda_4 + \lambda_5) \langle \phi^0 \rangle^2 \\ m_I^2 &= m_\eta^2 + (\lambda_3 + \lambda_4 - \lambda_5) \langle \phi^0 \rangle^2 \end{aligned} \quad \Rightarrow \quad m_R^2 - m_I^2 = 2\lambda_5 \langle \phi^0 \rangle^2$$

Radiative neutrino masses

[Ernest Ma, 2006]

Tree-level:

Forbidden by the Z_2 symmetry

Radiative generation of
neutrino masses

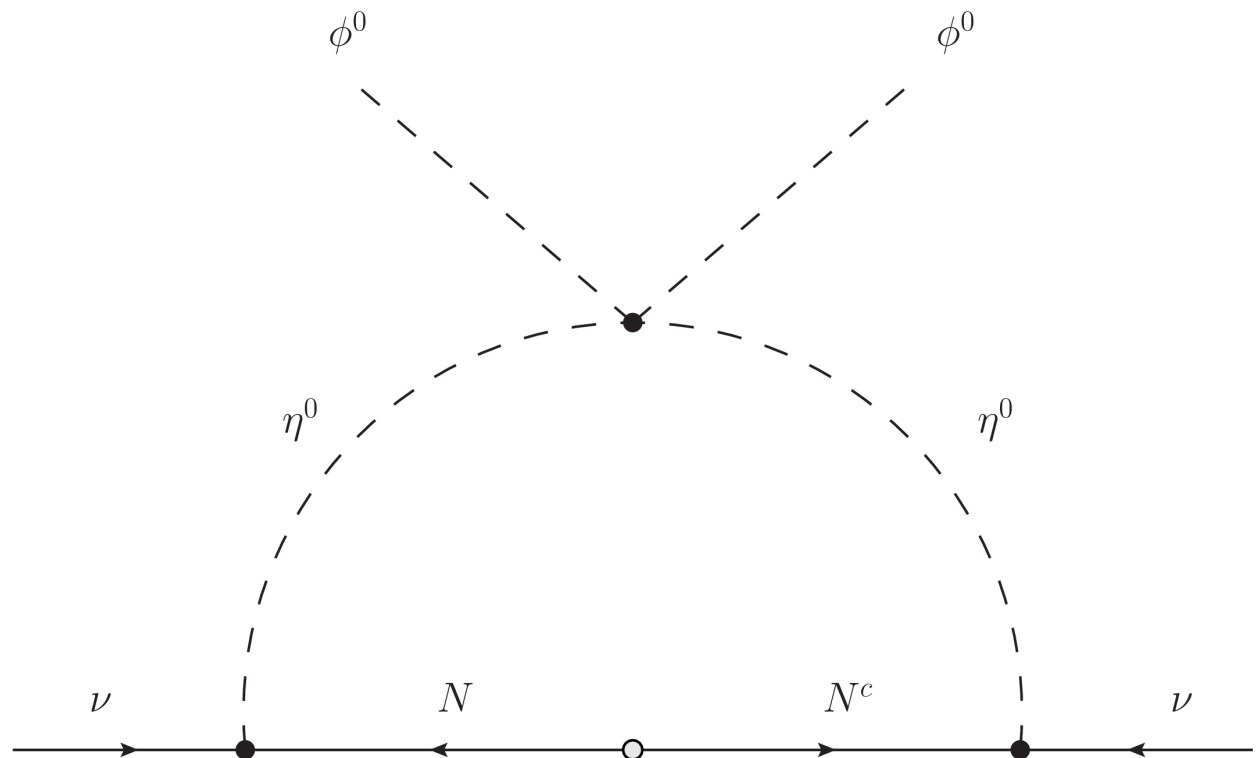


Additional
loop suppression

Dark particles in
the loop

[Other variations in
Restrepo et al, 2013]

1-loop neutrino masses:



Dark matter

The lightest particle charged under Z_2 is stable: **dark matter candidate**

Fermion Dark Matter: N_1

- It can only be produced via **Yukawa** interactions
- Potential problems with lepton flavor violation: is it compatible with the current bounds?

J. Kubo, E. Ma, D. Suematsu, PLB 642 (2006) 18, D. Aristizabal Sierra, J. Kubo, D. Restrepo, D. Suematsu, O. Zapata, PRD 79 (2009) 013011, D. Suematsu, T. Toma, T. Yoshida, PRD 79 (2009) 093004, D. Schmidt, T. Schwetz, T. Toma, PRD 85 (2012) 073009, A. Vicente, C. E. Yaguna, JHEP 1502 (2015) 144, A. Ibarra, C. E. Yaguna, O. Zapata, arXiv:1601.01163, ...

Scalar Dark Matter: the lightest neutral η scalar, η_R or η_I

- It also has **gauge** interactions
- Not correlated to lepton flavor violation

R. Barbieri, L. J. Hall, V. S. Rychkov, PRD 74 (2006) 015007, M. Cirelli, N. Fornengo, A. Strumia, NPB 753 (2006) 178, L. L. Honorez, E. Nezri, J. F. Oliver, M. H. G. Tytgat, JCAP 0702 (2007) 028, Q.-H. Cao, E. Ma, PRD (2007) 095011, S. Andreas, M. H. G. Tytgat, Q. Swillens, JCAP 0904 (2009) 004, E. Nezri, M. H. G. Tytgat, G. Vertongen, JCAP 0904 (2009) 014, T. Hambye, F.-S. Ling, L. L. Honorez, J. Roche, JHEP 07 (2009) 090, L. L. Honorez, C. E. Yaguna, JHEP 1009 (2010) 046 and JCAP 1101 (2011) 002, S. Kashiwase, D. Suematsu, PRD 86 (2012) 053001, A. Goudelis, B. Herrman, O. Stål, JHEP 1309 (2013) 106, M. Klasen, C. E. Yaguna, J. D. Ruiz-Alvarez, D. Restrepo, O. Zapata, JCAP 1304 (2013) 044, J. Racker, JCAP 1403 (2014) 02, ...

Scotogenic: implementation

$\text{FermionFields}[[1]] = \{q, 3, u_L, d_L, 1/6, 2, 3, 1\};$

generations



$U(1)_Y$



$SU(3)_c$



$SU(2)_L$



$\mathbb{Z}_2$



$$q = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$$

Quark doublet

Scotogenic: implementation

Yukawa Lagrangian

$$\text{LagFer} \equiv \mathcal{L}_Y = Y_d H^\dagger \bar{d} q + Y_e H^\dagger \bar{e} \ell + Y_u H \bar{u} q + Y_N \eta \bar{N} \ell$$

Scotogenic: implementation

Scalar decomposition

$$\begin{aligned} H^0 &= \frac{1}{\sqrt{2}} (v + h + iA) \\ \eta^0 &= \frac{1}{\sqrt{2}} (\eta_R + i\eta_I) \end{aligned}$$

Scotogenic: exploration

Tadpole equations

$$\frac{\partial \mathcal{V}}{\partial v} = 0$$

Scotogenic: exploration

Tadpole equations

$$\frac{\partial \mathcal{V}}{\partial v} = 0$$

$$\frac{1}{2}\lambda_1 v^3 - m_H^2 v = 0 \quad \Rightarrow \quad m_H^2 = \frac{1}{2}\lambda_1 v^2$$

Scotogenic: exploration

Mass matrices

- Charged leptons

$$\begin{pmatrix} -\frac{v(Y_e)_{11}}{\sqrt{2}} & -\frac{v(Y_e)_{21}}{\sqrt{2}} & -\frac{v(Y_e)_{31}}{\sqrt{2}} \\ -\frac{v(Y_e)_{12}}{\sqrt{2}} & -\frac{v(Y_e)_{22}}{\sqrt{2}} & -\frac{v(Y_e)_{32}}{\sqrt{2}} \\ -\frac{v(Y_e)_{13}}{\sqrt{2}} & -\frac{v(Y_e)_{23}}{\sqrt{2}} & -\frac{v(Y_e)_{33}}{\sqrt{2}} \end{pmatrix}$$

Chuck Norris fact of the day

When Chuck Norris crosses the street, cars look both ways for Chuck Norris.



Scotogenic: exploration

Mass matrices

- Right-handed neutrinos

$$\begin{pmatrix} -(M_N)_{11} & -\frac{1}{2} (M_N)_{12} - \frac{1}{2} (M_N)_{21} & -\frac{1}{2} (M_N)_{13} - \frac{1}{2} (M_N)_{31} \\ -\frac{1}{2} (M_N)_{12} - \frac{1}{2} (M_N)_{21} & -(M_N)_{22} & -\frac{1}{2} (M_N)_{23} - \frac{1}{2} (M_N)_{32} \\ -\frac{1}{2} (M_N)_{13} - \frac{1}{2} (M_N)_{31} & -\frac{1}{2} (M_N)_{23} - \frac{1}{2} (M_N)_{32} & -(M_N)_{33} \end{pmatrix}$$

$$\begin{pmatrix} -(M_N)_{11} & -(M_N)_{12} & -(M_N)_{13} \\ -(M_N)_{12} & -(M_N)_{22} & -(M_N)_{23} \\ -(M_N)_{13} & -(M_N)_{23} & -(M_N)_{33} \end{pmatrix}$$

Scotogenic: exploration

Higgs boson mass

$$m_h^2 = \frac{3}{2}\lambda_1 v^2 - m_H^2 \Rightarrow m_h^2 = \lambda_1 v^2$$

Tadpole equations

Scotogenic: exploration

$$\ell_i^+ - \ell_j^- - h$$

Vertices

$$\frac{i}{\sqrt{2}} \sum_{m,n=1}^3 (V_e)_{jn}^* (Y_e)_{mn} (U_e)_{im}^* P_L + \frac{i}{\sqrt{2}} \sum_{m,n=1}^3 (V_e)_{in} (Y_e)_{mn}^* (U_e)_{jm} P_R$$

$$\ell_i^+ - \nu_j - W_\mu^-$$

$$-i \frac{g_2}{\sqrt{2}} \sum_{m=1}^3 (V_e)_{im} (V_\nu)_{jm}^* \gamma_\mu P_L = i \frac{g_2}{\sqrt{2}} \sum_{m=1}^3 K_{ij} \gamma_\mu P_L$$

$$\nu_i - \chi_j - \eta_R$$

$$-\frac{i}{\sqrt{2}} \sum_{m,n=1}^3 (V_\nu)_{in}^* (Y_N)_{mn} (Z_X)_{jm}^* P_L - \frac{i}{\sqrt{2}} \sum_{m,n=1}^3 (V_\nu)_{in} (Y_N)_{mn}^* (Z_X)_{jm} P_R$$

Scotogenic: exploration

Renormalization group equations

$$\frac{dc}{dt} = \beta_c = \frac{1}{16\pi^2} \beta_c^{(1)} + \frac{1}{(16\pi^2)^2} \beta_c^{(2)} + \dots$$

Scotogenic: exploration

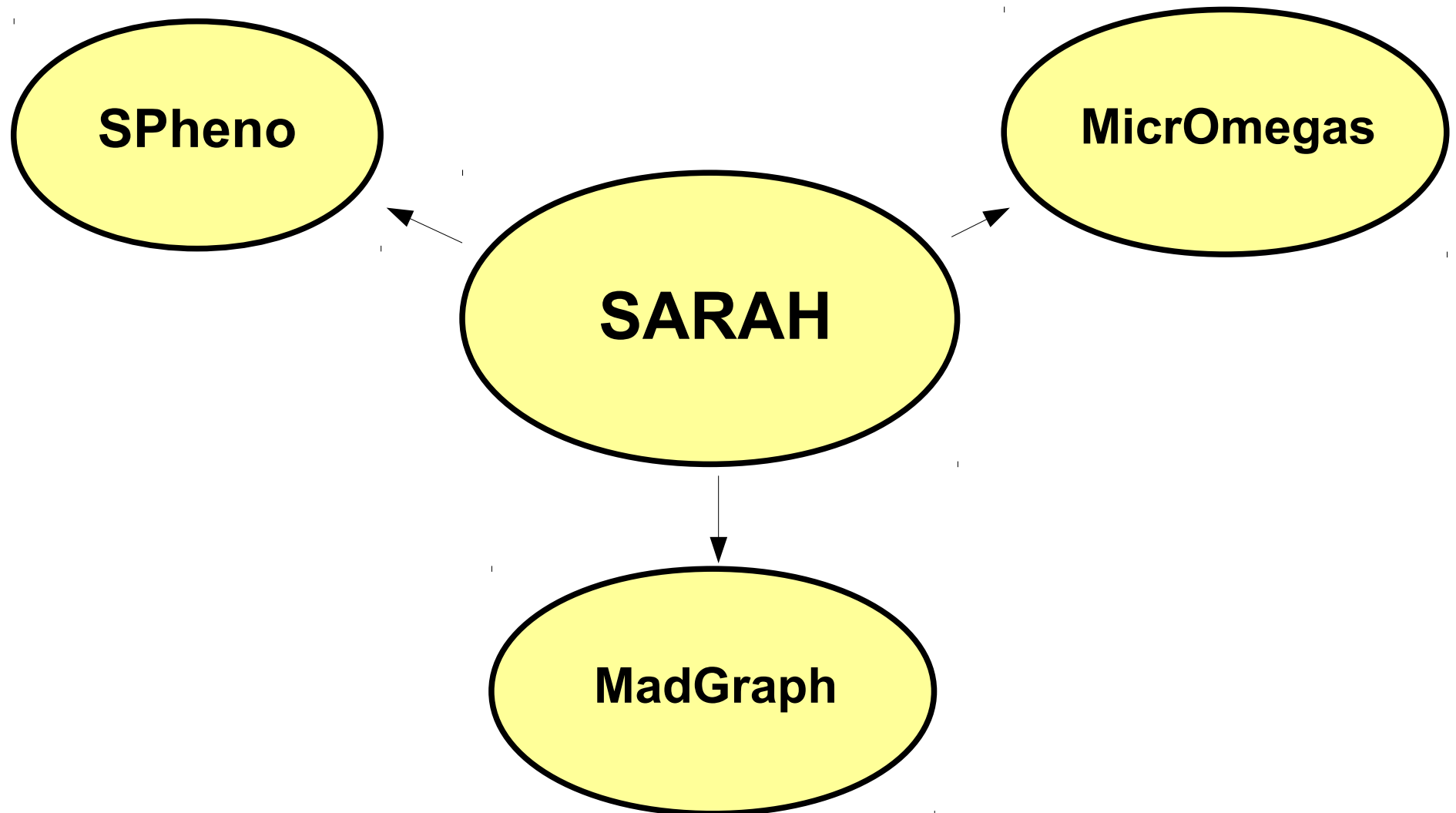
Renormalization group equations

$$\frac{dc}{dt} = \beta_c = \frac{1}{16\pi^2} \beta_c^{(1)} + \frac{1}{(16\pi^2)^2} \beta_c^{(2)} + \dots$$

$$\beta_{g_i}^{(1)} = \left(\frac{21}{5} g_1^3, -3g_2^3, -7g_3^3 \right)$$

$$\begin{aligned} \beta_{m_\eta^2}^{(1)} = & -\frac{9}{2} \left(\frac{1}{5} g_1^2 + g_2^2 \right) m_\eta^2 + 6\lambda_2 m_\eta^2 - 2(2\lambda_3 + \lambda_4) m_H^2 \\ & + 2m_\eta^2 \text{Tr} \left(Y_N Y_N^\dagger \right) - 4 \text{Tr} \left(M_N M_N^* Y_N Y_N^\dagger \right) \end{aligned}$$

SARAH: Input for other codes



SPheno

SPheno

[Porod, Staub]

- **Name of the tool:** SPheno
- **Authors:** Werner Porod (porod@physik.uni-wuerzburg.de) and Florian Staub (florian.staub@cern.ch)
- **Type of code:** Fortran
- **Website:** <http://spheno.hepforge.org/>

SPheno

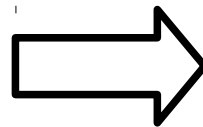
SPheno

[Porod, Staub]

SPheno is a **Fortran code**. It provides routines for the **numerical evaluation** of all vertices, masses and decay modes in a given model.

SARAH

Analytics



SPheno

Numerics

<http://spheno.hepforge.org/>

Scotogenic: benchmark point

BS1 benchmark point

$$\lambda_1 = 0.25$$

$$\lambda_2 = 0.5$$

$$\lambda_3 = 0.5$$

$$\lambda_4 = -0.5$$

$$\lambda_5 = 8 \cdot 10^{-11}$$

$$m_\eta^2 = 1.85 \cdot 10^5 \text{ GeV}^2$$

$$M_N = \begin{pmatrix} 345 \text{ GeV} & 0 & 0 \\ 0 & 4800 \text{ GeV} & 0 \\ 0 & 0 & 6800 \text{ GeV} \end{pmatrix}$$

$$Y_N = \begin{pmatrix} 0.0172495 & 0.300325 & 0.558132 \\ -0.891595 & 1.00089 & 0.744033 \\ -1.39359 & 0.207173 & 0.253824 \end{pmatrix}$$

Backup

Radiative neutrino masses

[Ernest Ma, 2006]

[See also Merle, Platscher, 2015]

$$m_\nu = y^T \Lambda y$$

$$\Lambda_{ij} = \frac{m_{N_i}}{2(4\pi)^2} \left[\frac{m_R^2}{m_R^2 - m_{N_i}^2} \log \left(\frac{m_R^2}{m_{N_i}^2} \right) - \frac{m_I^2}{m_I^2 - m_{N_i}^2} \log \left(\frac{m_I^2}{m_{N_i}^2} \right) \right] \delta_{ij}$$

$$y = \sqrt{\Lambda}^{-1} R \sqrt{\hat{m}_\nu} U_{\text{PMNS}}^\dagger$$

Modified
Casas-Ibarra parameterization
[Toma, Vicente, 2013]

$$U_{\text{PMNS}}^T m_\nu U_{\text{PMNS}} = m_\nu^{\text{diag}}$$



Mixing angles θ_{ij} $m_{\nu_1}, \Delta m_{sol}^2, \Delta m_{atm}^2$