

Standard-model prediction of ϵ_K with manifest CKM unitarity

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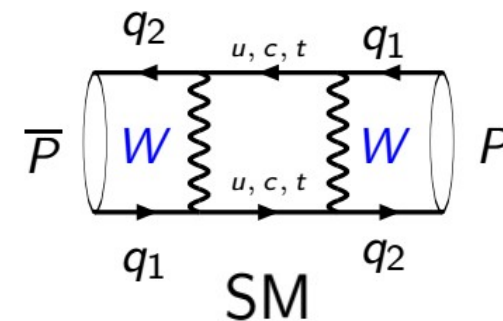
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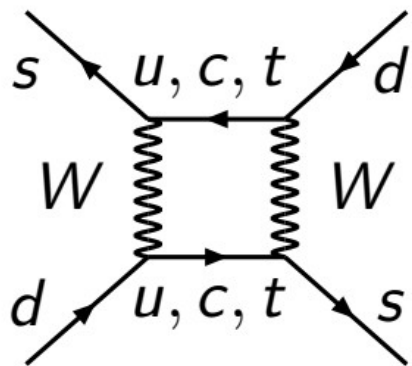
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The parameter ϵ_K describes CP violation in the neutral kaon system and is one of the most sensitive probes of new physics. The large uncertainties related to the charm-quark contribution to ϵ_K have so far prevented a reliable standard-model prediction. We show that CKM unitarity enforces a unique form of the $|\Delta S=2|$ weak effective Lagrangian in which the short-distance theory uncertainty of the imaginary part is dramatically reduced. The uncertainty related to the charm-quark contribution is now at the percent level. We present the updated standard-model prediction $\epsilon_K = 2.16(6)(8)(15) \times 10^{-3}$, where the errors in brackets correspond to QCD short-distance and long-distance, and parametric uncertainties, respectively.

- ϵ_K : discovery of CP Violation
- ϵ_K : CP Violation in the mixing of kaons
- BSM: new contributions to ϵ_K



ϵK w/o perturbative QCD corrections



$$= i \frac{G_F^2 M_W^2}{4\pi^2} \sum_{i,j=c,t} \underbrace{V_{is}^L V_{id}^{L*} V_{js}^L V_{jd}^{L*}}_{\text{unitarity: no } u \text{ index}} \underbrace{S(x_i, x_j)}_{\text{loop function}} Q_V, \quad x_i = \frac{m_i^2}{M_W^2}$$

Weak Left-Handed Currents: $Q_V = \bar{s} \gamma^\mu P_L d \cdot \bar{s} \gamma_\mu P_L d$

$$\begin{aligned} x_u &\sim 0 \\ x_c &\sim 3 \cdot 10^{-4} \\ x_t &\sim 4 \end{aligned}$$

i, j	t, t	c, t	c, c
$S(x_i, x_j)$	$f_{tt}(x_t)$	$x_c [-\log x_c + f_{ct}(x_t)]$	x_c

- CKM unitarity: $V_{us} \cdot (V_{ud})^* = -V_{ts} \cdot (V_{td})^* - V_{cs} \cdot (V_{cd})^*$
- S: Inami-Lim functions
- GIM mechanism: $S \propto x_t, x_c$ ($x_u=0$)

Taking into account QCD corrections

$\text{Tree-level} + \text{One-loop} + \dots \stackrel{\text{OPE}}{=} \sum_i \underbrace{C_i(\mu)}_{\text{short distances}} \times \underbrace{Q_i(\mu)}_{\text{long distances}}$

$C_i(\mu_{\text{high}})$ at high energies,
 $\mu_{\text{high}} = \mathcal{O}(M_{W',W}, M_H, m_t)$
 perturbative

$\langle P | Q_i | \bar{P} \rangle \equiv \langle Q_i(\mu_{\text{low}}) \rangle$ at low energies
 μ_{low} a few GeV
 non-perturbative, Lattice QCD

$$\frac{d}{d \log \mu} \begin{pmatrix} C_1 \\ C_2 \\ \vdots \end{pmatrix} \stackrel{\text{RGE}}{=} \gamma^T \cdot \begin{pmatrix} C_1 \\ C_2 \\ \vdots \end{pmatrix}$$

γ : anom. dim. matrix of Q_i

$$\frac{d}{d \log \mu} \alpha_s \stackrel{\text{LO}}{=} -\beta_0 \frac{\alpha_s^2}{2\pi}$$

$$C(\mu_{\text{low}}) \stackrel{\text{LO}}{=} \left(\frac{\alpha_s(\mu_{\text{low}})}{\alpha_s(\mu_{\text{high}})} \right)^{\frac{\gamma}{2\beta_0}} C(\mu_{\text{high}})$$

↓ running

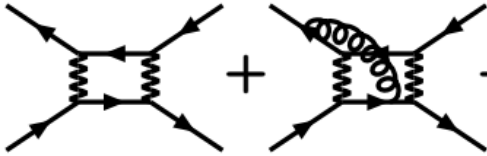
$$\left[\sum_{n=0}^{\infty} \left(\beta_0 \frac{\alpha_s(\mu_{\text{low}})}{2\pi} \log \left(\frac{\mu_{\text{low}}}{\mu_{\text{high}}} \right) \right)^n \right]^{\frac{\gamma}{2\beta_0}}$$

large $\alpha_s \cdot \log(\mu_{\text{low}}/\mu_{\text{high}})$ factors: Short-Distance QCD effects

Impact of the SD QCD corrections in the SM

BSM JC, LVS

The quest for SM precision has required to calculate Short-Distance QCD corrections up to high orders



$$= i \frac{G_F^2 M_W^2}{4\pi^2} \sum_{i,j=c,t} V_{iq_1}^L V_{iq_2}^{L*} V_{jq_1}^L V_{jq_2}^{L*} \times S(x_i, x_j) \bar{\eta}_{ij} Q_V$$

$\bar{\eta}_{ij}$ Short-Distance, perturbative QCD corrections

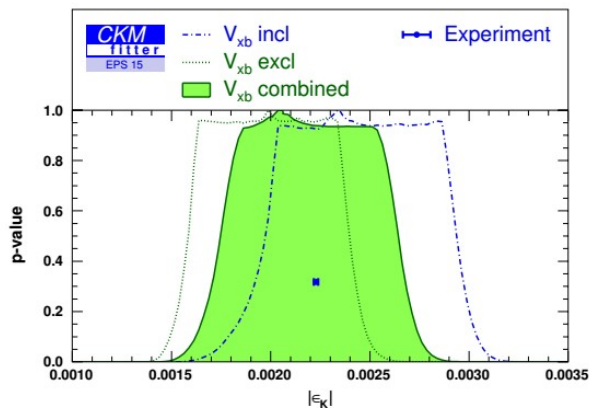
$\bar{\eta}_{ij} = 1 + \mathcal{O}(\alpha_s^m \cdot \log^n x_c) + \dots$: mandatory for phenomenology

SM	t,t	c,t	c,c
η_{LO}^K [Gilman, Wise '83]	0.59	0.37	0.74
η_{NLO}^K [Buras etal '90, Herrlich, Nierste '94 '96]	0.577(7)	$0.47_{-0.04}^{+0.03}$	$1.3_{-0.2}^{+0.3}$
η_{NNLO}^K [Brod, Gorbahn '10 '12]	-	0.50 ± 0.05	1.9 ± 0.8 😞
η_{NLO}^B [Buchalla etal '96]	0.551(2)		

Prediction of ϵ_K

$$\epsilon_K = \frac{G_F^2 F_K^2 M_K M_W^2}{12\sqrt{2}\pi^2 \Delta M_K} \kappa_\epsilon e^{i\phi_\epsilon} \hat{B}_K \sum_{i,j=c,t} \text{Im}\{\lambda_i^{sd} \lambda_j^{sd}\} \eta_{ij} S\left(\frac{\bar{m}_i^2}{M_W^2}, \frac{\bar{m}_j^2}{M_W^2}\right)$$

- QCD: perturbative & non-perturbative (e.g., BK) effects
- $|\epsilon_K| \propto 1 \cdot \eta_{tt} + 0.7 \cdot \eta_{ct} + 0.08 \cdot \eta_{cc}$ (roughly)



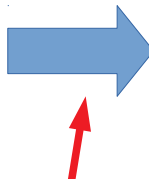
Error budget “c-t” (old evaluation!):

$$|\epsilon_K|_{\text{SM}} = 2.09 \cdot 10^{-3} (1 \pm 4 \% |_{A} \pm 2 \% |_{\lambda} \pm 1 \% |_{\bar{\rho}} \pm 3 \% |_{\bar{\eta}} \pm 2 \% |_{\hat{B}_K} \pm 7 \% |_{\eta_{cc}} \pm 4 \% |_{\eta_{ct}} \pm 1 \% |_{\eta_{tt}} \pm 2 \% |_{\bar{m}_c} \pm 0.5 \% |_{\bar{m}_t} \pm 1 \% |_{F_K} \pm 3 \% |_{\kappa_\epsilon})$$

CKM unitarity: go from “c-t”, to “u-t”

1.) $\lambda_c = -\lambda_t - \lambda_u$ ☹️ $\mathcal{C}_{S_2}^{uu} \equiv \mathcal{C}_1$, $\mathcal{C}_{S_2}^{tt} \equiv \mathcal{C}_2$, and $\mathcal{C}_{S_2}^{ut} \equiv \mathcal{C}_3$.

$$\mathcal{L}_{f=3}^{\Delta S=2} = -\frac{G_F^2 M_W^2}{4\pi^2} [\lambda_c^2 C_{S_2}^{cc}(\mu) + \lambda_t^2 C_{S_2}^{tt}(\mu) + \lambda_c \lambda_t C_{S_2}^{ct}(\mu)] Q_{S_2} + \text{h.c.} + \dots,$$

(5) 

$$\mathcal{L}_{f=3}^{\Delta S=2} = -\frac{G_F^2 M_W^2}{4\pi^2} [\lambda_u^2 \mathcal{C}_{S_2}^{uu}(\mu) + \lambda_t^2 \mathcal{C}_{S_2}^{tt}(\mu) + \lambda_u \lambda_t \mathcal{C}_{S_2}^{ut}(\mu)] Q_{S_2} + \text{h.c.} + \dots, \quad (4)$$

$C_{S_2}^{cc} \equiv \underline{\mathcal{C}_1}$, $C_{S_2}^{ct} \equiv \underline{2\mathcal{C}_1} + \mathcal{C}_3$, and $2C_{S_2}^{tt} \equiv \underline{2\mathcal{C}_1} + \mathcal{C}_2 + \mathcal{C}_3$.

2.) It turns out that η_{uu} , or C1, has a bad convergence (unknown reason)

3.) Same phase convention for $|\Delta S|=1,2$ transitions: “uu” does not contr. to ϵK

- The calculation of $\tilde{\eta}_{tt}$ is essentially the same as for η_{tt} , i.e., $\tilde{\eta}_{tt} \sim \eta_{tt}$
- The calculation of η_{ut} is analogous to the one of η_{ct}
- $|\epsilon K| \propto 1 \cdot \eta_{tt} + 0.6 \cdot \eta_{ut}$ (roughly)

“c-t” vs. “u-t” parameterizations

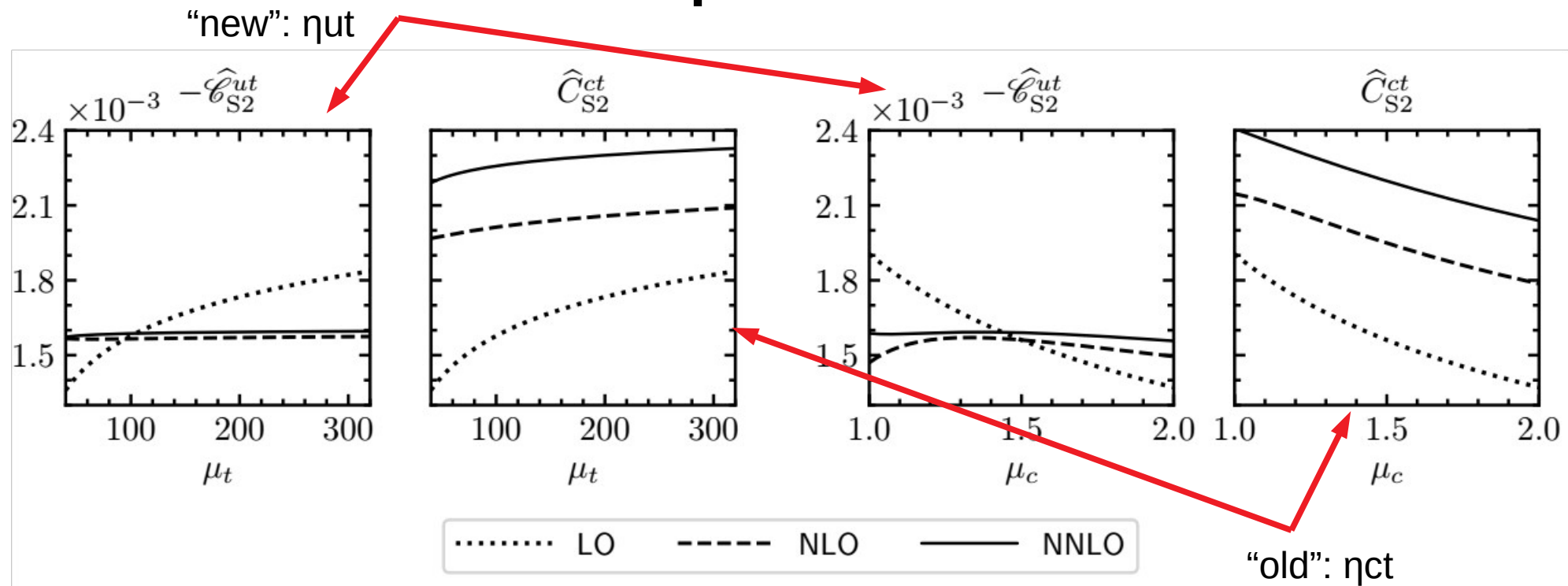


FIG. 1: Comparison of Wilson coefficients in $u-t$ (1st and 3rd plot) and $c-t$ unitarity (2nd and 4th plot). Shown is the residual renormalization-scale dependence of the RI Wilson coefficients as a proxy for their theory uncertainty. In the two plots on the left the five-flavour threshold, μ_t , is varied, while in the two on the right the three-flavour threshold, μ_c , is varied (see text for further details).

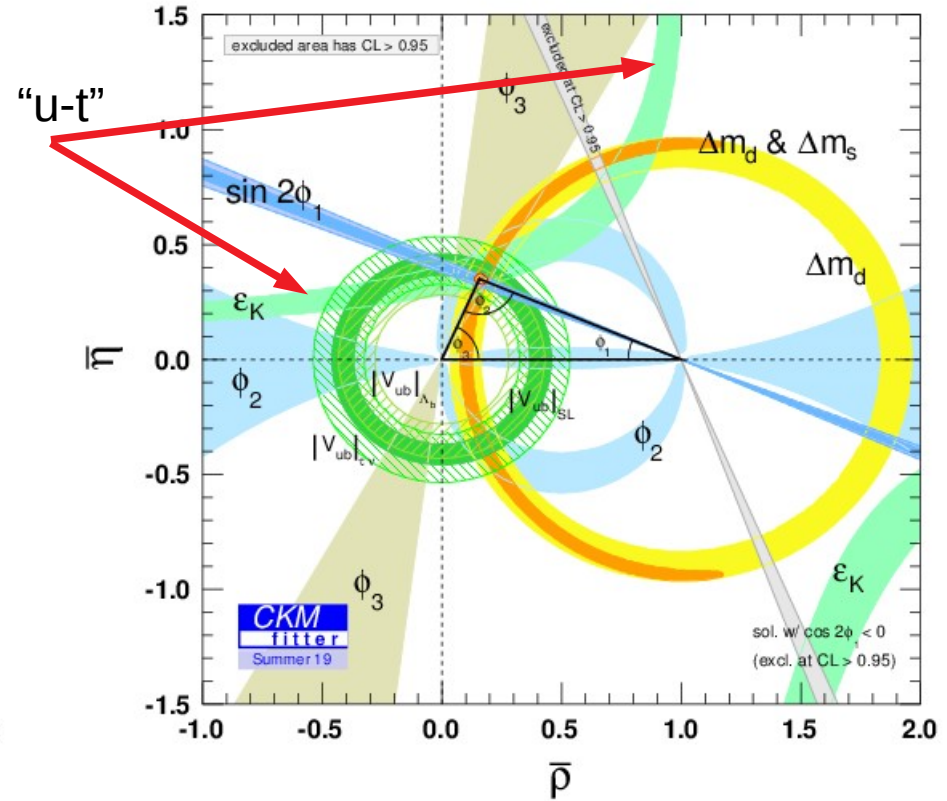
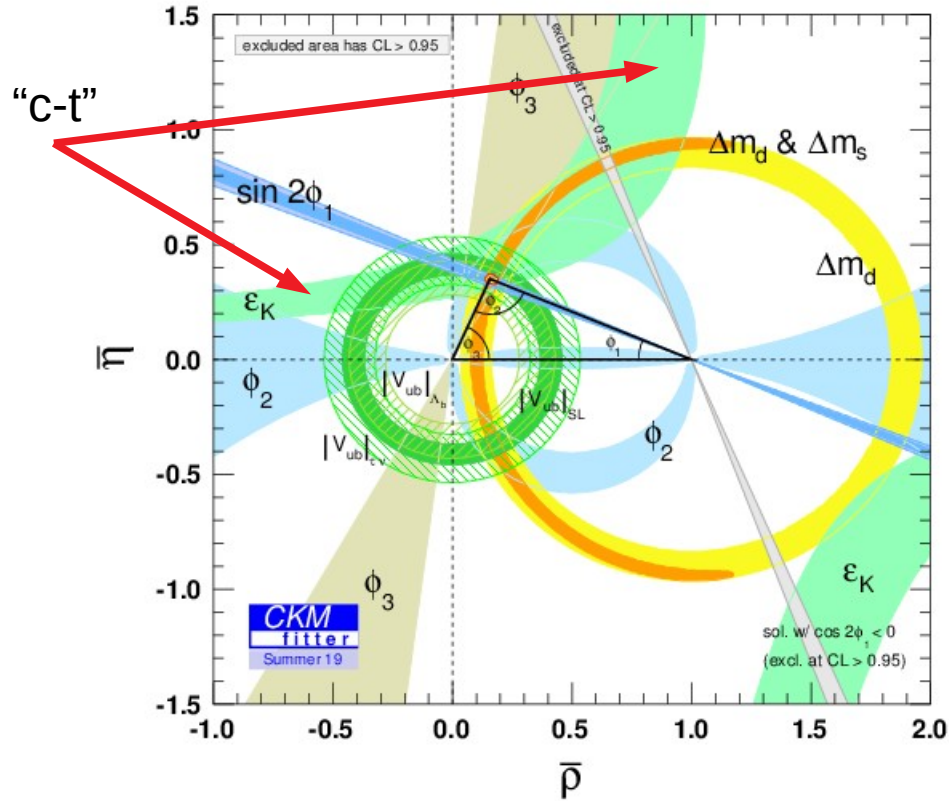
Error budget “u-t”

$$\begin{aligned}\eta_{tt}^{\text{NLL}} &= 0.55(1 \pm 4.2\%_{\text{scales}} \pm 0.1\%_{\alpha_s}), & (16) \\ \eta_{ut}^{\text{NNLL}} &= 0.402(1 \pm 1.3\%_{\text{scales}} \pm 0.2\%_{\alpha_s} \pm 0.2\%_{m_c}).\end{aligned}$$

$$\epsilon_K : \quad \text{CV} \times \left(1 \pm 7.1\%_{\text{param}} \pm 3.0\%_{\eta_{tt}} \pm 0.4\%_{\eta_{ut}} \pm 1.2\%_{\hat{B}_K} \pm 2.4\%_{\xi_s} \pm 2.1\%_{\kappa_\epsilon} \right)$$

Room for improvement: ArXiv: 1911.06822 [hep-ph]

- param: e.g., CKM matrix elements
- Likely (?), Brod & Gorbahn et al. will calculate η_{tt} to the NNLO
- BK: bag parameter
- $\xi_s = (F_{B_s} \sqrt{\hat{B}_s}) / (F_{B_d} \sqrt{\hat{B}_d}) = 1.206(17)$

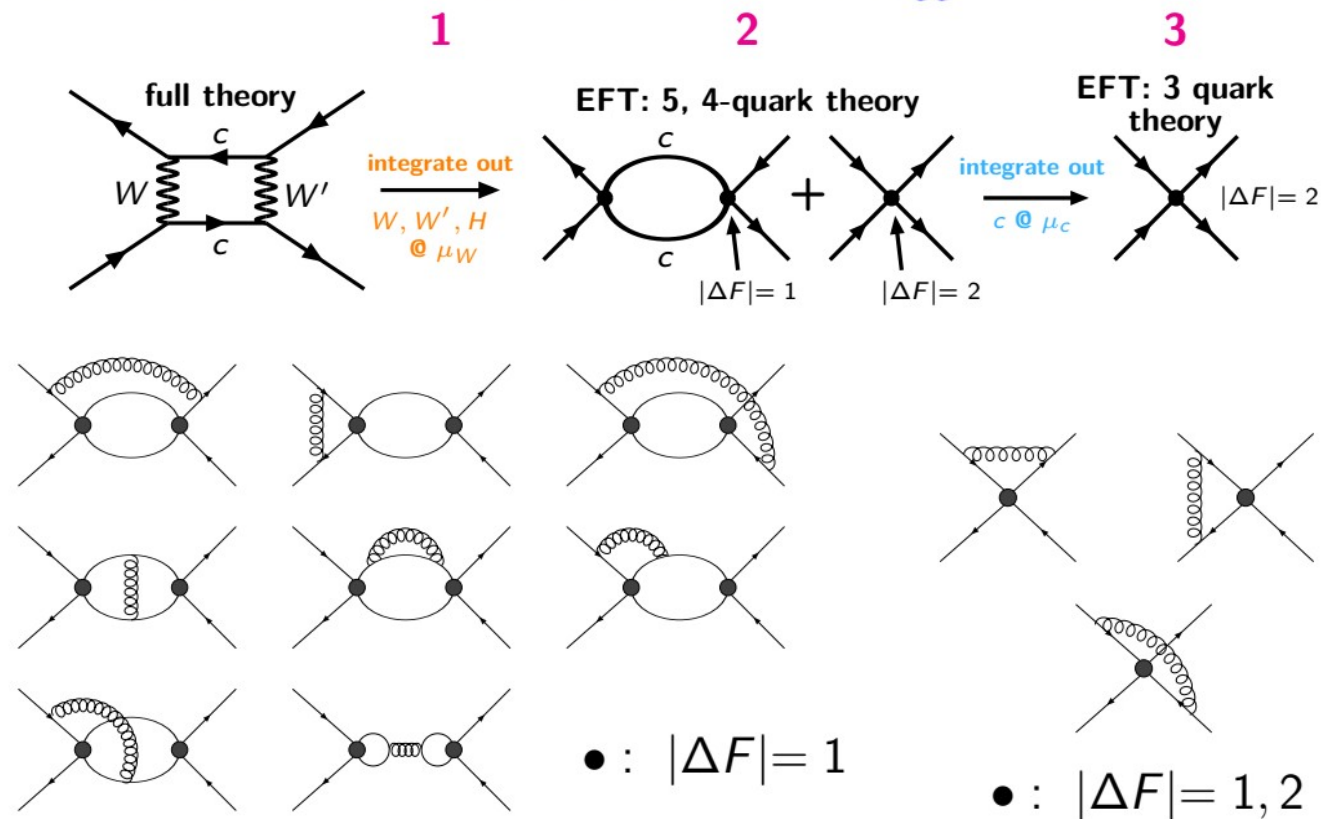


- Experimental measurement: $2.228 \cdot (1 \pm 0.5\%) / 10^3$
- ϵ_K with new param. still in very good agreement with global fit

More details @
arXiv: 1911.06822 [hep-ph]

Spare

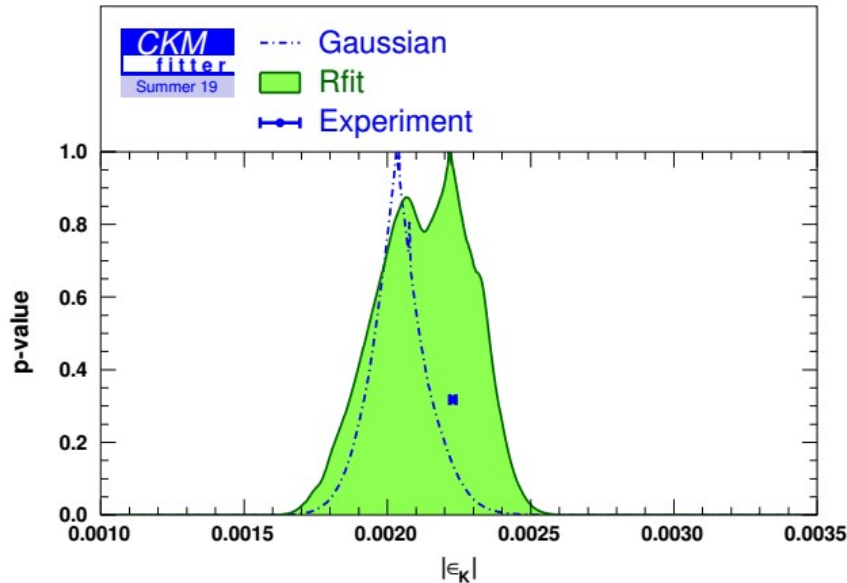
Overview of the calculation for $\bar{\eta}_{cc}^{LR}$ @ NLO



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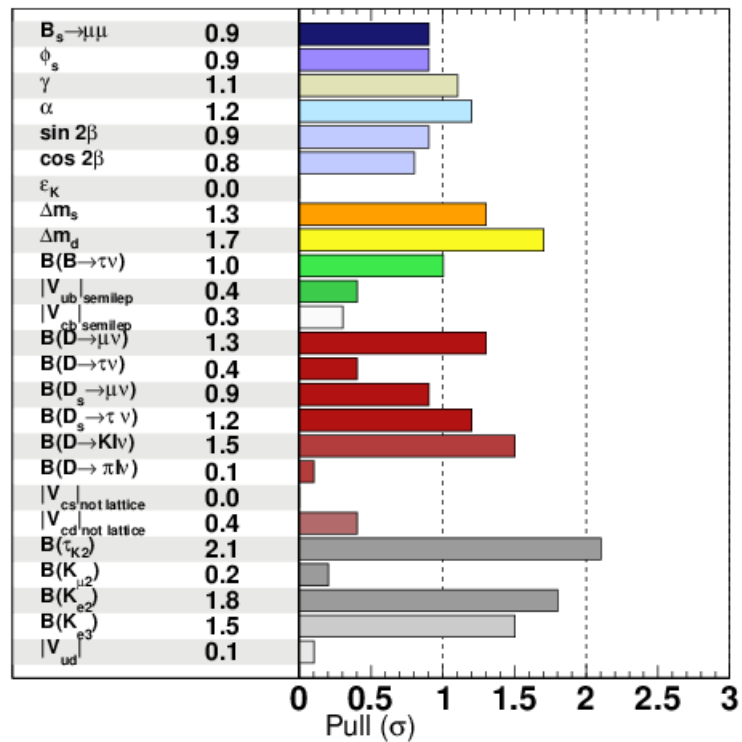
From time to time, question about compatibility of ϵ_K with the rest of the fit, related to the fact that ϵ_K has a strong dependence on

- B_K : role of theoretical uncertainties
- $|V_{cb}|$: inclusive, exclusive or average



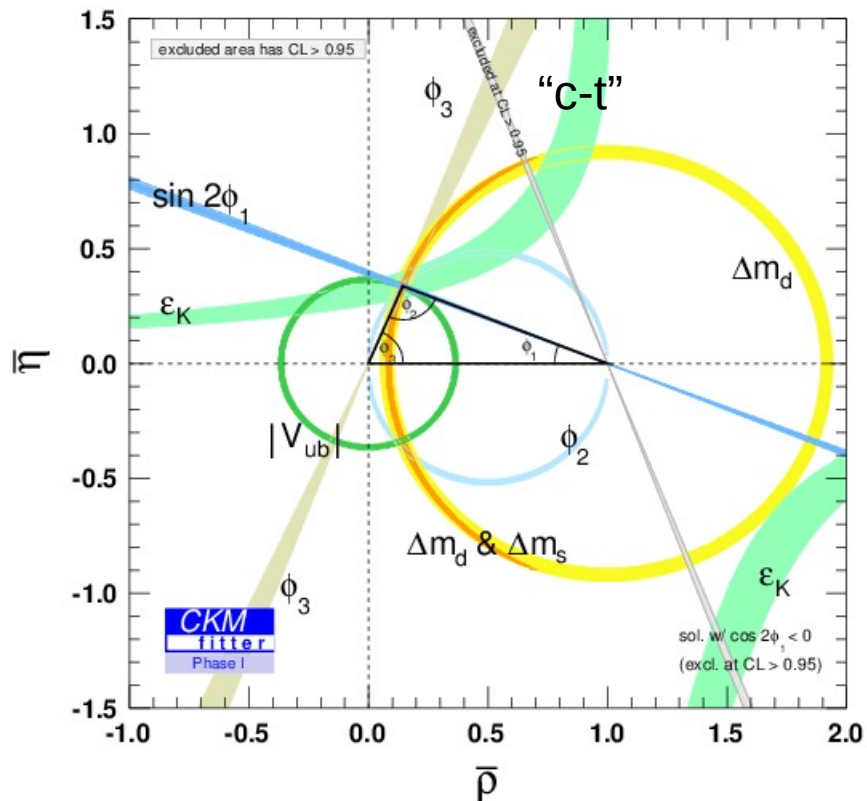
- Rfit versus Gaussian treatment of theoretical uncertainties
- agreement of prediction with experiment in both cases

Spare



Future bounds (LHCb, Belle II, ...)

LHCb at 23 fb⁻¹, CMS/ATLAS at 300 fb⁻¹,
and Belle II at 50 ab⁻¹.



LHCb at 300 fb⁻¹, CMS/ATLAS at 3000 fb⁻¹,
and Belle II at 50 ab⁻¹.

