#### Bounds on Neutrino-Scalar Coupling

Pedro S. Pasquini, Orlando L. G. Peres

Universidade Estadual de Campinas (Unicamp)

#### Encontro Nacional de Física de Partículas e Campos

14 a 18 Setembro 2015, Caxambu, MG



http://www.particlezoo.net/

Index





Index

http://www.particlezoo.net/



http://www.particlezoo.net/



# $\nu$ -Scalar interactions are interesting!

We assume a general Yukawa interaction with (only) neutrinos:

# $\nu$ -Scalar interactions are interesting!

We assume a general Yukawa interaction with (only) neutrinos:

$$\mathcal{L}_{\text{Yukawa}} = \frac{1}{2} g_{\alpha\beta} \overline{\nu}_{\alpha} \nu_{\beta} \phi_1 + \frac{i}{2} g_{\alpha\beta}' \overline{\nu}_{\alpha} \gamma_5 \nu_{\beta} \phi_2$$

# *v*-Scalar interactions are interesting!

We assume a general Yukawa interaction with (only) neutrinos:

$$\mathcal{L}_{\text{Yukawa}} = \frac{1}{2} g_{\alpha\beta} \overline{\nu}_{\alpha} \nu_{\beta} \phi_1 + \frac{i}{2} g'_{\alpha\beta} \overline{\nu}_{\alpha} \gamma_5 \nu_{\beta} \phi_2$$

This interactions can be found in many models, e. g.

Majoron and Related [1]

[1] Y. Chikashige, R. Mohapatra, and R. Peccei, Phys. Rev. Lett. 45, 1926 (1980).

# v-Scalar interactions are interesting!

We assume a general Yukawa interaction with (only) neutrinos:

 $\mathcal{L}_{\text{Yukawa}} = \frac{1}{2} g_{\alpha\beta} \overline{\nu}_{\alpha} \nu_{\beta} \phi_1 + \frac{i}{2} g'_{\alpha\beta} \overline{\nu}_{\alpha} \gamma_5 \nu_{\beta} \phi_2$ 

This interactions can be found in many models, e. g.

Majoron and Related [1] [1] Y. Chikashige, R. Mohapatra, and R. Peccei, Phys. Rev. Lett. 45, 1926 (1980). Two Higgs Models [2]: [2] arXiv:1507.07550 [hep-ph]

$$P^{\pm} \rightarrow l^{\pm} + \stackrel{(-)}{\nu_l}$$

Meson ←  $R^{\pm} \rightarrow l^{\pm} + \overset{(-)}{\nu_l}$ 





$$P^{\pm} \rightarrow l^{\pm} + \stackrel{(-)}{\nu_l}$$



$$P^{\pm} \rightarrow l^{\pm} + \stackrel{(-)}{\nu_l}$$



$$P \rightarrow l + \nu_m + \phi$$



$$P \rightarrow l + \nu_m + \phi$$



$$\Gamma_{\rm tot} = \Gamma(P \to l\nu) + \Gamma(P \to l\nu\phi)$$

Decay Rate (*P* rest Frame)

$$\Gamma_{\rm tot} = \Gamma(P \to l\nu) + \Gamma(P \to l\nu\phi)$$

Can't be observed on this experiment

$$\Gamma_{\rm tot} = \Gamma(P \to l\nu) + \Gamma(P \to l\nu\phi)$$

$$\Gamma_{\rm tot} = \Gamma(P \to l\nu) \left( 1 + A_{Pl} \sum_{\alpha} |g_{l\alpha}|^2 \right)$$

$$\Gamma_{\rm tot} = \Gamma(P \to l\nu) + \Gamma(P \to l\nu\phi)$$

$$\Gamma_{\rm tot} = \Gamma(P \to l\nu) \left(1 + A_{Pl} g_l^2\right)$$

$$\Gamma_{\rm tot} = \Gamma(P \to l\nu) + \Gamma(P \to l\nu\phi)$$

$$\int_{\Gamma_{\rm tot}} SM$$

$$\int_{\Gamma_{\rm tot}} \Gamma(P \to l\nu) (1 + A_{Pl} g_l^2)$$

$$\begin{split} \Gamma_{\rm tot} &= \Gamma(P \to l\nu) + \Gamma(P \to l\nu\phi) \\ & \text{New Physics} \\ & \Gamma_{\rm tot} &= \Gamma(P \to l\nu) \left(1 + A_{Pl} g_l^2\right) \end{split}$$

$$\Gamma\left(P \to l\nu\right) = \frac{G_F^2 f_P^2 |V_{qq'}|^2 m_p^3}{8\pi} \left(\frac{m_l}{m_p}\right)^2 \left[1 - \left(\frac{m_l}{m_p}\right)^2\right]^2$$



$$\Gamma(P \to l\nu) = \frac{G_F^2 f_P^2 |V_{qq'}|^2 m_p^3}{8\pi} \left(\frac{m_l}{m_p}\right)^2 \left[1 - \left(\frac{m_l}{m_p}\right)^2\right]^2$$











More on that later...  
Meson Decay Constant
$$\Gamma\left(P \to l\nu\right) = \frac{G_{I}^{2}(f_{P}^{2})Y_{qq'}|^{2}m_{p}^{3}}{8\pi} \left(\frac{m_{l}}{m_{p}}\right)^{2} \left[1 - \left(\frac{m_{l}}{m_{p}}\right)^{2}\right]^{2}$$







Reach up to 25%...








9 / 27











### Lattice $\mathsf{QCD}^3$

### Lattice $\mathsf{QCD}^3$

	$f_P[\text{MeV}]$
$\pi$	130.2(1.4)
K	156.3(0.9)
D	209(3.3)
$D_s$	250(7)
B	186(4)

### Lattice QCD<sup>3</sup>



### It is a Three body Decay

Finally, Beyond SM:



Following Barger, Keung and Pakvasa $^4$ 

Following Barger, Keung and Pakvasa $^4$ 

Calculated the corrections  $A_{Pl}(m_{\phi})$ 

Following Barger, Keung and Pakvasa<sup>4</sup> 
$$\beta = \frac{m_{\phi}^2}{m_P^2}, \epsilon = \frac{m_{\nu}^2}{m_P^2}$$

$$d\Gamma(P \to l\nu\phi) = \Gamma_{l\nu} \frac{\left(x^2 + \epsilon^2 + 6x\epsilon - \beta x - \beta\epsilon\right)\lambda^{1/2}(x,\epsilon,\beta)}{x^2(x-\epsilon)^2} \frac{g_l^2}{32\pi^2} dx$$

Following Barger, Keung and Pakvasa<sup>4</sup> 
$$\beta = \frac{m_{\phi}^2}{m_P^2}, \epsilon = \frac{m_{\nu}^2}{m_P^2}$$

$$d\Gamma(P \to l\nu\phi) = \Gamma_{l\nu} \frac{\left(x^2 + \epsilon^2 + 6x\epsilon - \beta x - \beta\epsilon\right)\lambda^{1/2}(x,\epsilon,\beta)}{x^2(x-\epsilon)^2} \frac{g_l^2}{32\pi^2} dx$$
$$\to \lambda(x,y,z) = x^2 + y^2 + z^2 - 2xy - 2xz - 2zy$$

Kinematic Function

Following Barger, Keung and Pakvasa
$$^4$$
  $eta=rac{m_{\phi}^2}{m_P^2}$ ,  $\epsilon=rac{m_{
u}^2}{m_P^2}$ 

$$d\Gamma(P \to l\nu\phi) = \Gamma_{l\nu} \underbrace{\frac{(x^2 + \epsilon^2 + 6x\epsilon - \beta x - \beta\epsilon)\lambda^{1/2}(x,\epsilon,\beta)}{x^2(x-\epsilon)^2} \frac{g_l^2}{32\pi^2} dx}_{\text{Leptonic Decay rate}}$$
for neutrino with Squared mass  $x$ 

$$|V_{eqt}|^{2m_{p}^2} t = t = (x-\epsilon)^{2(\lambda+1/2)} t^{1/2} t = 0$$

$$\begin{split} \Gamma_{l\nu} &= \frac{G_F^2 f_F^2 |V_{qq'}|^2 m_p^3}{8\pi} [x + \alpha - (x - \alpha)^2] \lambda^{1/2}(1, x, \alpha) \\ &\alpha = \frac{m_l^2}{m_P^2} \end{split}$$
4 Phys. Rev. D 25, 907 (1982).

13 / 27

Following Barger, Keung and Pakvasa<sup>4</sup> 
$$\beta = rac{m_{\phi}^2}{m_P^2}, \epsilon = rac{m_{\nu}^2}{m_P^2}$$

$$d\Gamma(P \to l\nu\phi) = \Gamma_{l\nu} \frac{(x^2 + \epsilon^2 + 6x\epsilon - \beta x - \beta\epsilon)\lambda^{1/2}(x,\epsilon,\beta)}{x^2(x-\epsilon)^2} \frac{g_l^2}{32\pi^2} dx$$

#### Three body and Yukawa

<sup>4</sup>Phys. Rev. D **25**, 907 (1982).

# It is like a virtual twobody!

# x varies angles and $M^2_{ m miss}$

Following Barger, Keung and Pakvasa<sup>4</sup> 
$$\beta = \frac{m_{\phi}^2}{m_P^2}, \epsilon = \frac{m_{\nu}^2}{m_P^2}$$

$$d\Gamma(P \to l\nu\phi) = \Gamma_{l\nu} \frac{(x^2 + \epsilon^2 + 6x\epsilon - \beta x - \beta\epsilon)\lambda^{1/2}(x,\epsilon,\beta)}{x^2(x-\epsilon)^2} \frac{g_l^2}{32\pi^2} dx$$

<sup>4</sup>Phys. Rev. D **25**, 907 (1982).

Following Barger, Keung and Pakvasa<sup>4</sup> 
$$\beta = \frac{m_{\phi}^2}{m_P^2}, \epsilon = \frac{m_{\nu}^2}{m_P^2}$$

$$d\Gamma(P \to l\nu\phi) = \Gamma_{l\nu} \frac{(x^2 + \epsilon^2 + 6x\epsilon - \beta x - \beta \epsilon)\lambda^{1/2}(x,\epsilon,\beta)}{x^2(x-\epsilon)} \frac{g_l^2}{32\pi^2} dx$$
Infrared Divergent!!
$${}^4_{\text{Phys. Rev. D 25, 907 (1982).}}$$



## It is not Divergent!



**Decays from**  $\pi, K, D^+, D_s, B$ 

Using a  $\chi^2$ :

$$\chi^{2} = \sum_{P,l} \frac{\left[\Gamma_{Pl}(1 + A_{Pl}g_{l}^{2}) - a_{Pl}\right]^{2}}{\sigma_{Pl}^{2}}$$

Using a  $\chi^2$ :

$$\chi^{2} = \sum_{P,l} \frac{\left[\Gamma_{Pl}(1 + A_{Pl}g_{l}^{2}) - a_{Pl}\right]^{2}}{\sigma_{Pl}^{2}}$$
Experimental Data 8 Decays:

### $\pi, K, D^+, D_s, B$

[5] PDG Chin. Phys. C 38, 090001 (2014).

[6] BESIII Collaboration, Phys. Rev. D 89, no. 5, 051104 (2014)

[7] BaBar Collaboration, arXiv:1003.3063 [hep-ex].

[8] Belle Collaboration, JHEP 1309, 139(2013).

# Bounds on $\pi, K, D^+, D_s, B$



## **Continuous peaks** = $l\nu\phi$ decay

Heavy Neutrino Search allows to use this calculations too!

Heavy Neutrino Search allows to use this calculations too!

The three body decay mimic a continuos mass spectrum of heavy neutrino

Heavy Neutrino Search allows to use this calculations too!

The three body decay mimic a continuos mass spectrum of heavy neutrino



They look for peaks on the SM decay spectrum

# Using Heavy Search from $\pi, K$ [8,9]





22 / 27

## Heavy Search Uses Lepton Spectrum $\rightarrow$ much better

### Conclusions

We improved the bounds on the constants!

Our Results are valid up to  $m_{\phi} \sim 100 \,\,\mathrm{MeV}$ 

Constants	Our Results	[10]	[11]
$ g_e ^2$	$ <1.9\times10^{-6}$	$< 4.4 \times 10^{-5}$	$< 1.6 \times 10^{-5}$
$ g_{\mu} ^2$	$ <1.9\times10^{-7}$	$< 3.6 \times 10^{-4}$	
$ g_{ au} ^2$	< 7.5	$<2.2\times10^{-1}$	

**Table:** Limits setting  $m_{\phi} = 0$ .

[10] A. P. Lessa and O. L. G. Peres, Phys. Rev. D 75, 094001 (2007).

[11] J. B. Albert et al. [EXO-200], Phys. Rev. D 90, no. 9, 092004 (2014).



## Heavy Search Uses Lepton Spectrum $\rightarrow$ much better

### Conclusions

We improved the bounds on the constants!

Our Results are valid up to  $m_{\phi} \sim 100 \text{ MeV}$ 

Constants	Our Results	[10]	[11]
$ g_e ^2$	$ <1.9\times10^{-6}$	$<4.4\times10^{-5}$	$< 1.6 \times 10^{-5}$
$ g_{\mu} ^2$	$ <1.9\times10^{-7}$	$< 3.6 \times 10^{-4}$	
$ g_{ au} ^2$	< 7.5	$< 2.2  imes 10^{-1}$	

**Table:** Limits setting  $m_{\phi} = 0$ .

**Agradecimentos:** 

[10] A. P. Lessa and O. L. G. Peres, Phys. Rev. D 75, 094001 (2007).

[11] J. B. Albert et al. [EXO-200], Phys. Rev. D 90, no. 9, 092004 (2014).

Fapesp CNPQ Unicamp (IFGW) Advisor, Professors and friends



Backup Slides



### **Missing Energy Allows Probe New Physics**

 $x \neq 0$ 



26 / 27

