



UNIVERSITÀ  
DI TORINO



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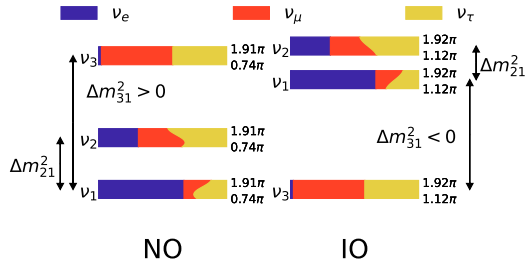
`stefano.gariazzo@unito.it`

## Relic neutrinos: decoupling and direct detection perspectives

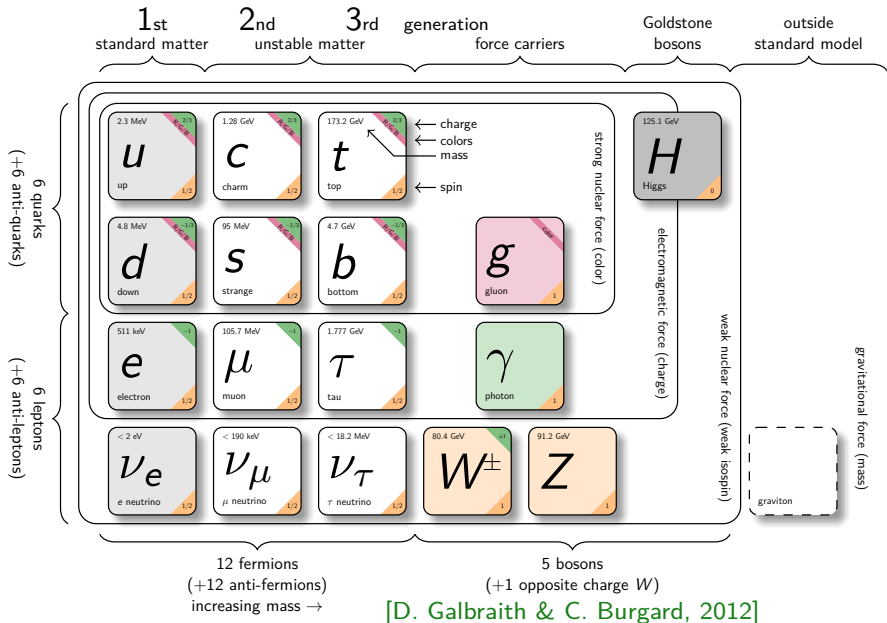
# N Neutrinos

Based on:

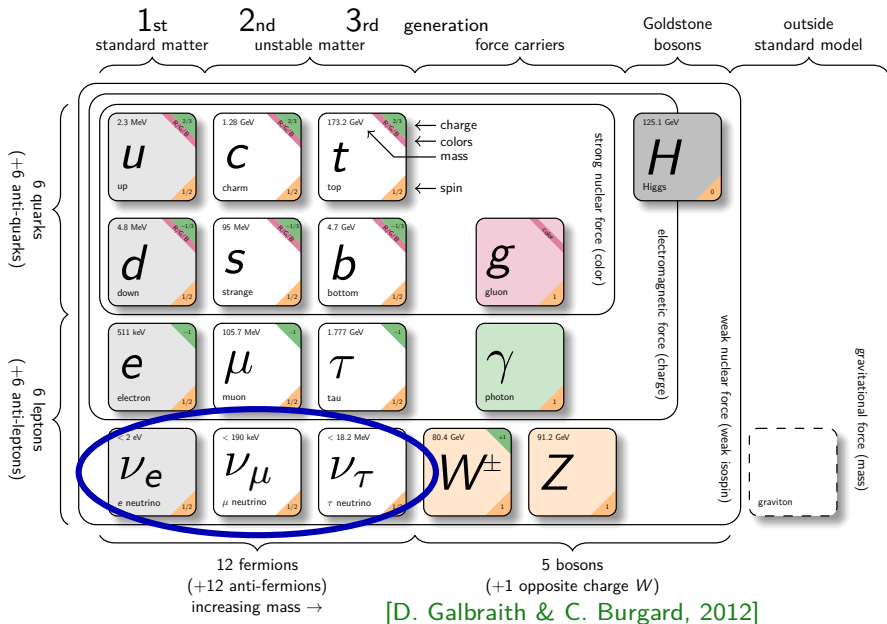
- JHEP 02 (2021) 071 and update

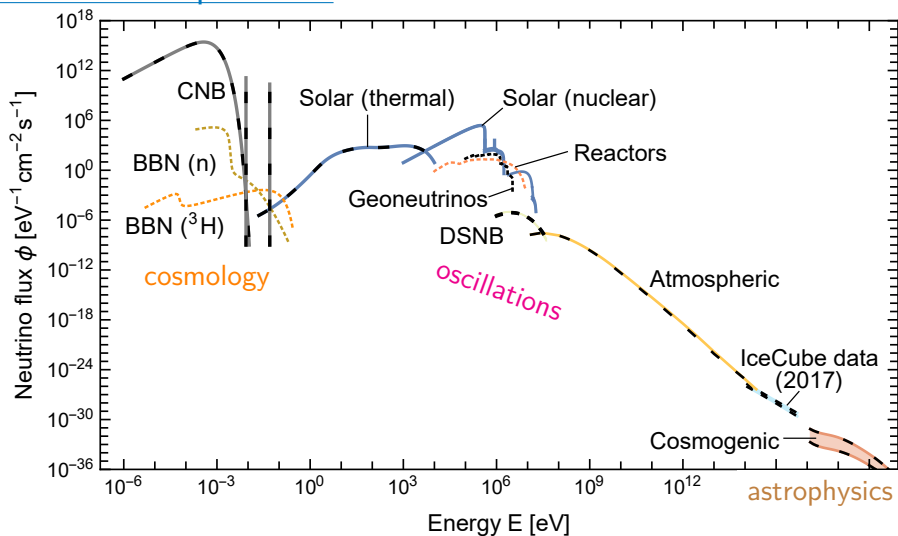


# The Standard Model of Particle Physics



# The Standard Model of Particle Physics





neutrinos at all energies provide valuable information!

# Three Neutrino Oscillations

$$\nu_\alpha = \sum_{k=1}^3 U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau)$$

$U_{\alpha k}$  described by 3 mixing angles  $\theta_{12}$ ,  $\theta_{13}$ ,  $\theta_{23}$  and one CP phase  $\delta$

Current knowledge of the 3 active  $\nu$  mixing: [JHEP 02 (2021) update]

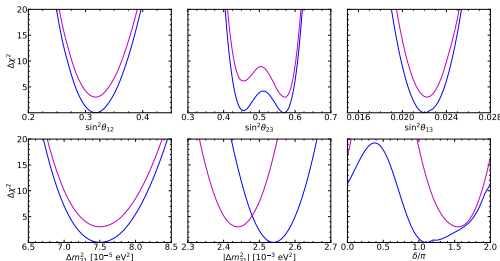
NO/NH: Normal Ordering/Hierarchy,  $m_1 < m_2 < m_3$

IO/IH: Inverted O/H,  $m_3 < m_1 < m_2$

$$\begin{aligned} \Delta m_{21}^2 &= (7.50^{+0.22}_{-0.20}) \cdot 10^{-5} \text{ eV}^2 \\ |\Delta m_{31}^2| &= (2.54 \pm 0.03) \cdot 10^{-3} \text{ eV}^2 \text{ (NO)} \\ &= (2.44 \pm 0.03) \cdot 10^{-3} \text{ eV}^2 \text{ (IO)} \end{aligned}$$

$$\begin{aligned} 10 \sin^2(\theta_{12}) &= 3.18 \pm 0.16 \\ 10^2 \sin^2(\theta_{13}) &= 2.200^{+0.069}_{-0.062} \text{ (NO)} \\ &= 2.225^{+0.064}_{-0.070} \text{ (IO)} \\ 10 \sin^2(\theta_{23}) &= 4.55 \pm 0.13 \text{ (NO)} \\ &= 5.71^{+0.14}_{-0.17} \text{ (IO)} \end{aligned}$$

$$\begin{aligned} \delta/\pi &= 1.10^{+0.27}_{-0.12} \text{ (NO)} \\ &= 1.54 \pm 0.14 \text{ (IO)} \end{aligned}$$



mass ordering  
still unknown

$\delta$  still unknown

see also: <http://globalfit.astroparticles.es>

## Normal ordering (NO)

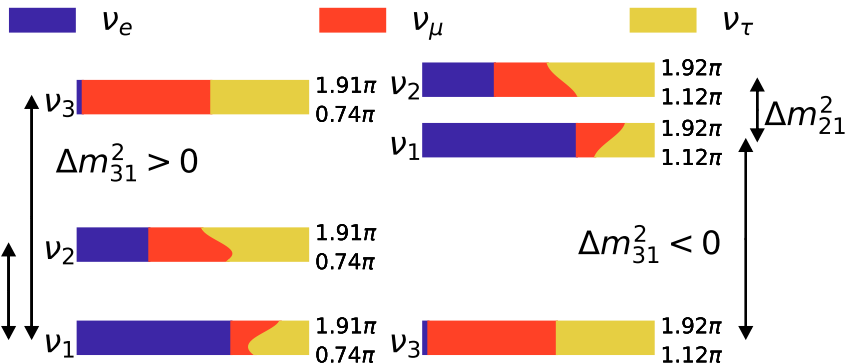
$$m_1 < m_2 < m_3$$

$$\sum m_k \gtrsim 0.06 \text{ eV}$$

## Inverted ordering (IO)

$$m_3 < m_1 < m_2$$

$$\sum m_k \gtrsim 0.1 \text{ eV}$$



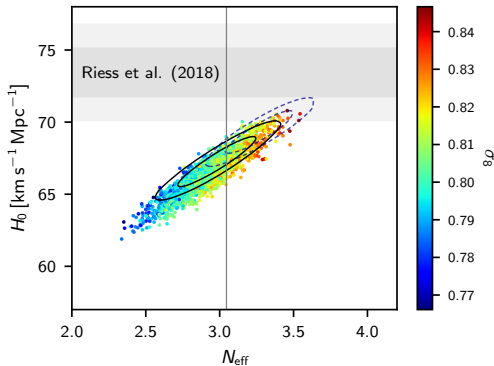
Absolute scale unknown!

Can we constrain the mass ordering using bounds on  $\sum m_\nu$ ?

# E Neutrinos in the Early Universe

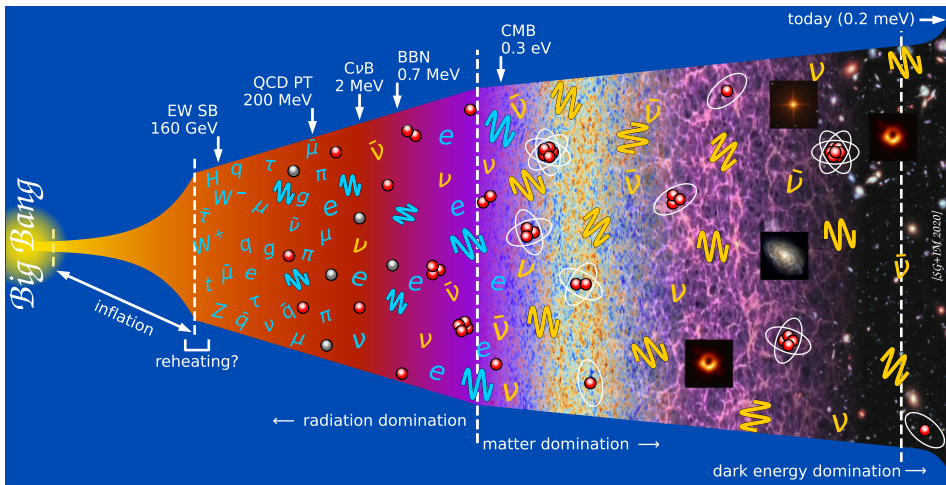
Based on:

- Planck 2018
- JCAP 04 (2021) 073

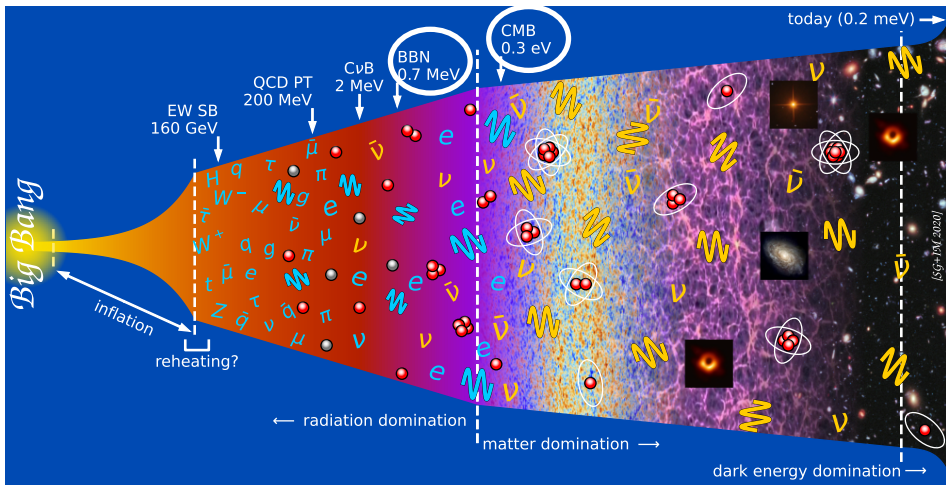




# History of the universe

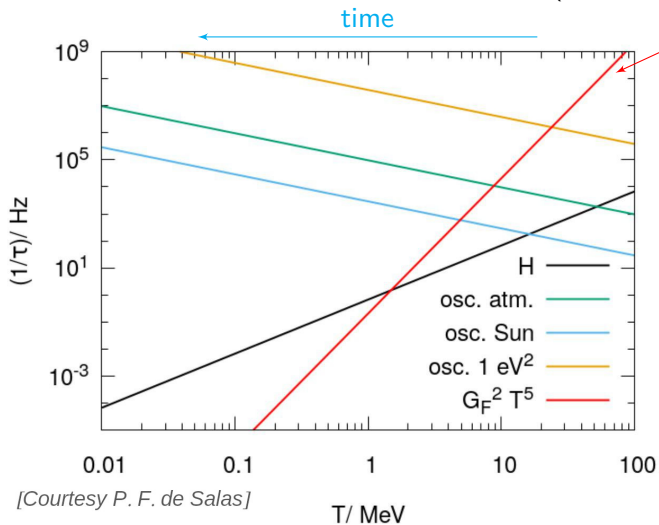


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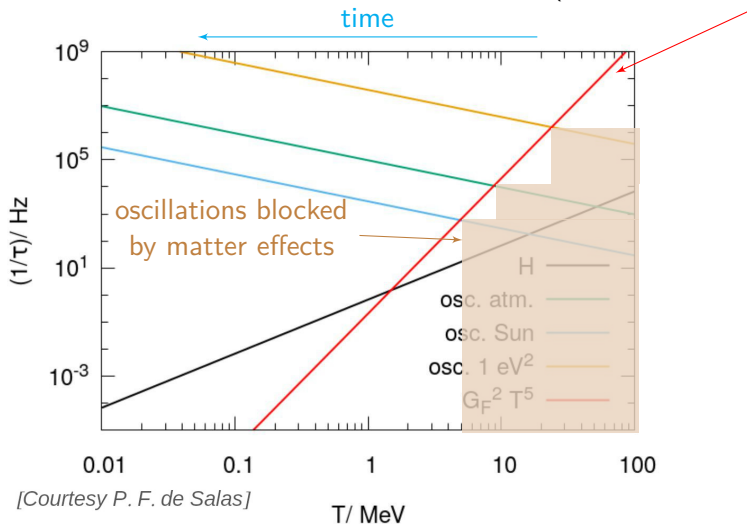
# Neutrinos in the early Universe

before BBN: neutrinos coupled to plasma ( $\nu_\alpha \bar{\nu}_\alpha \leftrightarrow e^+ e^-$ ,  $\nu e \leftrightarrow \nu e$ )



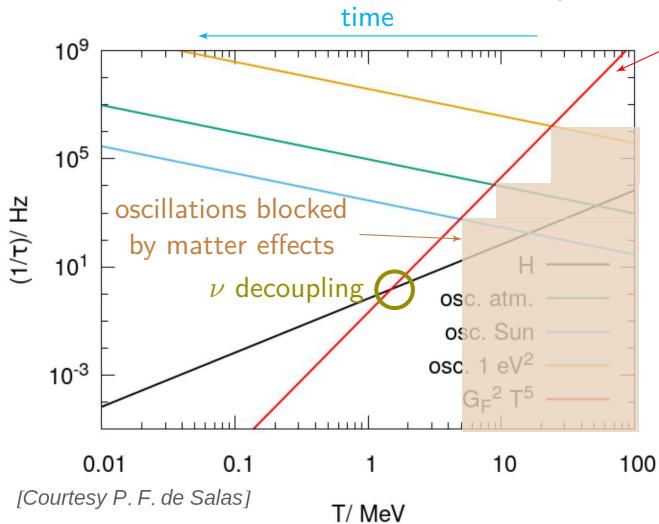
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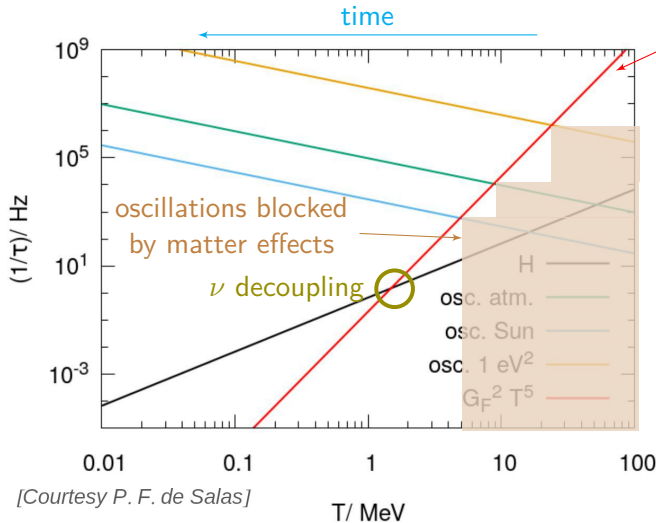
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$\nu$  decouple mostly before  $e^+ e^- \rightarrow \gamma\gamma$  annihilation!

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$$T_\nu \simeq (4/11)^{1/3} T_\gamma$$

after  $e^+ e^- \rightarrow \gamma\gamma$

$f_\nu$ : frozen Fermi-Dirac distribution

Today:

$$T_{\nu,0} = 1.945 \text{ K} \simeq 1.676 \times 10^{-4} \text{ eV}$$

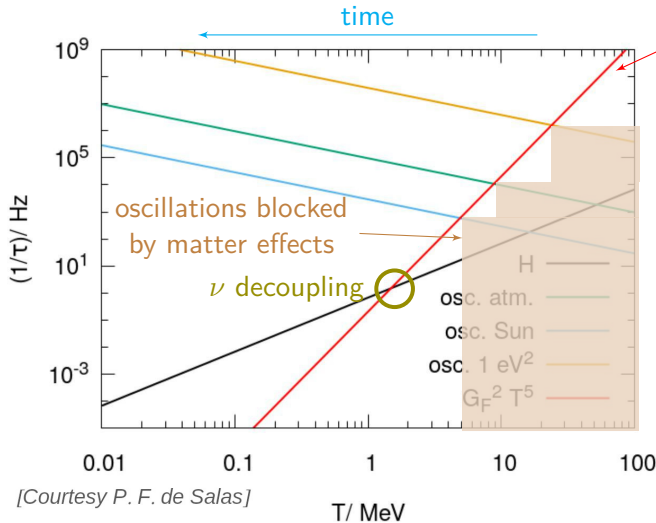
$$\langle E_\nu \rangle \simeq 3.1 T_{\nu,0} \simeq 5 \times 10^{-4} \text{ eV}$$

$$n_0 = n_{\nu,0} = n_{\bar{\nu},0} \simeq 56 \text{ cm}^{-3} \text{ per family}$$

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$\nu$  decouple mostly before  $e^+ e^- \rightarrow \gamma\gamma$  annihilation!  
 actually, the decoupling  $T$  is momentum dependent!

distortions to equilibrium  $f_\nu$ !

# $\nu$ oscillations in the early universe

[Bennett, SG+, JCAP 2021]  
[Sigl, Raffelt, 1993]

comoving coordinates:  $a = 1/T$   $x \equiv m_e a$   $y \equiv p a$   $z \equiv T_\gamma a$   $w \equiv T_\nu a$

density matrix:  $\varrho(x, y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} \end{pmatrix}$   
 $\propto \langle a_j^\dagger(p, t) a_i(p, t) \rangle$   
 off-diagonals to take into account coherency in the neutrino system

$$\varrho \text{ evolution from } x \text{ to } y: \frac{d\varrho(y, x)}{dx} = -ia[\mathcal{H}_{\text{eff}}, \varrho] + b\mathcal{I}$$

$H$  Hubble factor  $\rightarrow$  expansion (depends on universe content)

effective Hamiltonian  $\mathcal{H}_{\text{eff}} = \frac{M_F}{2y} - \frac{2\sqrt{2}G_F y m_e^6}{x^6} \left( \frac{E_\ell + P_\ell}{m_W^2} + \frac{4}{3} \frac{E_\nu}{m_Z^2} \right)$

vacuum oscillations  $\longleftarrow$   $\longrightarrow$  matter effects

$\mathcal{I}$  collision integrals

take into account  $\nu$ -e scattering and pair annihilation,  $\nu$ - $\nu$  interactions

2D integrals over momentum, take most of the computation time

$$\text{solve together with } z \text{ evolution, from } x \frac{d\rho(x)}{dx} = \rho - 3P$$

$\rho, P$  total energy density and pressure, also take into account FTQED corrections



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FORTran-Evolved Primordial Neutrino Oscillations  
(FortEPiano)

[https://bitbucket.org/ahep\\_cosmo/fortepiano\\_public](https://bitbucket.org/ahep_cosmo/fortepiano_public)

vacuum oscillations

matter effects

$\mathcal{I}$  collision integrals

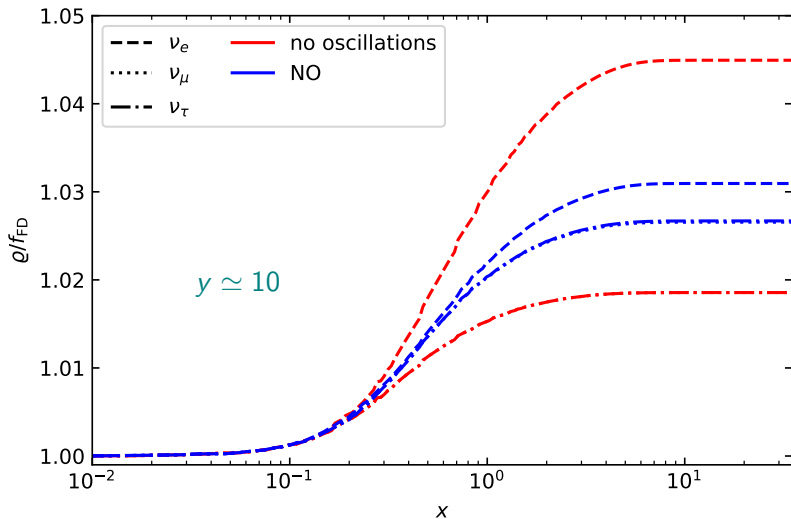
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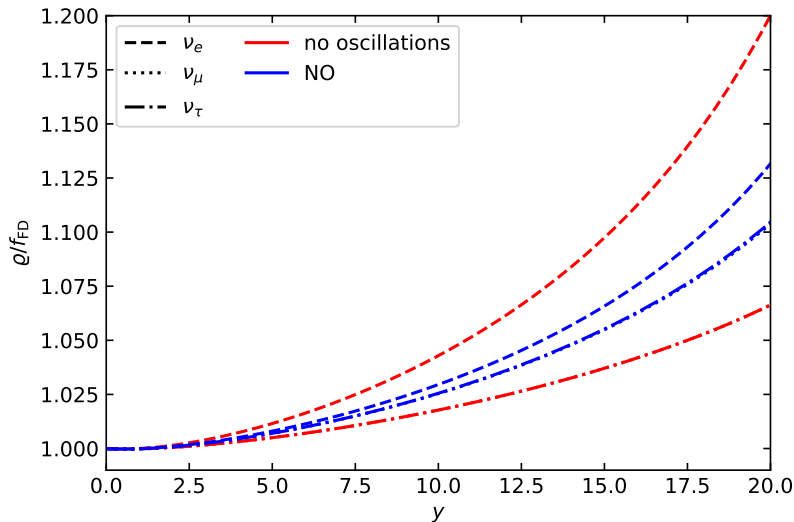
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Distortion of the momentum distribution ( $f_{\text{FD}}$ : Fermi-Dirac at equilibrium)

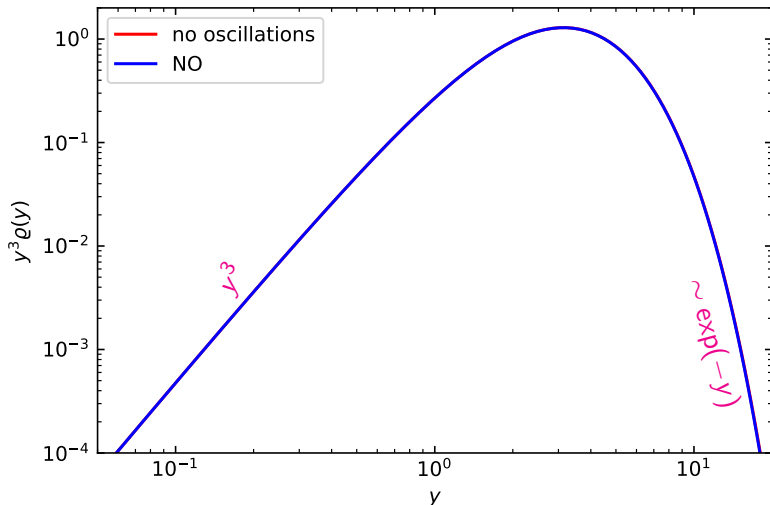


Distortion of the momentum distribution ( $f_{\text{FD}}$ : Fermi-Dirac at equilibrium)

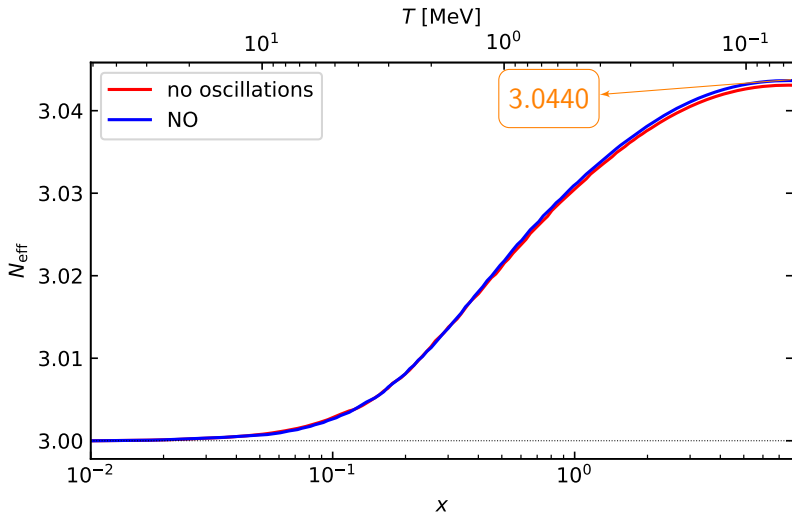


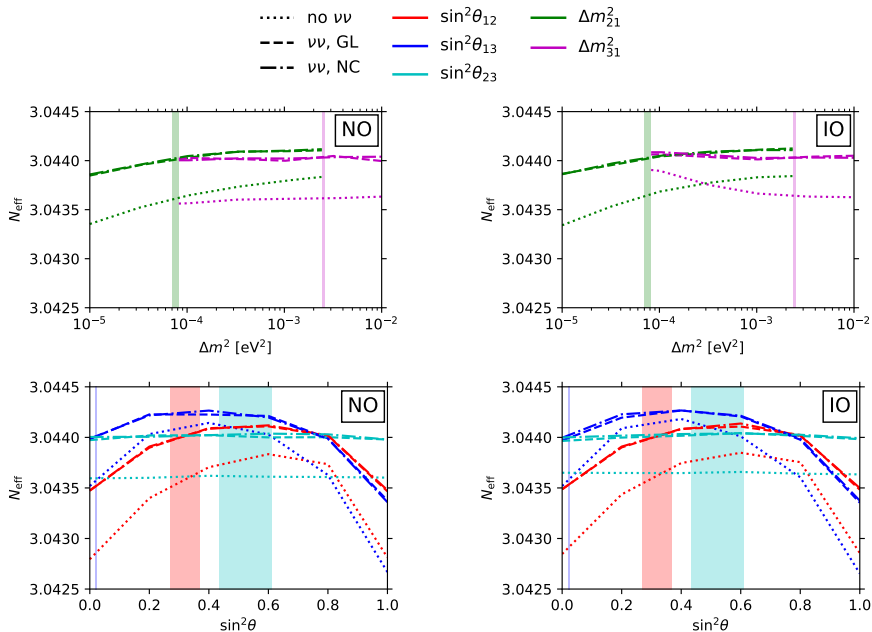
$$N_{\text{eff}}^{\text{final}} = \frac{8}{7} \left( \frac{11}{4} \right)^{4/3} \frac{\rho_\nu}{\rho_\gamma} = \frac{8}{7} \left( \frac{11}{4} \right)^{4/3} \frac{1}{\rho_\gamma} \sum_i g_i \int \frac{d^3 p}{(2\pi)^3} E(p) f_{\nu,i}(p)$$

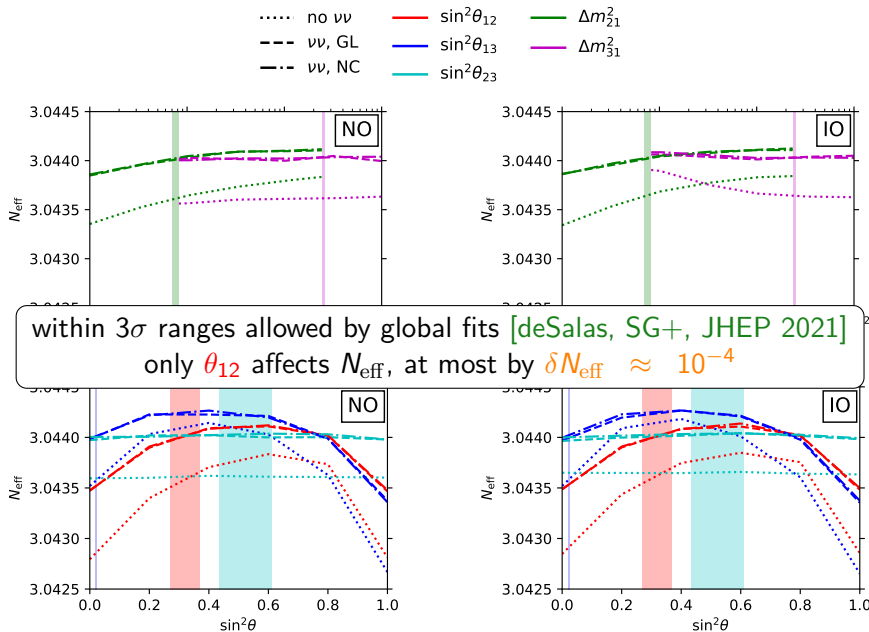
$(11/4)^{1/3} = (T_\gamma/T_\nu)^{\text{fin}}$ 
 $\hookrightarrow \propto y^3 \varrho_{ii}(y)$



$$N_{\text{eff}}^{\text{any time}} = \frac{8}{7} \left( \frac{T_\gamma}{T_\nu} \right)^4 \frac{\rho_\nu}{\rho_\gamma} = \frac{8}{7} \left( \frac{T_\gamma}{T_\nu} \right)^4 \frac{1}{\rho_\gamma} \sum_i g_i \int \frac{d^3 p}{(2\pi)^3} E(p) f_{\nu,i}(p)$$

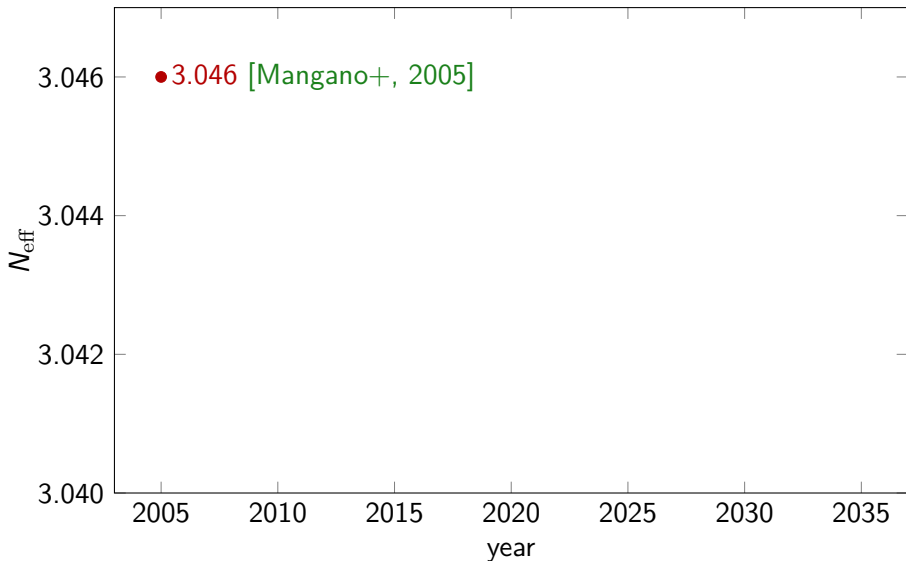






## How precise is $N_{\text{eff}} = 3.04\dots$ ?

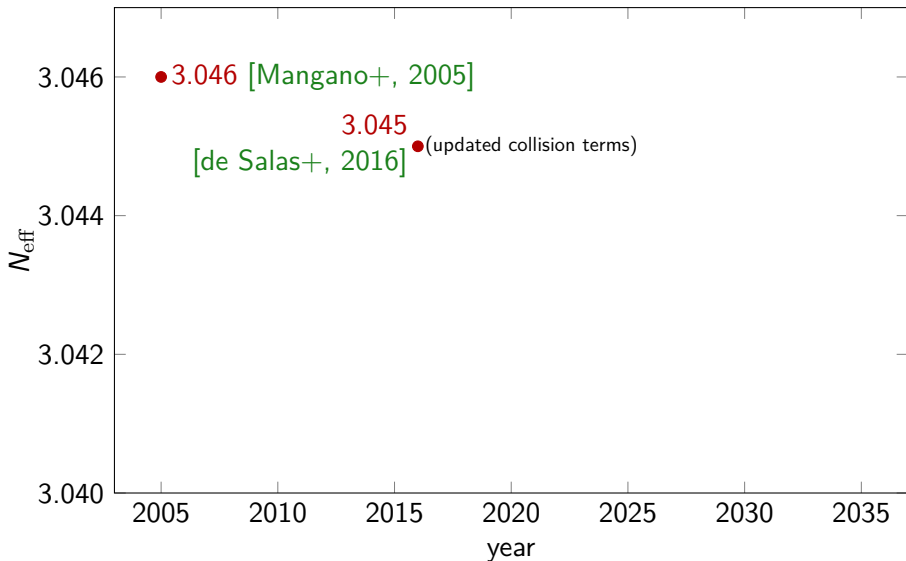
Full  $3\nu$  mixing results:





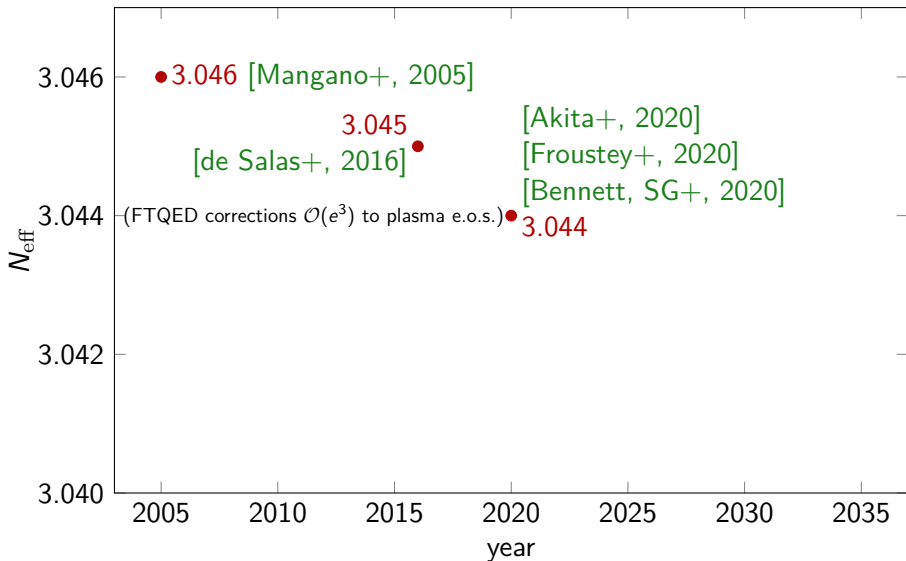
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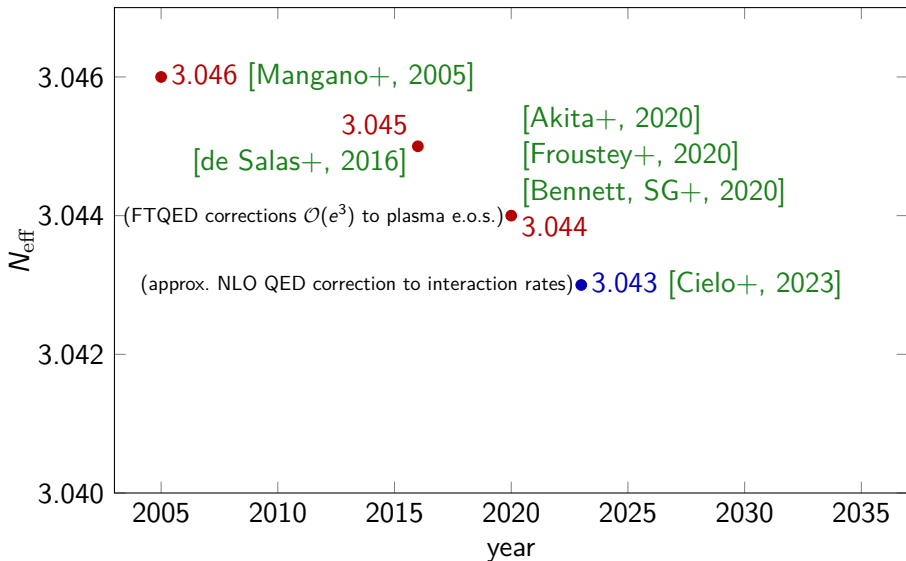
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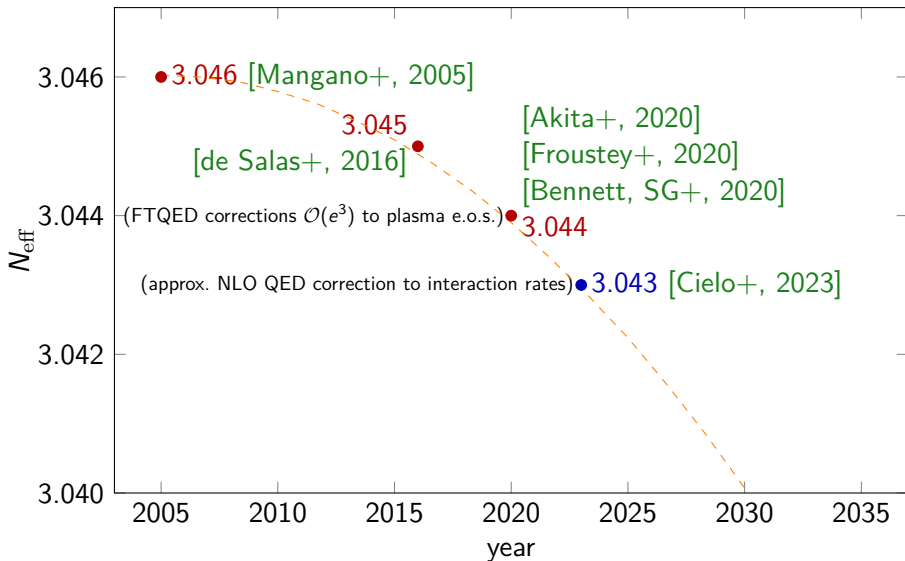
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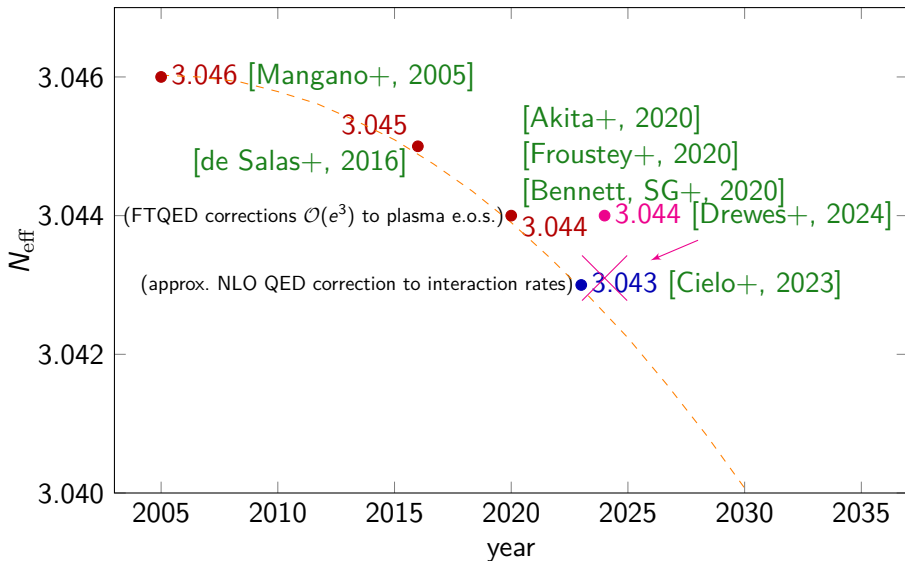
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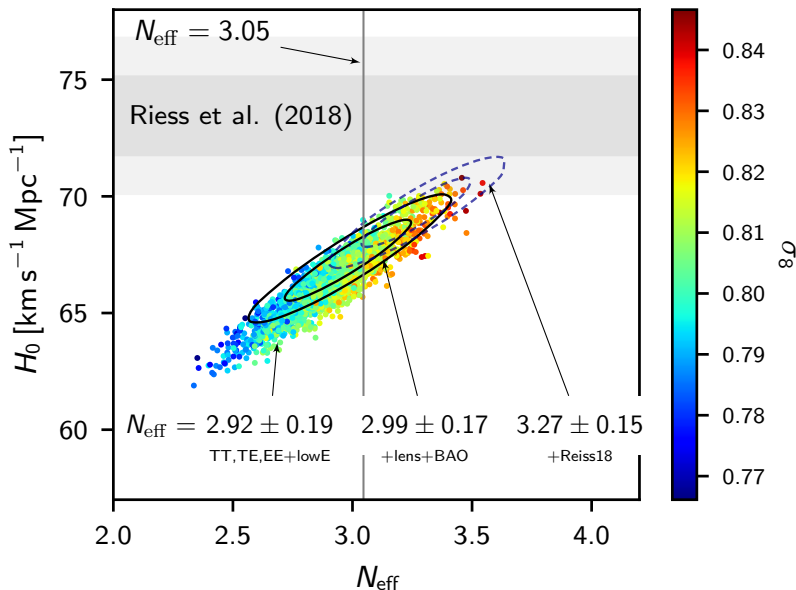
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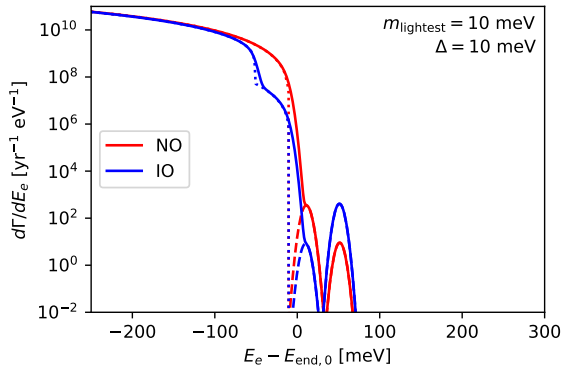


## D

## Direct detection of relic neutrinos

Based on:

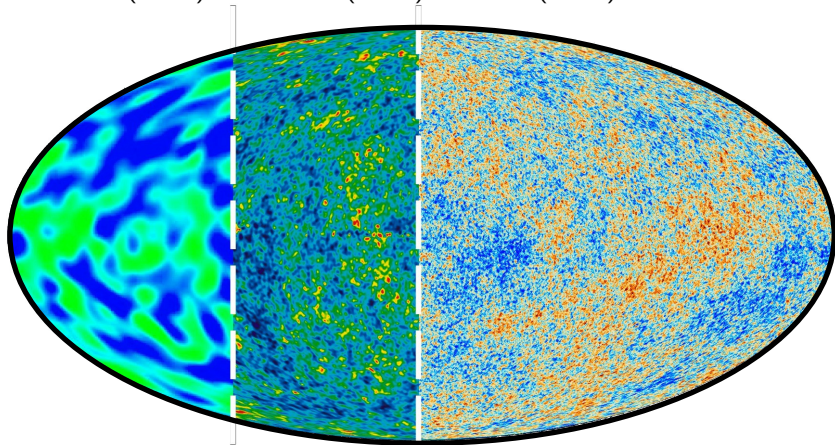
- JCAP 01 (2023) 003
- JCAP 08 (2014) 038
- JCAP 07 (2019) 047



# The oldest picture of the Universe

The Cosmic Microwave Background, generated at  $t \simeq 4 \times 10^5$  years

COBE (1992)    WMAP (2003)    Planck (2013)

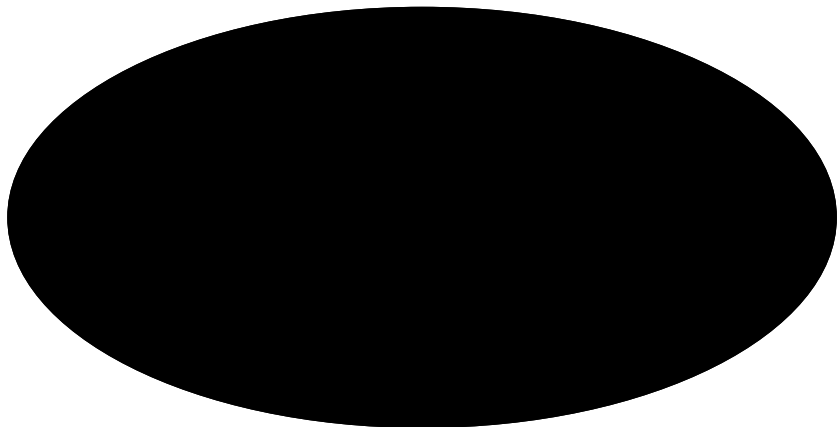




# The oldest picture of the Universe

The Cosmic Neutrino Background, generated at  $t \simeq 1$  s

...  $\rightarrow$  2024  $\rightarrow$  ...



$T_\nu \sim 10^{-4}$  eV,  $E_\nu \sim 5 \times 10^{-4}$  eV today!

We need **thresholdless detection process...** **How do we get them?**

# Stodolsky effect?

How to directly detect non-relativistic neutrinos?

Stodolsky effect

[Stodolsky, 1974][Duda+, 2001]

(only if there is  
lepton asymmetry)

energy splitting of  $e^-$  spin states due to  
coherent scattering with relic neutrinos



torque on  $e^-$  in lab rest frame



use a ferromagnet to build detector



measure torque with a torsion balance

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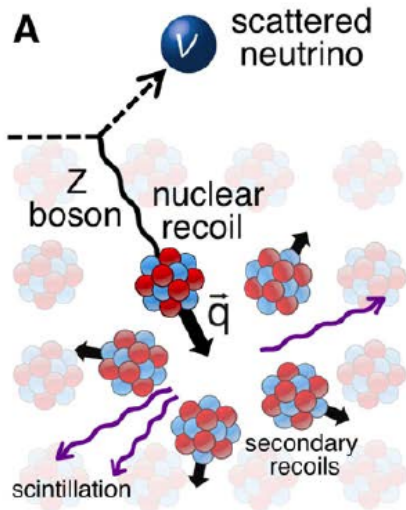
$$\text{expected } a_\nu \simeq \mathcal{O}(10^{-26}) \text{ cm/s}^2$$



$$a_{\text{exp}} \simeq \mathcal{O}(10^{-12}) \text{ cm/s}^2$$

First of all: what's **Coherent Elastic  $\nu$ -Nucleous Scattering**?

**elastic scattering** where  $\nu$  interacts with **nucleous** "as a whole"



Predicted for  $|\vec{q}|R \lesssim 1$   
by [Freedman, PRD 1974]

**small recoil energies!**  $\lesssim 10$  keV...  
difficult to measure

$$\frac{d\sigma}{dT}(E_\nu, T) \sim \frac{G_F^2 M}{4\pi} N^2$$

[Drukier, Stodolsky, PRD 1984]

**enhancement  $N^2$**  because  
 $\nu$  interacts  
**coherently with all nucleons**

may give huge cross  
section enhancement

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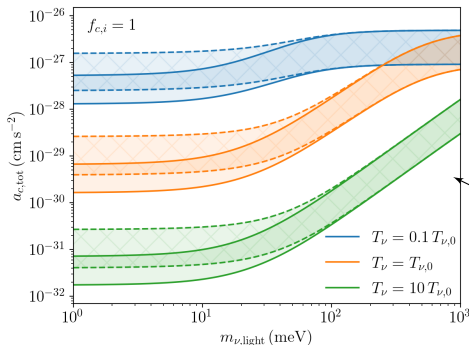
elastic scattering where  $\nu$  interacts with nucleous "as a whole"

Can we detect relic neutrinos with CE $\nu$ NS?

relic neutrinos have de Broglie length  $\lambda \sim 2\pi/p_\nu$



enhancement in interactions due to coherence with nuclei in volume  $\lambda^3$



Acceleration induced by CE $\nu$ NS  
of relic  $\nu$  on test mass  $M$ :

$$a^N \propto ((A - Z)/A)^2 E_{\nu} / p_{\nu}^2 \Delta p_{\nu} n_{\nu} \rho$$

$A, Z$  mass, atomic numbers  
 $p_{\nu}, E_{\nu}$  neutrino momentum and energy  
 $\Delta p_{\nu}$  net momentum transfer  
 $n_{\nu}$  neutrino number density  
 $\rho$  target mass density

unclustered relic  $\nu$ s,  $n_{\nu} = n_0$

$a^N$  of atoms in silicon target

First of all: what's **Coherent Elastic  $\nu$ -Nucleous Scattering**?

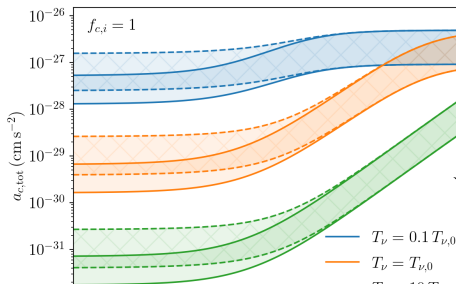
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 $n_\nu$  neutrino number density  
 $\rho$  target mass density

unclustered relic  $\nu$ s.  $n_\nu = n_0$

proposed torsion balances can most optimistically reach  $a \sim 10^{-23} \text{ cm s}^{-2}$

# At interferometers?

How to directly detect non-relativistic neutrinos?

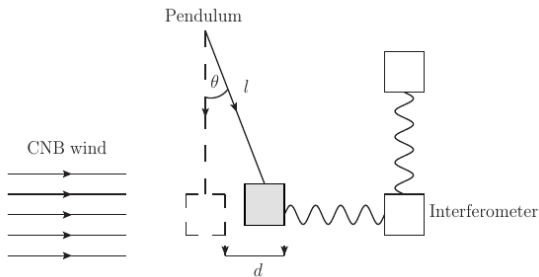
At interferometers

[Domcke+, 2017] [Shergold, 2021]

coherent scattering of  
relic  $\nu$  on a pendulum



measure oscillations  
at interferometers



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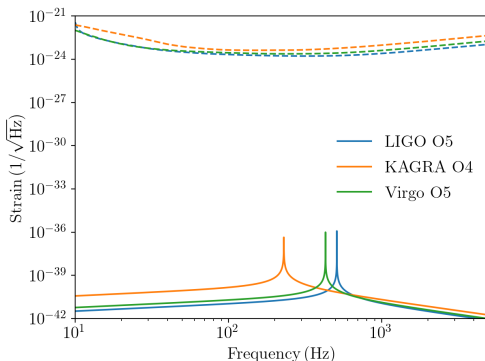
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measure oscillations  
at interferometers



Sensitivity

Example signal  
 $m_\nu = 0.8$  eV



How to directly detect non-relativistic neutrinos?

Remember that  
 $\langle E_\nu \rangle \simeq \mathcal{O}(10^{-4})$  eV today



a process without energy  
 threshold is necessary

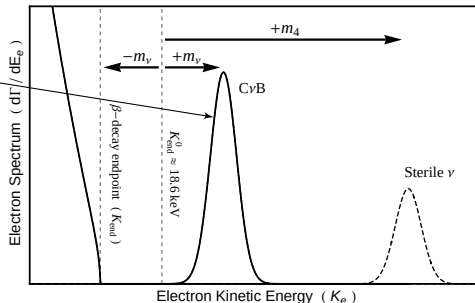
[Weinberg, 1962]: neutrino capture in  $\beta$ -decaying nuclei  $\nu + n \rightarrow p + e^-$

Main background:  $\beta$  decay  $n \rightarrow p + e^- + \bar{\nu}$ !

signal is a peak at  $2m_\nu$   
 above  $\beta$ -decay endpoint

only with a lot of material

need a very good energy resolution



best element has highest  $\sigma_{\text{NCB}}(v_\nu/c) \cdot t_{1/2}$

to minimize contamination from  $\beta$  decay background

| Isotope           | Decay     | $Q_\beta$ (keV) | Half-life (s)           | $\sigma_{\text{NCB}}(v_\nu/c) (10^{-41} \text{ cm}^2)$ |
|-------------------|-----------|-----------------|-------------------------|--------------------------------------------------------|
| $^3\text{H}$      | $\beta^-$ | 18.591          | $3.8878 \times 10^8$    | $7.84 \times 10^{-4}$                                  |
| $^{63}\text{Ni}$  | $\beta^-$ | 66.945          | $3.1588 \times 10^9$    | $1.38 \times 10^{-6}$                                  |
| $^{93}\text{Zr}$  | $\beta^-$ | 60.63           | $4.952 \times 10^{13}$  | $2.39 \times 10^{-10}$                                 |
| $^{106}\text{Ru}$ | $\beta^-$ | 39.4            | $3.2278 \times 10^7$    | $5.88 \times 10^{-4}$                                  |
| $^{107}\text{Pd}$ | $\beta^-$ | 33              | $2.0512 \times 10^{14}$ | $2.58 \times 10^{-10}$                                 |
| $^{187}\text{Re}$ | $\beta^-$ | 2.64            | $1.3727 \times 10^{18}$ | $4.32 \times 10^{-11}$                                 |
| $^{11}\text{C}$   | $\beta^+$ | 960.2           | $1.226 \times 10^3$     | $4.66 \times 10^{-3}$                                  |
| $^{13}\text{N}$   | $\beta^+$ | 1198.5          | $5.99 \times 10^2$      | $5.3 \times 10^{-3}$                                   |
| $^{15}\text{O}$   | $\beta^+$ | 1732            | $1.224 \times 10^2$     | $9.75 \times 10^{-3}$                                  |
| $^{18}\text{F}$   | $\beta^+$ | 633.5           | $6.809 \times 10^3$     | $2.63 \times 10^{-3}$                                  |
| $^{22}\text{Na}$  | $\beta^+$ | 545.6           | $9.07 \times 10^7$      | $3.04 \times 10^{-7}$                                  |
| $^{45}\text{Ti}$  | $\beta^+$ | 1040.4          | $1.307 \times 10^4$     | $3.87 \times 10^{-4}$                                  |

best element has highest  $\sigma_{\text{NCB}}(v_\nu/c) \cdot t_{1/2}$

to minimize contamination from  $\beta$  decay background

| Isotope           | Decay     | $Q_\beta$ (keV) | Half-life (s)           | $\sigma_{\text{NCB}}(v_\nu/c)$ ( $10^{-41} \text{ cm}^2$ ) |
|-------------------|-----------|-----------------|-------------------------|------------------------------------------------------------|
| $^3\text{H}$      | $\beta^-$ | 18.591          | $3.8878 \times 10^8$    | $7.84 \times 10^{-4}$                                      |
| $^{63}\text{Ni}$  | $\beta^-$ | 66.945          | $3.1588 \times 10^9$    | $1.38 \times 10^{-6}$                                      |
| $^{93}\text{Zr}$  | $\beta^-$ | 60.63           | $4.952 \times 10^{13}$  | $2.39 \times 10^{-10}$                                     |
| $^{106}\text{Ru}$ | $\beta^-$ | 39.4            | $3.2278 \times 10^7$    | $5.88 \times 10^{-4}$                                      |
| $^{107}\text{Pd}$ | $\beta^-$ | 33              | $2.0512 \times 10^{14}$ | $2.58 \times 10^{-10}$                                     |
| $^{187}\text{Re}$ | $\beta^-$ | 2.64            | $1.3727 \times 10^{18}$ | $4.32 \times 10^{-11}$                                     |
| $^{11}\text{C}$   | $\beta^+$ | 960.2           | $1.226 \times 10^3$     | $4.66 \times 10^{-3}$                                      |
| $^{13}\text{N}$   | $\beta^+$ | 1198.5          | $5.99 \times 10^2$      | $5.3 \times 10^{-3}$                                       |
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<sup>3</sup>H better because the cross section ( $\rightarrow$  event rate) is higher

$$\frac{d\tilde{\Gamma}_{\text{CNB}}}{dE_e}(E_e) = \frac{1}{\sqrt{2\pi}\sigma} \sum_{i=1}^{N_\nu} \bar{\sigma} N_T |U_{ei}|^2 n_0 f_c(m_i) \times e^{-\frac{[E_e - (E_{\text{end}} + m_i + m_{\text{lightest}})]^2}{2\sigma^2}}$$

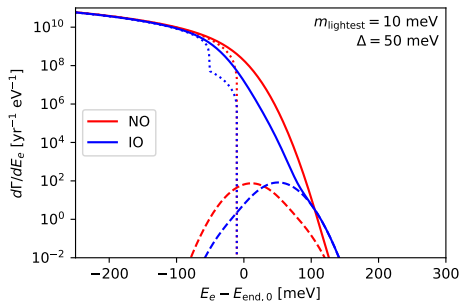
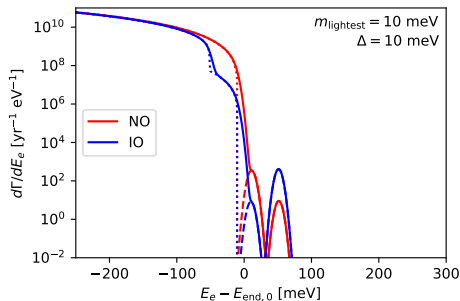
$$\frac{d\Gamma_\beta}{dE_e} = \frac{\bar{\sigma}}{\pi^2} N_T \sum_{i=1}^{N_\nu} |U_{ei}|^2 H(E_e, m_i)$$

$$\frac{d\tilde{\Gamma}_\beta}{dE_e}(E_e) = \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{+\infty} dx \frac{d\Gamma_\beta}{dE_e}(x) \exp\left[-\frac{(E_e - x)^2}{2\sigma^2}\right]$$

$$\frac{d\tilde{\Gamma}_{\text{CNB}}}{dE_e}(E_e) = \frac{1}{\sqrt{2\pi}\sigma} \sum_{i=1}^{N_\nu} \bar{\sigma} N_T |U_{ei}|^2 n_0 f_c(m_i) \times e^{-\frac{[E_e - (E_{\text{end}} + m_i + m_{\text{lightest}})]^2}{2\sigma^2}}$$

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Pontecorvo Tritium Observatory for Light, Early-universe, Massive-neutrino Yield (PTOLEMY)

expected resolution  $\Delta \simeq 0.1 \text{ eV?}$   
 $0.05 \text{ eV?}$

can probe  $m_\nu \simeq 1.4\Delta \simeq 0.1 \text{ eV}$

built mainly for CNB

$M_T = 100 \text{ g of atomic } ^3\text{H}$

$$\Gamma_{\text{CNB}} = \sum_{i=1}^3 |U_{ei}|^2 [n_i(\nu_{h_R}) + n_i(\nu_{h_L})] N_T \bar{\sigma} \sim \mathcal{O}(10) \text{ yr}^{-1}$$

$N_T$  number of  $^3\text{H}$  nuclei in a sample of mass  $M_T$      $\bar{\sigma} \simeq 3.834 \times 10^{-45} \text{ cm}^2$      $n_i$  number density of neutrino  $i$

(without clustering)

# Pontecorvo Tritium Observatory for Light, Early-universe, Massive-neutrino Yield (PTOLEMY)

expected resolution  $\Delta \simeq 0.1 \text{ eV}$ ?  
0.05 eV?

built mainly for CNB

can probe  $m_\nu \simeq 1.4\Delta \simeq 0.1 \text{ eV}$

$M_T = 100 \text{ g}$  of atomic  $^3\text{H}$

enhancement from  
 $\nu$  clustering in the galaxy?

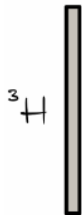
enhancement from  
other effects?

$$\Gamma_{\text{CNB}} = \sum_{i=1}^3 |U_{ei}|^2 [n_i(\nu_{h_R}) + n_i(\nu_{h_L})] N_T \bar{\sigma} \sim \mathcal{O}(10) \text{ yr}^{-1}$$

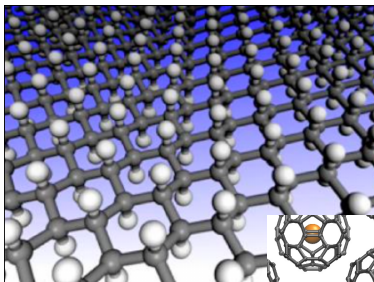
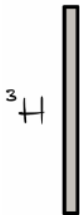
$N_T$  number of  $^3\text{H}$  nuclei in a sample of mass  $M_T$      $\bar{\sigma} \simeq 3.834 \times 10^{-45} \text{ cm}^2$      $n_i$  number density of neutrino  $i$

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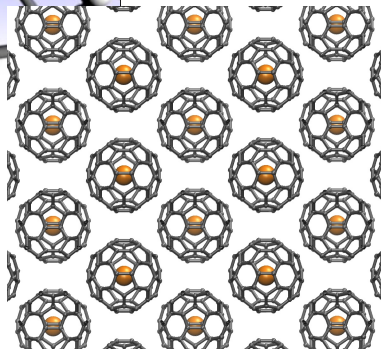


[Courtesy A. Esposito]



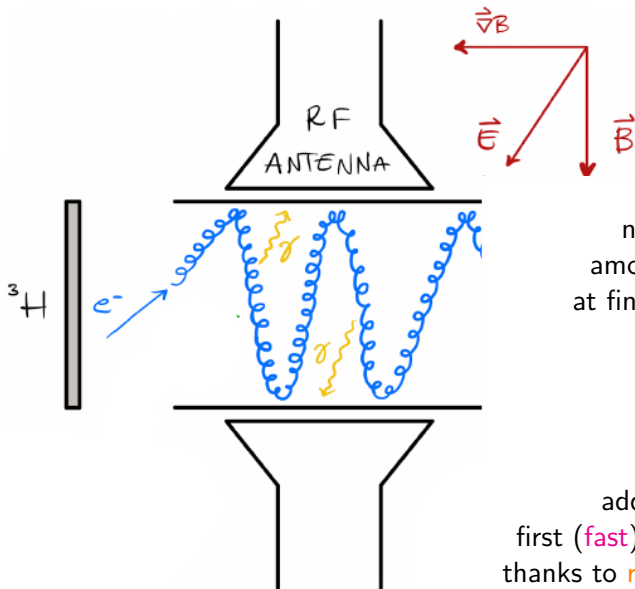
tritium on  
graphene?

tritium on  
fullerene?



[Courtesy A. Esposito]

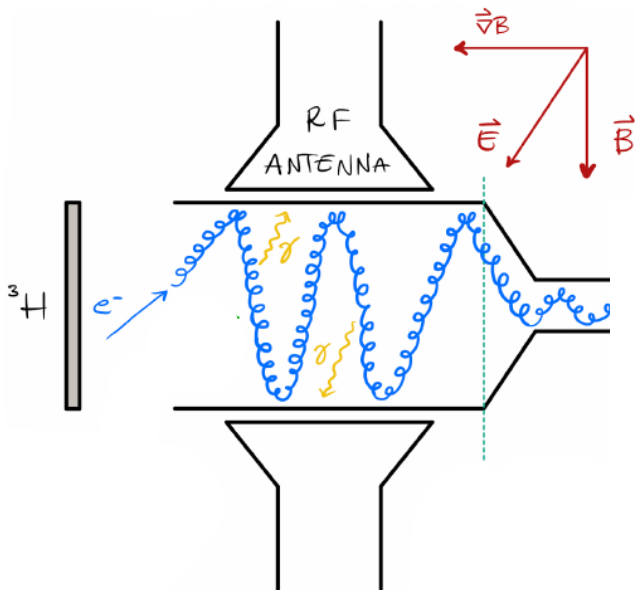
[Courtesy V. Tozzini]



need to reduce  
amount of **electrons**  
at final energy sensors:  
**EM filter**

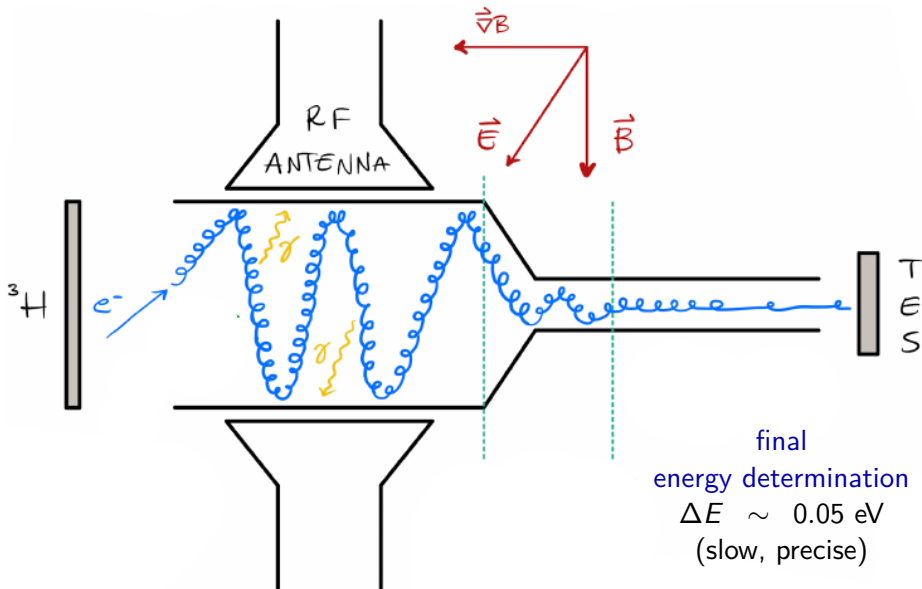
additional benefit:  
first (**fast**) energy determination  
thanks to **radio-frequency antenna**

[Courtesy A. Esposito]



filter only events  
close to endpoint  
( $E \gtrsim E_0 - 10 \text{ eV}$ )

[Courtesy A. Esposito]



[Courtesy A. Esposito]

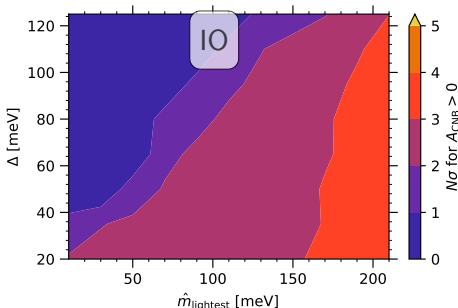
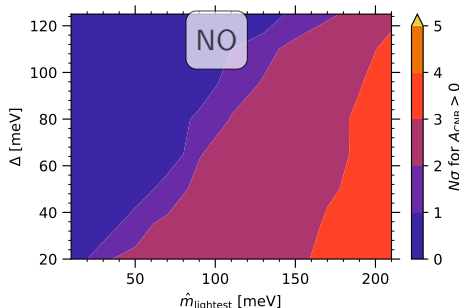
using the definition:

$$N_{\text{th}}^i(\theta) = A_\beta N_\beta^i(\hat{E}_{\text{end}} + \Delta E_{\text{end}}, m_i, U) + \mathbf{A}_{\text{CNB}} N_{\text{CNB}}^i(\hat{E}_{\text{end}} + \Delta E_{\text{end}}, m_i, U) + N_b$$

if  $\mathbf{A}_{\text{CNB}} > 0$  at  $N\sigma$ , direct detection of CNB accomplished at  $N\sigma$

statistical only!

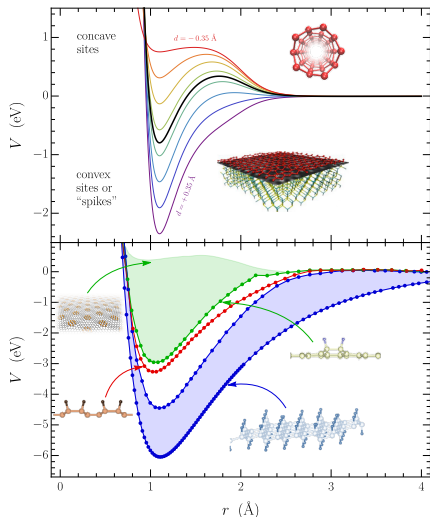
significance on  $A_{\text{CNB}} > 0$   
as a function of  $\hat{m}_{\text{lightest}}$ ,  $\Delta$



Tritium atoms will be attached to graphene sheets



Quantum uncertainty on electron energy due to condensed matter state



Binding potential depends on  
distance graphene-tritium  
and  
graphene sheet configuration

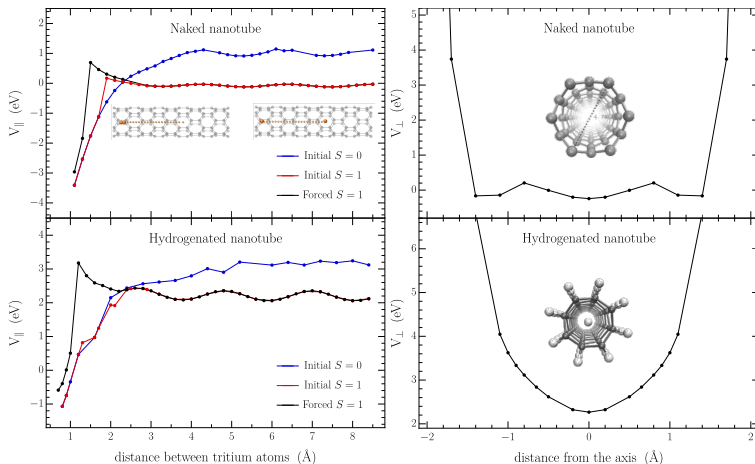
Concave sites  
behave differently than  
convex ones

Hydrogen coverage (amount  
of  $H$  atoms in the graphene  
sheet) is also important

Tritium atoms will be attached to graphene sheets



Quantum uncertainty on electron energy due to condensed matter state



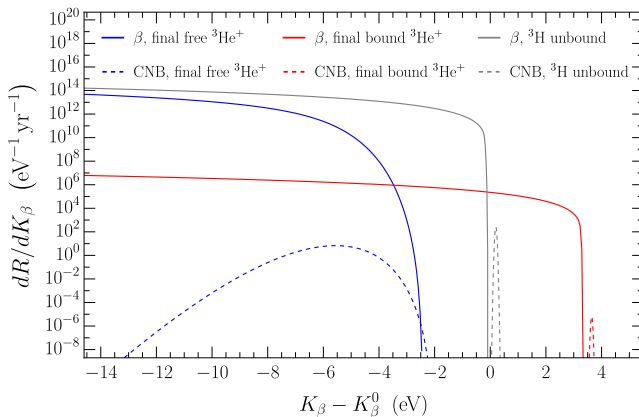
Nanotube configurations have flat potential parallel to nanotube axis



Tritium atoms will be attached to graphene sheets



Quantum uncertainty on electron energy due to condensed matter state



Graphene configuration can make observation range from  
completely impossible or suppressed but possible



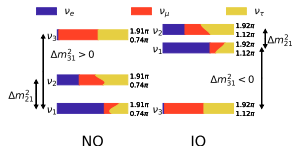
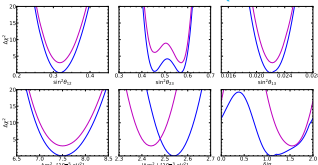
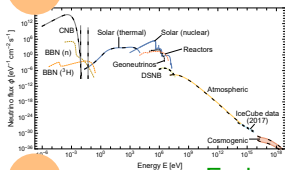
## Conclusions

almost there!

# What do we learn from relic neutrinos?

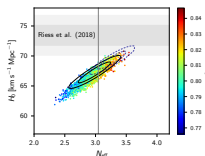
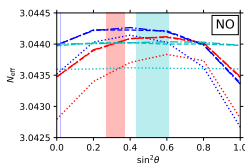
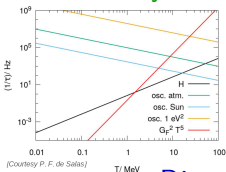
N

## Neutrinos: precision era (many sources)



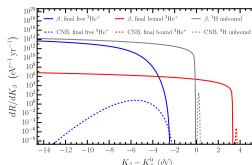
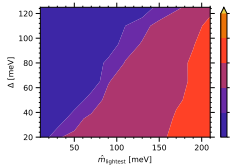
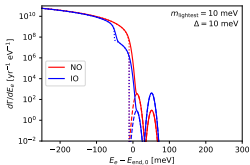
E

## Early universe effects: indirect indications



D

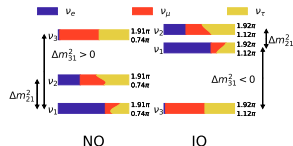
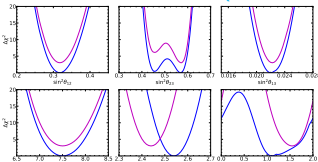
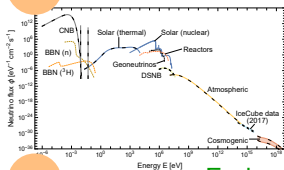
## Direct detection: still long to go



# What do we learn from relic neutrinos?

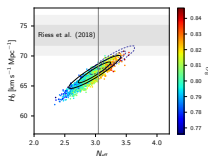
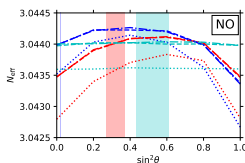
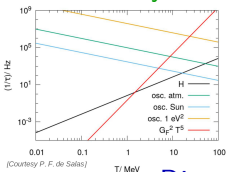
N

## Neutrinos: precision era (many sources)



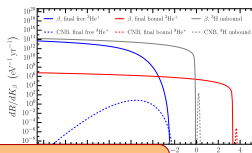
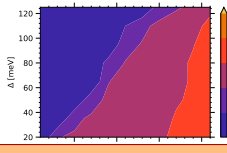
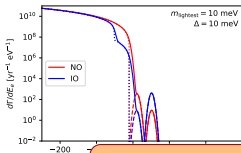
E

## Early universe effects: indirect indications



D

## Direct detection: still long to go



Thanks for your attention!