

Indirect and Direct Detection of Dark Matter and Flavor Symmetry

Takashi Toma

Kanazawa University
Max-Planck Institute for Nuclear Physics

FLASY11@Valencia in Spain
12th July

Based on arXiv:1104.0367 published in Eur. Phys. J. C.
In collaboration with: Yuji Kajiyama and Hiroshi Okada

Outline

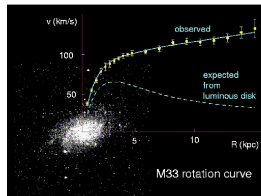
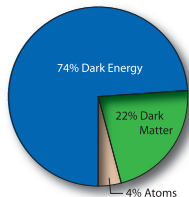
- 1 Introduction
- 2 Indirect and Direct Detection of Dark Matter
 - e^\pm excess in cosmic ray
 - Flavor dependence of e^\pm flux
 - Constraint from gamma ray
 - Direct detection of dark matter
- 3 D_6 Flavor Symmetric Model
 - Model
 - Relic abundance of DM
 - Cosmic ray analysis (Indirect detection)
 - Elastic cross section (Direct detection)
- 4 Summary

Introduction

There are many experimental evidences for dark matter (DM).

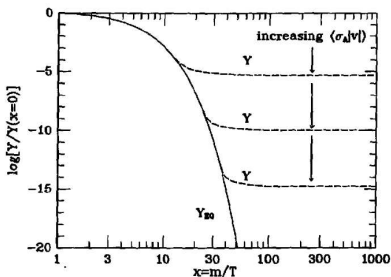
- rotation curves of spiral galaxy
- CMB observation by WMAP
- gravitational lensing
- large scale structure of the universe

The existence of DM is crucial.



The Nature of DM

- Zero electric charge
- Non-relativistic particle
- Stable or extremely long lifetime



$\chi\chi \rightarrow \text{SM particles}$

$$\Omega h^2 \simeq 0.11 \times \left(\frac{3.0 \times 10^{-26}}{\langle\sigma v\rangle [\text{cm}^3/\text{s}]} \right)$$

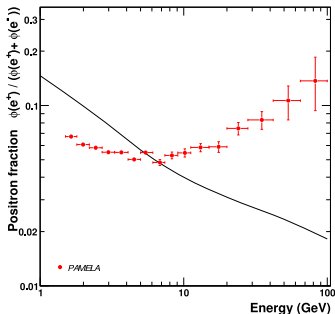
Required annihilation cross section

$$\begin{aligned} \langle\sigma v\rangle &\sim 3 \times 10^{-26} [\text{cm}^3/\text{s}] \\ &\sim 10^{-9} [\text{GeV}^{-2}] \end{aligned}$$

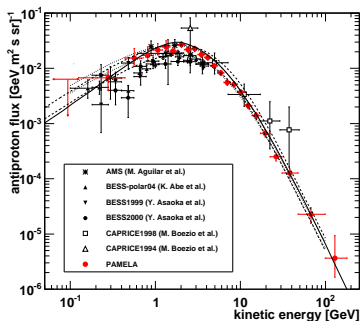
Indirect detection of DM ($DM, DM \rightarrow SM, SM$)
 DM annihilates to SM particles at the present universe.
 The produced particles are observed as the cosmic ray
 (e^\pm, γ, p, ν etc).

PAMELA O. Adriani et al. arXiv:0810.4995, arXiv:1007.0821

e^\pm observation



anti-proton observation



The positron excess might be produced by annihilation of DM. Leptophilic DM is favorable to explain the observation of the positron excess and no excess of anti-proton.

However large annihilation cross section is required.

$$\langle\sigma v\rangle\sim 10^{-7}[\text{GeV}^{-2}] \text{ for } e^\pm \text{ excess of PAMELA}$$

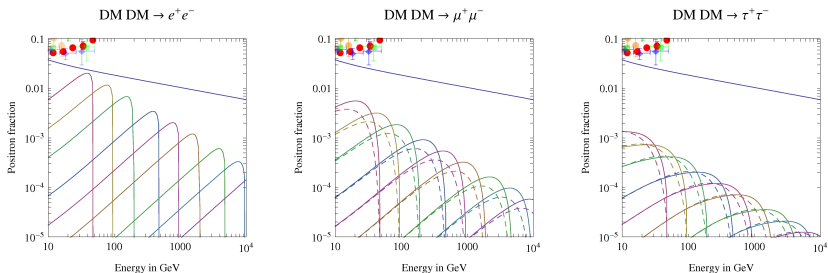
$$\langle\sigma v\rangle\sim 10^{-9}[\text{GeV}^{-2}] \text{ for correct relic abundance}$$

$\mathcal{O}(100)$ difference

How do you explain the difference?

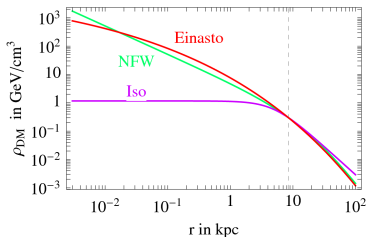
- Breit-Wigner enhancement
- Sommerfeld enhancement
- Non-thermal production of DM
- (Decaying DM)

Flavor dependence on e^\pm flux M. Cirelli et al., arXiv:0809.2409

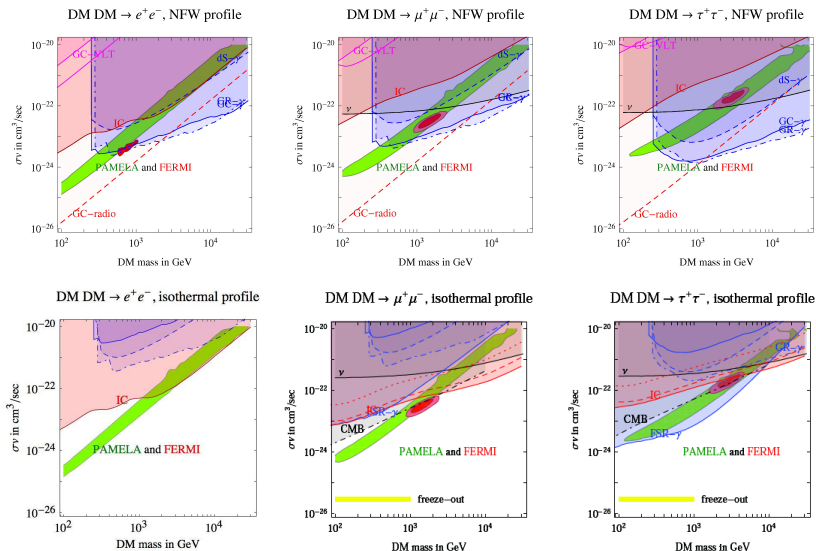


There is also DM density profile dependence for gamma ray flux.

M. Cirelli and P. Panci, arXiv:0904.3830



Constraint from gamma ray P. Meade et al., arXiv:0905.0480; M. Papucci et al., 0912.0742

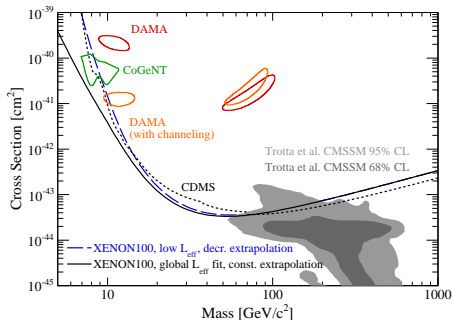


Direct Detection

Elastic scattering between DM and nucleus $\chi N \rightarrow \chi N$
 detection rate

$$\frac{dR}{dE_R} = \frac{\rho_0}{m_\chi m_N} \int v f(v) \frac{d\sigma}{dE_R} dv$$

$f(v)$: velocity distribution of DM, E_R : recoil energy



ρ_0 : local DM density
 $\simeq 0.3 [\text{GeV}/\text{cm}^3]$

Upper bound of
 elastic cross section
 $\sigma_0^{\text{SI}} \lesssim 3.4 \times 10^{-44} [\text{cm}^2]$

XENON100 Collaboration arXiv:1005.0380

D_6 Flavor Symmetric Model

- Field contents

SM + n_i , η and φ

$D_6 \times \hat{Z}_2 \times Z_2$ symmetry is imposed.

	L_S	n_S	e_S^c	L_I	n_I	e_I^c
$SU(2)_L \times U(1)_Y$	$(\mathbf{2}, -1/2)$	$(\mathbf{1}, 0)$	$(\mathbf{1}, 1)$	$(\mathbf{2}, -1/2)$	$(\mathbf{1}, 0)$	$(\mathbf{1}, 1)$
D_6	$\mathbf{1}$	$\mathbf{1}'''$	$\mathbf{1}$	$\mathbf{2}'$	$\mathbf{2}'$	$\mathbf{2}'$
\hat{Z}_2	+	+	-	+	+	-
Z_2	+	-	+	+	-	+

	ϕ_S	ϕ_I	η_S	η_I	φ
$SU(2)_L \times U(1)_Y$	$(\mathbf{2}, -1/2)$	$(\mathbf{2}, -1/2)$	$(\mathbf{2}, -1/2)$	$(\mathbf{2}, -1/2)$	$(\mathbf{1}, 0)$
D_6	$\mathbf{1}$	$\mathbf{2}'$	$\mathbf{1}'''$	$\mathbf{2}'$	$\mathbf{1}$
\hat{Z}_2	+	-	+	+	+
Z_2	+	+	-	-	+

- $D_6 \rightarrow$ predict the lepton mixing, **restrict annihilation channels of DM**
- $\hat{Z}_2 \rightarrow$ suppress FCNC of the quark sector
- $Z_2 \rightarrow$ forbid Dirac neutrino masses and stabilize DM candidate

$$\begin{aligned}
 \mathcal{L}_{\text{lepton}} = & \sum_{a,b,d=1,2,S} \left[Y_{ab}^{ed} (L_a i\sigma_2 \phi_d) e_b^c + Y_{ab}^{\nu d} (\eta_d^\dagger L_a) n_b \right] \\
 & - \sum_{I=1,2} \frac{M_I}{2} n_I n_I - \frac{M_S}{2} n_S n_S \\
 & - \sum_{I=1,2} \frac{\mathfrak{G}_I}{2} \varphi n_I n_I - \frac{\mathfrak{G}_S}{2} \varphi n_S n_S + \text{h.c.}
 \end{aligned}$$

In the model, neutrino masses are generated radiatively.

$$m_\nu \sim \frac{Y^\nu Y^\nu \kappa}{(4\pi)^2} \frac{v^2}{M_S} I_1 \left(\frac{M_S^2}{M_\eta^2} \right)$$
$$\kappa \ll 1, \quad Y^\nu \sim 1, \quad M_S, M_\eta \sim 1 \text{ [TeV]}.$$

MNS matrix at the leading order,

$$V_{MNS} \simeq \begin{pmatrix} \cos \theta_{12} & \sin \theta_{12} & 0 \\ -\frac{1}{\sqrt{2}} \sin \theta_{12} & \frac{1}{\sqrt{2}} \cos \theta_{12} & -\frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \sin \theta_{12} & \frac{1}{\sqrt{2}} \cos \theta_{12} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

- The maximal mixing of atmospheric neutrino is derived.
- Inverted hierarchy is only allowed. ($|\Delta m_{21}^2| < |\Delta m_{23}^2|$)

DM candidates : Z_2 odd particles

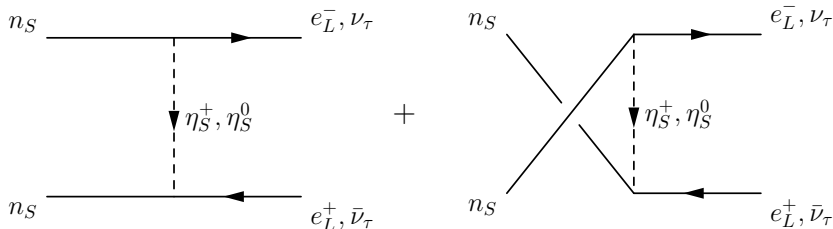
$\rightarrow \eta_I^0, \eta_S^0, n_I, n_S$

We assume n_S to be DM. It is interesting because there are only a few parameters due to D_6 flavor symmetry.

We have to study whether n_S can satisfy the correct relic abundance.

$$\begin{array}{l}
 n_1, n_2, n_S \\
 \mathcal{L} \supset \eta_S^+ \bar{\ell}_i Y_{ij} n_j, \quad Y \simeq \begin{pmatrix} 0 & 0 & h \\ 0 & 0 & \frac{m_e}{m_\mu} h \\ 0 & 0 & 0 \end{pmatrix} \begin{array}{l} e \\ \mu \\ \tau \end{array} \quad \text{for charged leptons} \\
 \mathcal{L} \supset \eta_S^0 \bar{\ell}_i Y_{ij} n_j, \quad Y = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & h \end{pmatrix} \begin{array}{l} \nu_e \\ \nu_\mu \\ \nu_\tau \end{array} \quad \text{for neutrinos}
 \end{array}$$

The annihilation processes of n_S



Only e^\pm are generated as charged leptons because of the D_6 flavor symmetry.

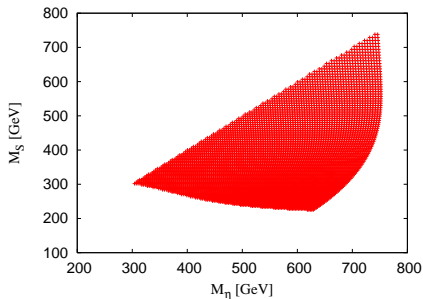
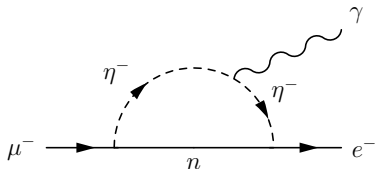
$$\Omega h^2 \propto \frac{1}{\langle \sigma v \rangle}, \quad \langle \sigma v \rangle \simeq \frac{|h|^4}{4\pi} \frac{M_S^2 (M_S^4 + M_\eta^4)}{(M_S^2 + M_\eta^2)^4} \frac{T}{M_S}$$

parameters : h, M_S, M_η

Constraint from $\mu \rightarrow e\gamma$

$$\text{Br}(\mu \rightarrow e\gamma) \simeq \frac{3\alpha_{\text{em}}}{64\pi G_F^2} \frac{m_e^2}{m_\mu^2} \frac{|h|^4}{M_\eta^4} F_2 \left(\frac{M_S^2}{M_\eta^2} \right)$$

$$F_2(x) = \frac{1 - 6x + 3x^2 + 2x^3 - 6x^2 \log x}{6(1-x)^4}$$

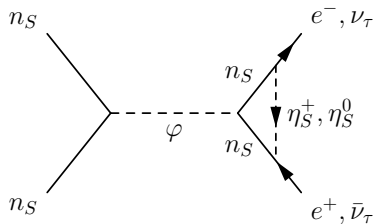


← The allowed region from the following constraints

- $\text{Br}(\mu \rightarrow e\gamma) < 1.2 \times 10^{-11}$
- $|h| < 1.5$
- the relic abundance of n_S ($\Omega h^2 \simeq 0.11$)
- $M_S < M_\eta$.

→ $230\text{GeV} < M_S < 750\text{GeV}$

Indirect Detection (e^\pm excess)



The singlet φ mixes with the Higgses ϕ_I and ϕ_S .

$$\sigma_{2v} \propto \frac{1}{(s - M_R^2)^2 + M_R^2 \Gamma_R^2},$$

$$s \simeq M_S^2 \left(1 + \frac{v^2}{4} \right),$$

R is the resonance particle $2M_S \simeq M_R$.

$$R = \mathcal{O}_I \phi_I + \mathcal{O}_S \phi_S + \mathcal{O}_\varphi \varphi \text{ (mass eigenstate)}$$

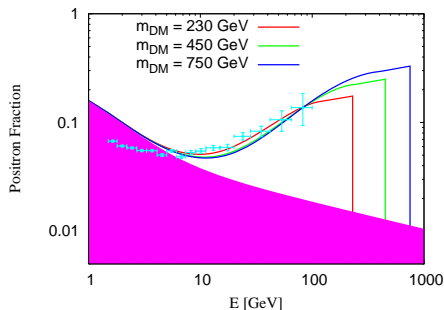
This process is enhanced by the Breit-Wigner enhancement, and effective at only the present universe, and neglected at the early universe.

Indirect Detection (e^\pm excess)

$\langle\sigma_2v\rangle \sim 10^{-7}[\text{GeV}^{-2}]$ is needed.

$$BF \equiv \frac{\langle\sigma_2v\rangle}{3.0 \times 10^{-9}[\text{GeV}^{-2}]} \sim 100$$

DM profile is assumed to be Isothermal.

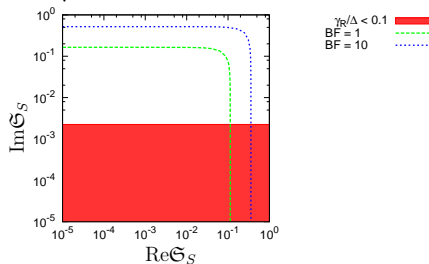


$$\langle\sigma_2v\rangle = \frac{\int(\sigma_2v)e^{-E_1/T}e^{-E_2/T}d^3p_1d^3p_2}{\int e^{-E_1/T}e^{-E_2/T}d^3p_1d^3p_2}$$

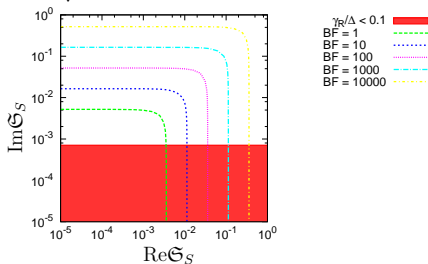
If $\Gamma_R/M_R \equiv \gamma_R \ll \Delta (\equiv 1 - 4M_S^2/M_R^2)$,

$$\langle\sigma_2v\rangle \simeq \frac{\sqrt{\pi}}{10(4\pi)^4}|h|^4\mathcal{O}_\varphi^2(\text{Re}\mathfrak{S}_S)^2\frac{m_e^2}{M_\eta^4}\left(\frac{M_S}{T}\right)^{3/2}e^{-\Delta M_S/T}$$

for $\sqrt{\Delta} = 10^{-6}$



for $\sqrt{\Delta} = 10^{-7}$



When $\text{Re}\mathcal{G}_S \ll \text{Im}\mathcal{G}_S$, Large boost factor is obtained.

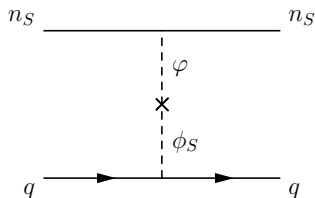
The constraint from no excess of anti-proton flux

Due to the constraint from no anti-proton excess,

$$\langle \sigma_3 v \rangle / \langle \sigma_2 v \rangle \lesssim 10^{-2} \text{ is needed.}$$

$$\longrightarrow \frac{|\mathcal{O}_S|}{|\mathcal{G}_S|} \lesssim 10^{-12} \quad \text{It is the severe constraint!}$$

Direct Detection



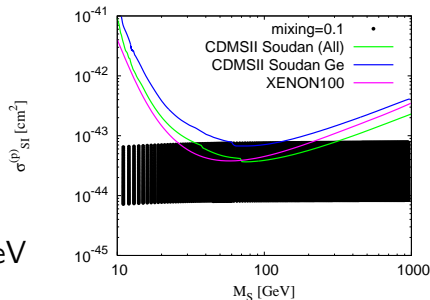
This process occurs through the Higgs mixing.
 The largest contribution comes from the t-channel exchange of SM-Higgs.

$$\text{SM-Higgs} = \mathcal{U}_I \phi_I + \mathcal{U}_S \phi_S + \mathcal{U}_\varphi \varphi$$

$$\sigma_{\text{SI}}^p \propto \left| \frac{\mathcal{U}_S \mathcal{U}_\varphi \mathfrak{G}_S Y_q}{m_{\text{SM-Higgs}}^2} \right|^2$$

$$\mathcal{U}_S \mathcal{U}_\varphi \mathfrak{G}_S = 0.1$$

$$115 \text{ GeV} < m_{\text{SM-Higgs}} < 200 \text{ GeV}$$



Summary

- 1 Leptophilic DM is favored from indirect detection.
- 2 e^\pm flux has flavor dependence, in particular annihilation to τ^\pm is excluded by the constraint from gamma ray.
- 3 DM n_S in the D_6 symmetric model dominantly couples to e^\pm .
- 4 The e^\pm excess in the cosmic ray is explained by the Breit-Wigner enhancement.
- 5 The elastic scattering between DM and quarks occurs through the Higgs mixing.
- 6 It is consistent with indirect and direct detection of DM since the mediate particle for indirect detection is different from direct detection.