

The achievements of the CERN proton – antiproton collider

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- **Motivation of the project**
- **The proton – antiproton collider**
- **UA1 and UA2 detectors**
- **Discovery of the W and Z bosons**
- **Hadronic jets at high transverse momentum**
- **First indirect evidence for $\bar{B}^0 - B^0$ mixing**

IMFP 2004, Alicante, 5 March 2004

**1973 Discovery of neutral – current neutrino interactions:
the first experimental evidence in favour of the unified
electro-weak theory;**

first measurement of the weak mixing angle θ_w

➡ first quantitative prediction of the W^\pm and Z mass values:

$$m_W = 60 - 80 \text{ GeV}$$

$$m_Z = 75 - 95 \text{ GeV}$$

too large to be produced by any existing accelerators

**The ideal machine to produce and study the W and Z bosons
in the most convenient experimental conditions: a high-energy e^+e^- collider**

$$e^+e^- \rightarrow Z$$

$$e^+e^- \rightarrow W^+W^-$$

still far in the future in the 1970's (first operation of LEP in 1989)

1976: the shortcut to W and Z production

(presented at the Neutrino 76 conference in Aachen)

PRODUCING MASSIVE NEUTRAL INTERMEDIATE VECTOR BOSONS WITH EXISTING ACCELERATORS*)

C. Rubbia and P. McIntyre

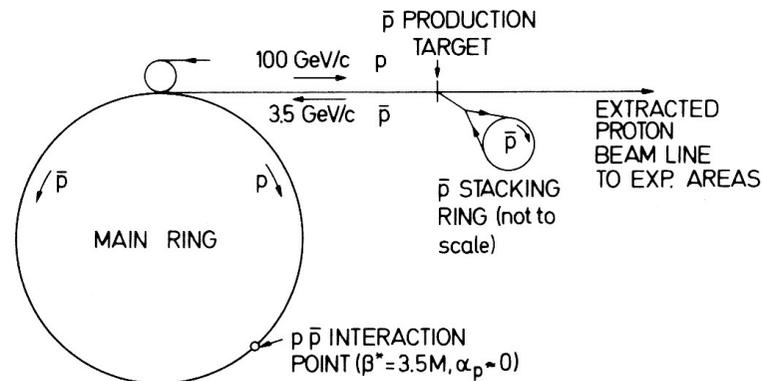
Department of Physics, Harvard University, Cambridge, Massachusetts 02138
and

D. Cline

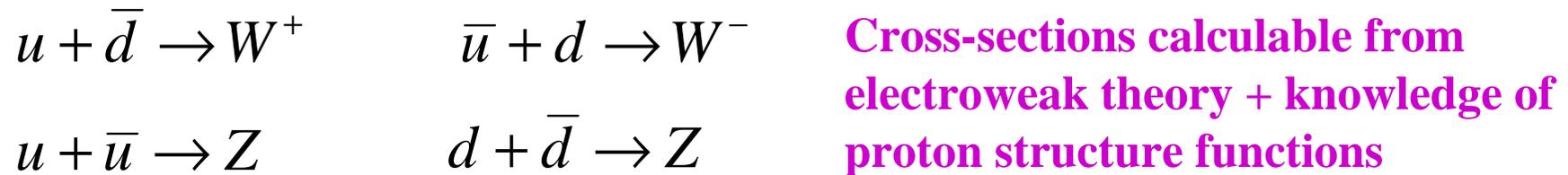
Department of Physics, University of Wisconsin, Madison, Wisconsin 53706

Presented by C. Rubbia

Abstract: We outline a scheme of searching for the massive weak boson ($M = 50 - 200 \text{ GeV}/c^2$). An antiproton source is added either to the Fermilab or the CERN SPS machines to transform a conventional 400 GeV accelerator into a $p\bar{p}$ colliding beam facility with 800 GeV in the center of mass ($E_{\text{eq}} = 320,000 \text{ GeV}$). Reliable estimates of production cross sections along with a high luminosity make the scheme feasible.



Dominant W and Z production processes at a proton – antiproton collider:



▪ **Energy requirements:**

proton (antiproton) momentum at high energies is carried by gluons (~ 50%) and valence quarks (antiquarks) (~ 50%)

On average: quark momentum $\approx \frac{1}{6}$ (proton momentum)

→ collider energy $\approx 6 \times$ boson mass $\approx 500 - 600$ GeV

▪ **Luminosity requirements:**

Inclusive cross-section for $\bar{p} + p \rightarrow Z + \text{anything}$ at ~ 600 GeV: $\sigma \approx 1.6$ nb

Branching ratio for $Z \rightarrow e^+ e^-$ decay $\approx 3\%$

$$\sigma(\bar{p}p \rightarrow Z \rightarrow e^+ e^-) \approx 50 \text{ pb} = 5 \times 10^{-35} \text{ cm}^2$$

Event rate = $L \sigma$ [s^{-1}] (L \equiv luminosity)

→ 1 event / day $\Rightarrow L \approx 2.5 \times 10^{29} \text{ cm}^2 \text{ s}^{-1}$

CERN accelerators in 1976

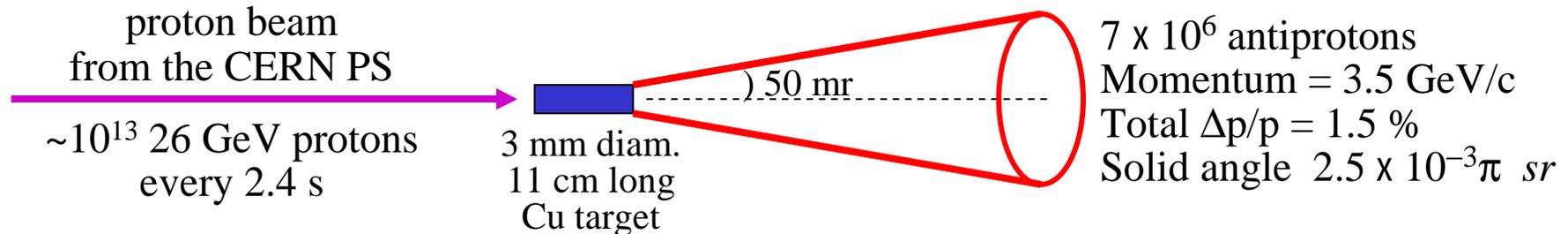
- 26 GeV proton synchrotron (PS) in operation since 1959
- 450 GeV proton synchrotron (SPS) just starting operation



A view of the CERN SPS

To achieve luminosities $\geq 10^{29} \text{ cm}^{-2} \text{ s}^{-1}$ need an antiproton source capable of delivering once per day $3 \times 10^{10} \bar{p}$ distributed into few (3 – 6) tightly collimated bunches within the angular and momentum acceptance of the SPS

Antiproton production:



Number of antiprotons / PS cycle OK

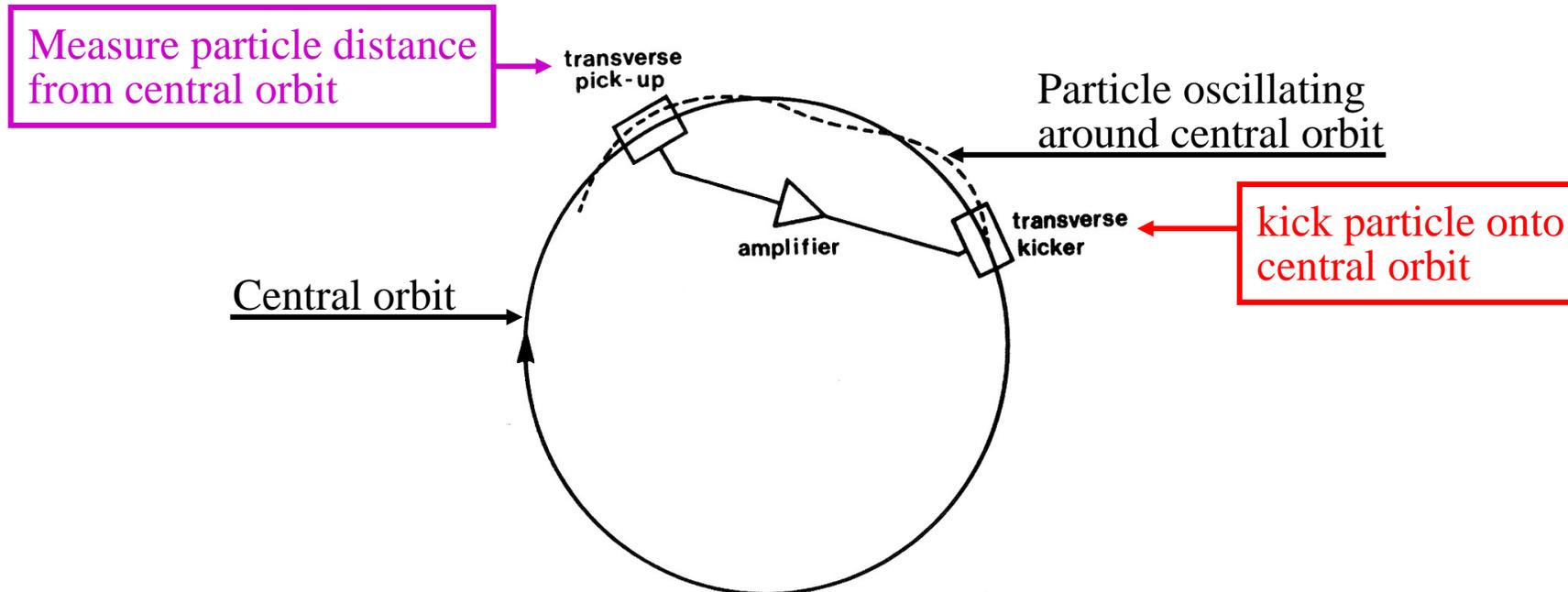
but phase space volume too large by a factor $\geq 10^8$ to fit into SPS acceptance even after acceleration to the injection energy of 26 GeV

→ must increase the antiproton phase space density by $\geq 10^8$ before sending them to the SPS (“cooling”)

“Stochastic” cooling

(invented at CERN by Simon van der Meer in 1972)

Example: cooling of the horizontal motion



In practice, the pick-up system measures the average distance from central orbit of a group of particles (depending on frequency response)

Independent pick-up – kicker systems to cool:

- horizontal motion
- vertical motion
- longitudinal motion (decrease of $\Delta p/p$)
(signal from pick-up system proportional to Δp)

The CERN Antiproton Accumulator (AA)

3.5 GeV/c large-aperture ring for antiproton storage and cooling



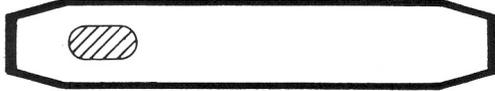
(during construction)

AA operation

Section of the AA vacuum chamber



The first pulse of $7 \times 10^6 \bar{p}$ has been injected



Precooling reduces momentum spread



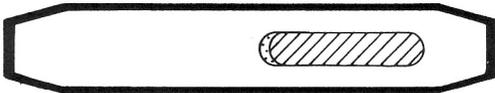
First pulse is moved to the stack region where cooling continues



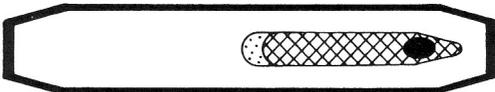
Injection of 2nd \bar{p} pulse 2.4 s later



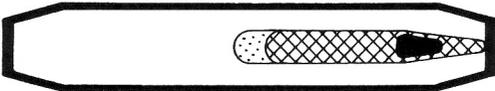
After precooling 2nd pulse is also stacked



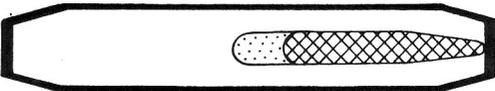
After 15 pulses the stack contains $10^8 \bar{p}$



After one hour a dense core has formed inside the stack



After one day the core contains enough \bar{p} 's for transfer to the SPS

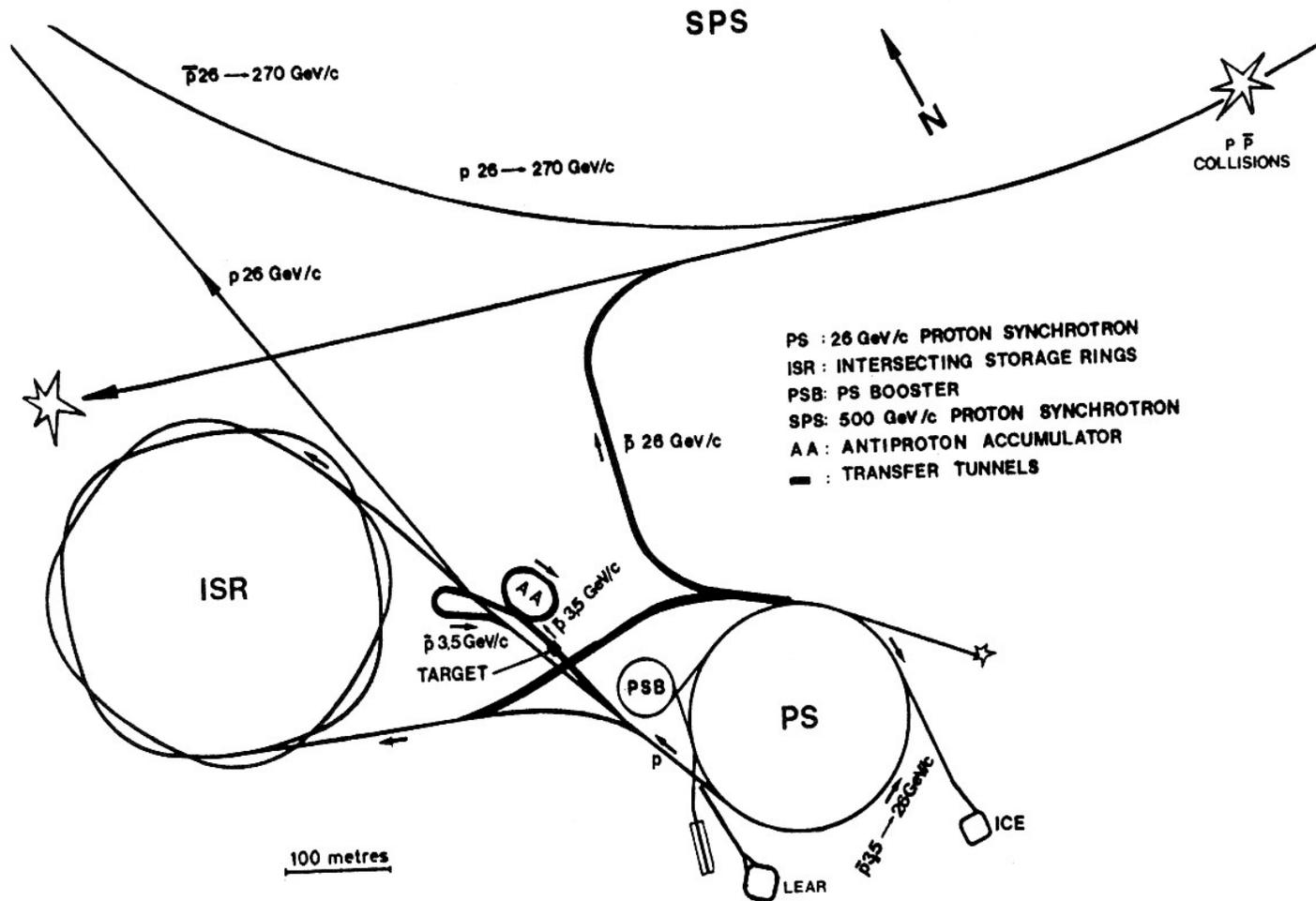


The remaining \bar{p} 's are used for next day accumulation



\bar{p} momentum

Sketch of the CERN accelerators in the early 1980's



1986 – 90: add another ring (“Antiproton Collector” AC)
around the AA – larger acceptance for single \bar{p} pulses
(7×10^7 p / pulse \Rightarrow \sim tenfold increase of stacking rate)



Proton – antiproton collider operation, 1981 - 90

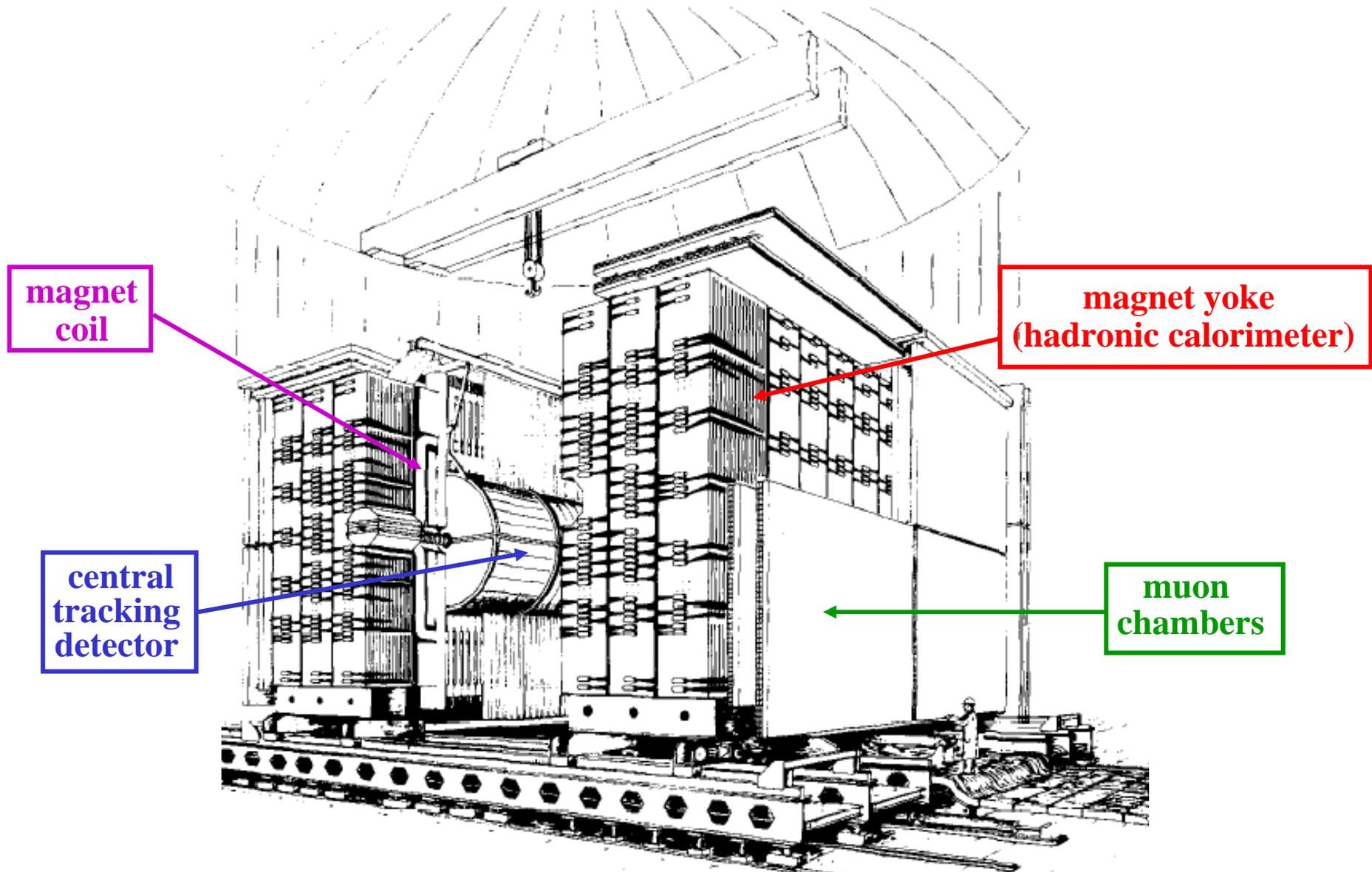
Year	Collision Energy (GeV)	Peak luminosity ($\text{cm}^{-2} \text{s}^{-1}$)	Integrated luminosity (cm^{-2})
1981	546	$\sim 10^{27}$	2.0×10^{32}
1982	546	5×10^{28}	2.8×10^{34}
1983	546	1.7×10^{29}	1.5×10^{35}
1984-85	630	3.9×10^{29}	1.0×10^{36}
1987-90	630	$\sim 2 \times 10^{30}$	1.6×10^{37}

← **W discovery**

← **Z discovery**

1991: end of collider operation

UA1 detector



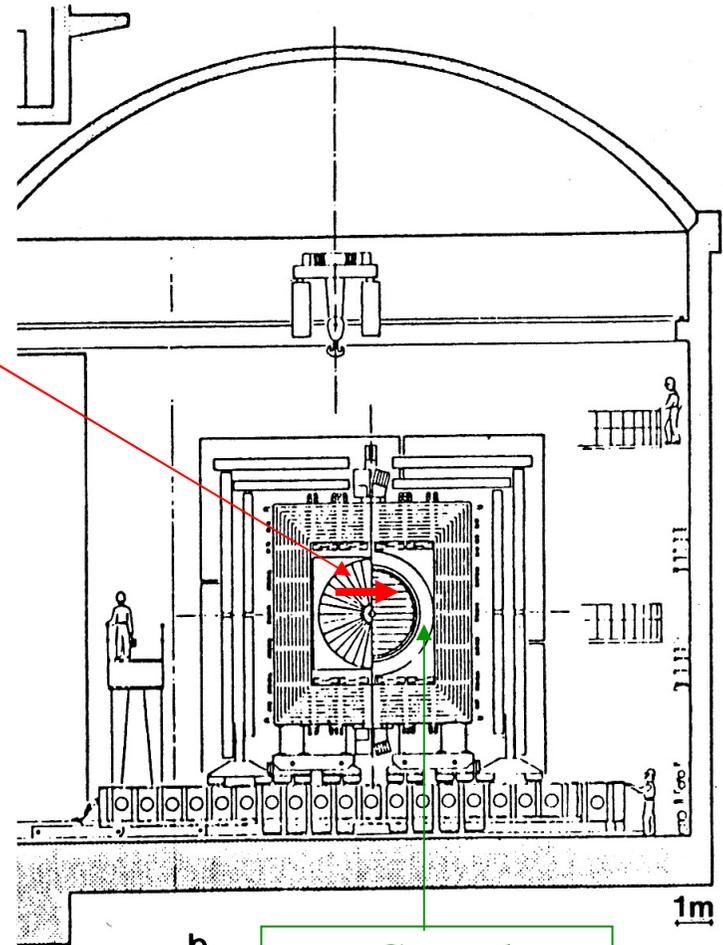
(shown with the two halves of the dipole magnet opened up)

View normal to beam axis



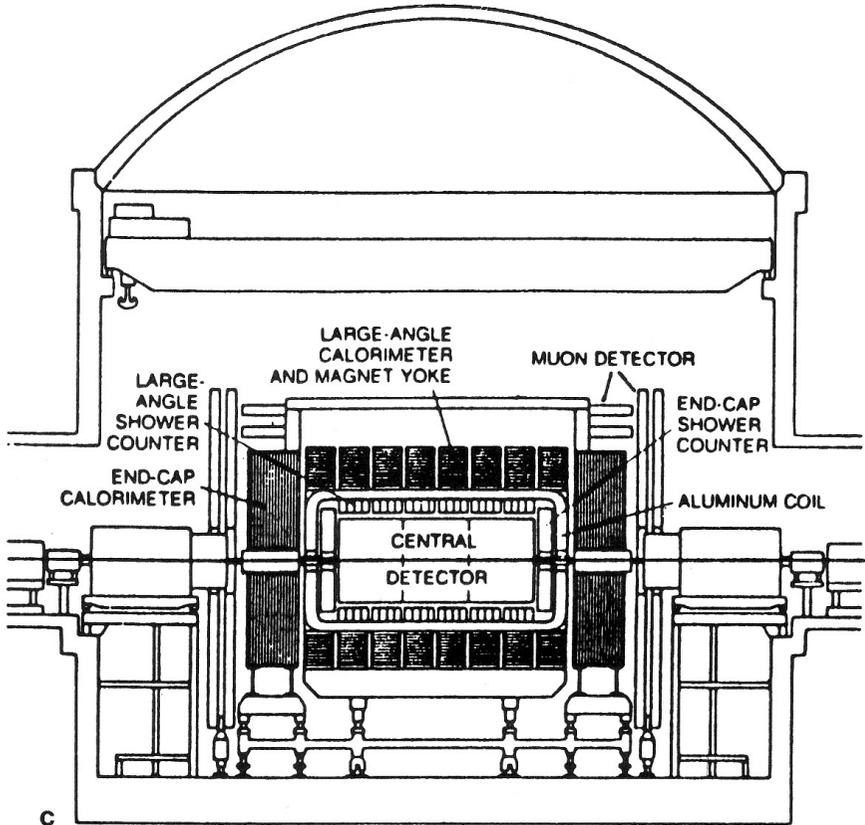
Magnetic field direction

Vertical section along beam axis



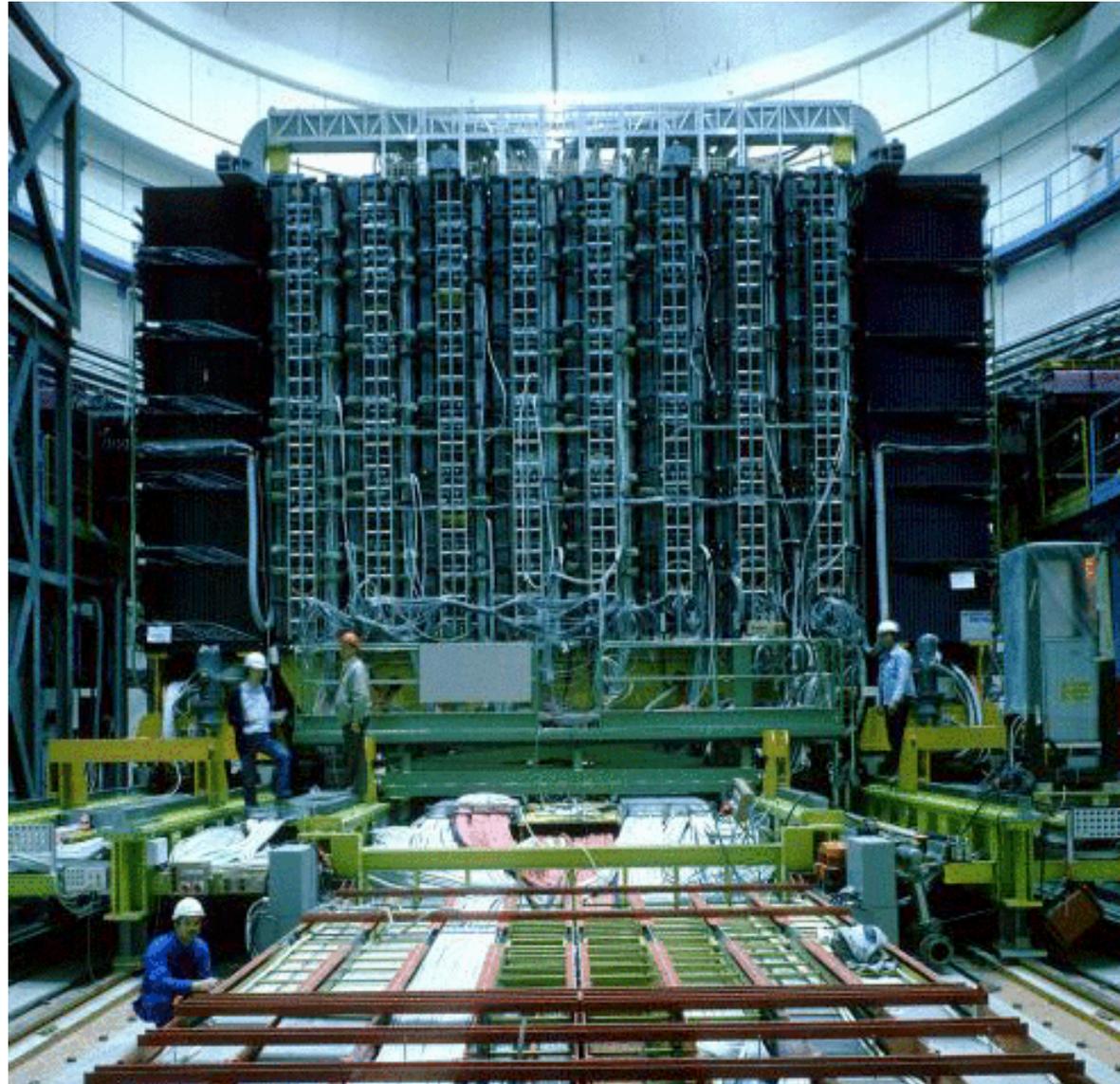
b

Central electromagnetic calorimeter

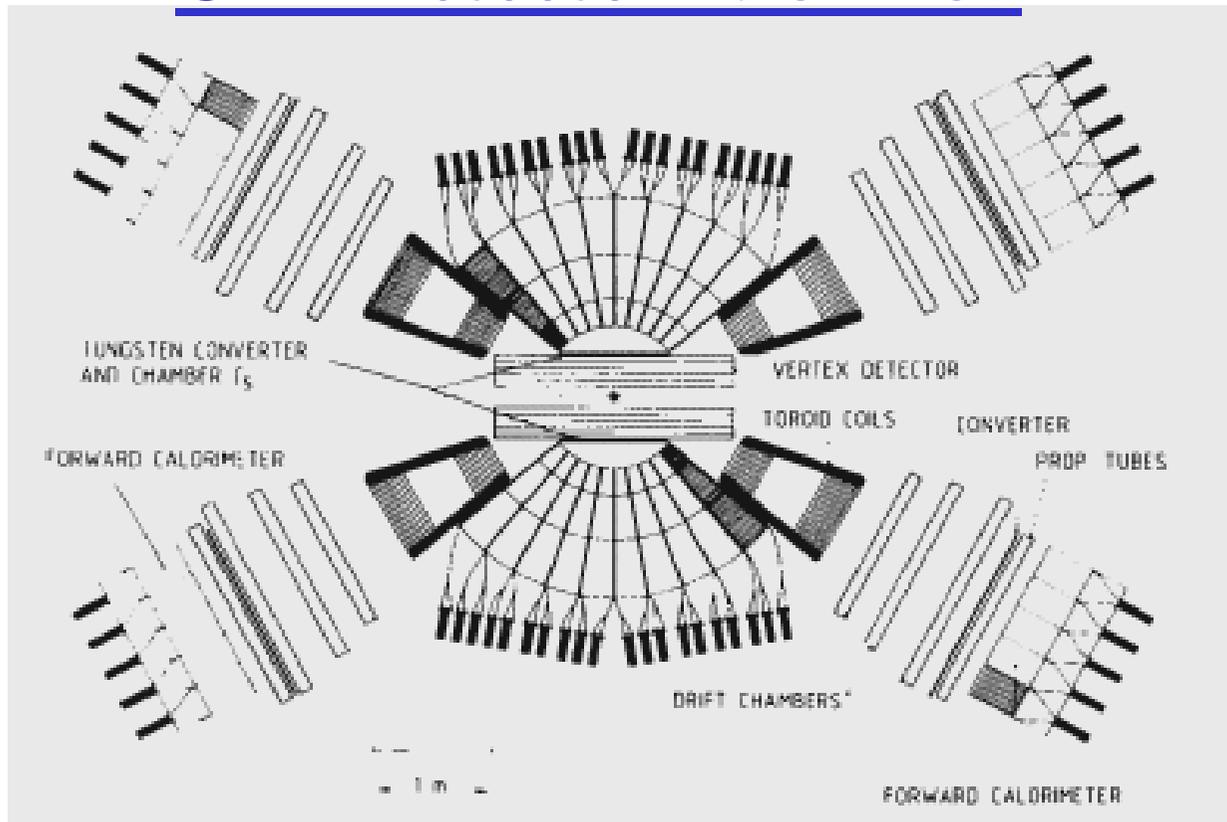


c

UA1 detector during assembly



UA2 Detector 1981 - 85

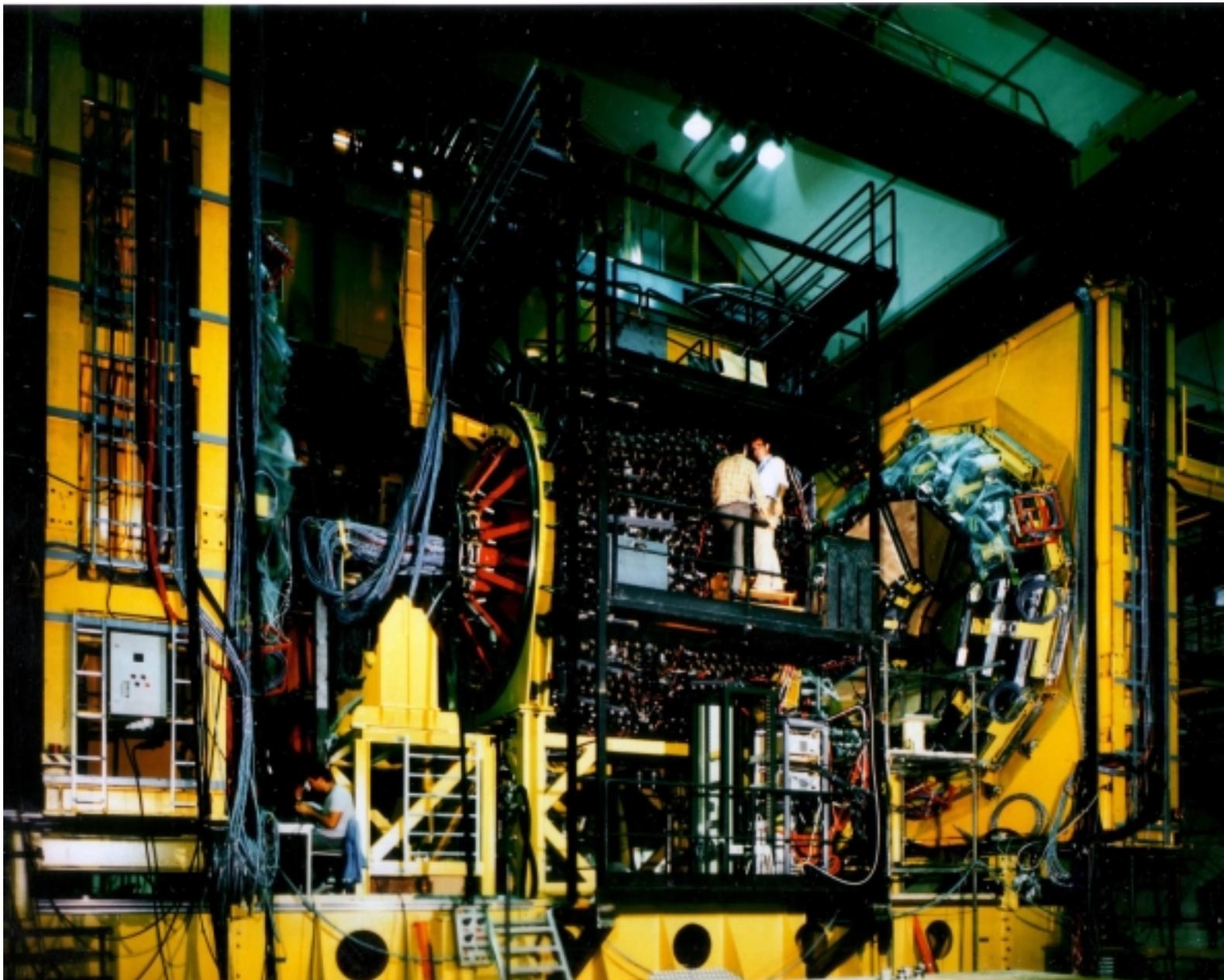


**Central region: tracking detector (“vertex detector”);
“pre-shower” detector
electromagnetic and hadronic calorimeters;
no magnetic field**

**20° – 40° regions : toroidal magnetic field;
tracking detectors;
“pre-shower” detector + electromagnetic calorimeter.**

No muon detector

UA2 detector during assembly



W discovery

Dominant decay mode (~70%) $W \rightarrow q \bar{q}' \rightarrow$ two hadronic jets overwhelmed by two-jet background from QCD processes

\Rightarrow search for leptonic decays:

$$W^+ \rightarrow e^+ + \nu_e$$

(UA1, UA2)

$$W^+ \rightarrow \mu^+ + \nu_\mu$$

(UA1 only) (and charge-conjugate decays)

Expected signal from $W \rightarrow e \nu$ decay:

- large transverse momentum (p_T) isolated electron
- p_T distribution peaks at $m_W / 2$ (“Jacobian peak”)
- large missing transverse momentum from the undetected neutrino

(W produced by quark-antiquark annihilation, e.g. $u + \bar{d} \rightarrow W^+$, is almost collinear with beam axis; decay electron and neutrino emitted at large angles to beam axis have large p_T)

NOTE

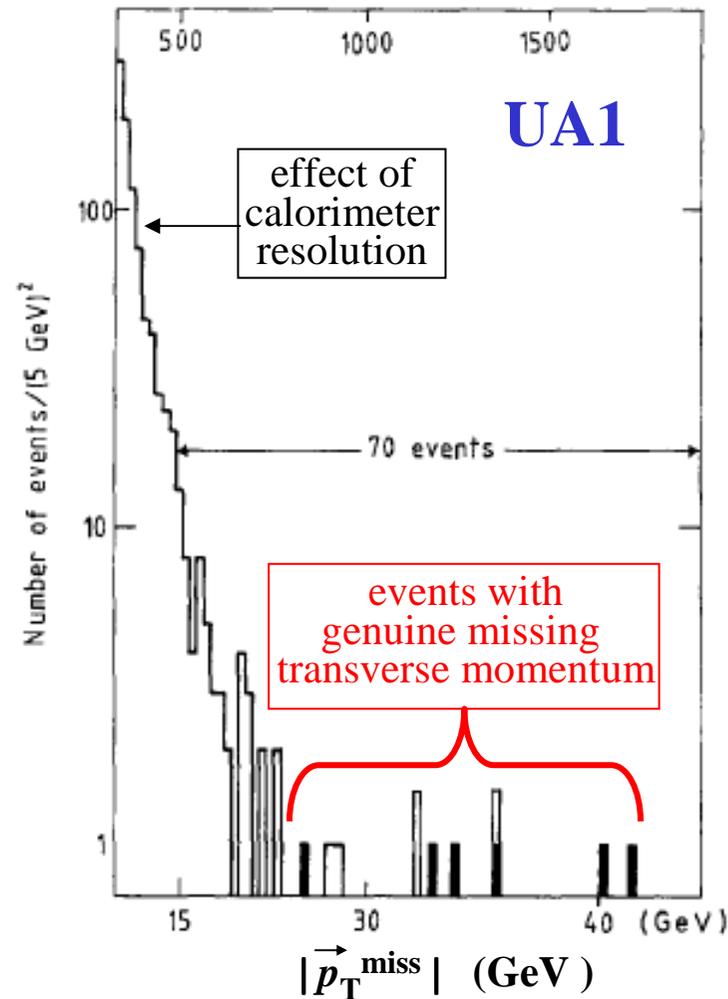
Missing longitudinal momentum cannot be measured at hadron colliders because of large number of high-energy secondary particles emitted at very small angles inside the machine vacuum pipe

Missing transverse momentum (\vec{p}_T^{miss})

- Associate momentum vector \vec{p} to each calorimeter cell with energy deposition > 0
- Direction of \vec{p} from event vertex to cell centre
- $|\vec{p}| =$ energy deposited in cell
- Definition:

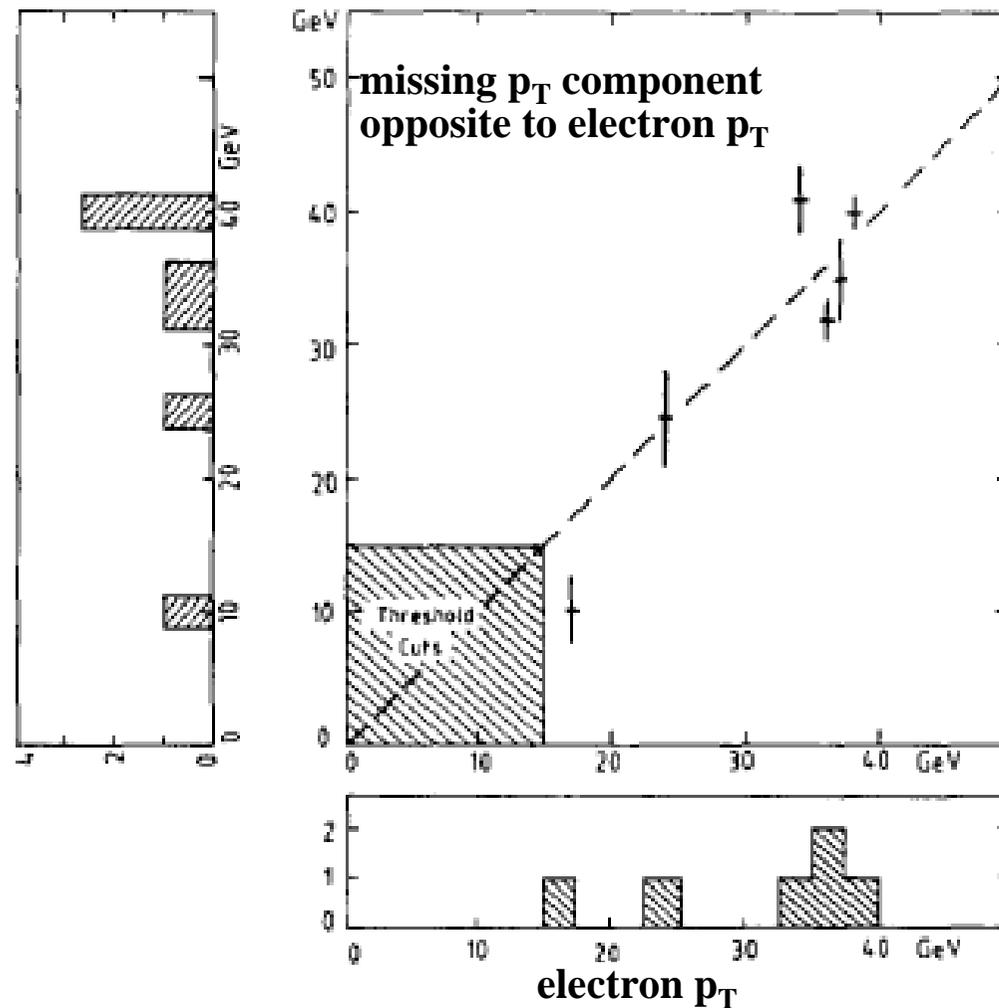
$$\vec{p}_T^{\text{miss}} + \sum_{\text{cells}} \vec{p}_T = 0$$

(momentum conservation in plane perpendicular to beam axis)



■ Six events containing a large p_T electron

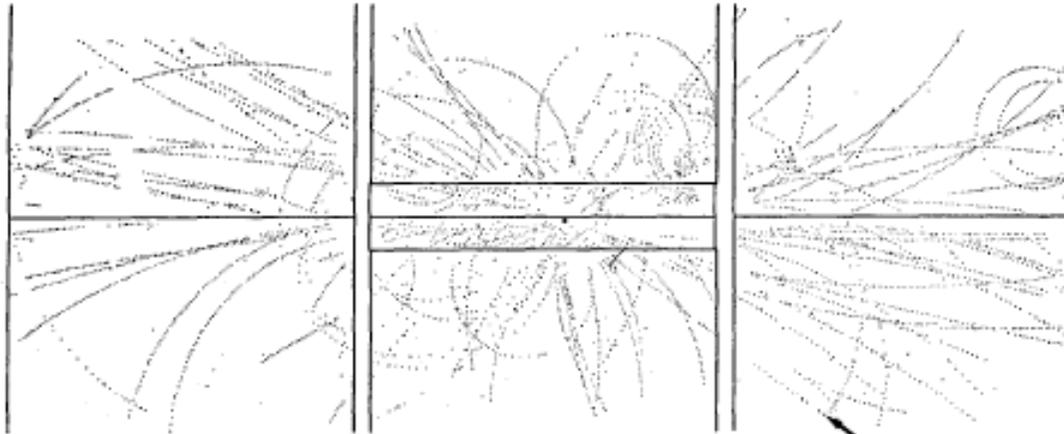
UA1: correlation between electron p_T and missing p_T



Six events with large p_T electron and large missing p_T opposite to electron p_T consistent with $W \rightarrow e \nu$ decay (result announced at a CERN seminar on January 20, 1983)

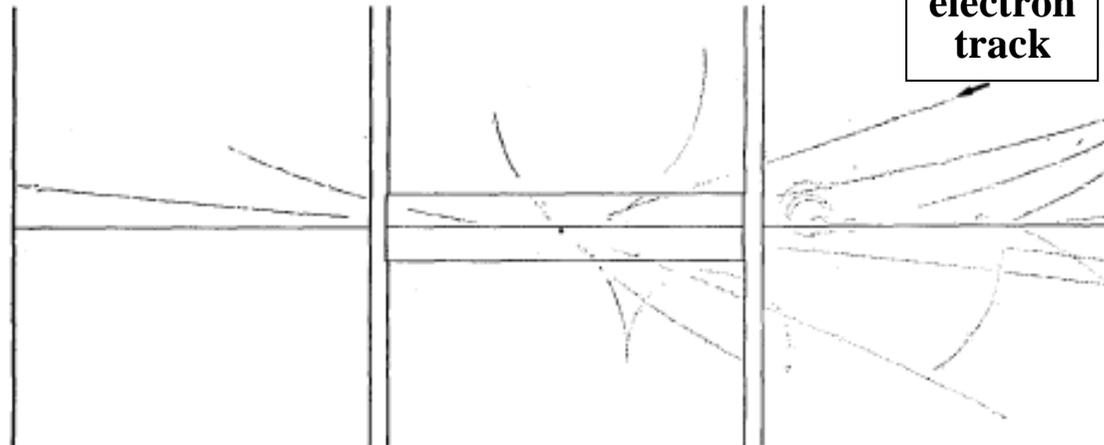
Two UA1 $W \rightarrow e \nu$ events

EVENT 2958. 1279.



electron track

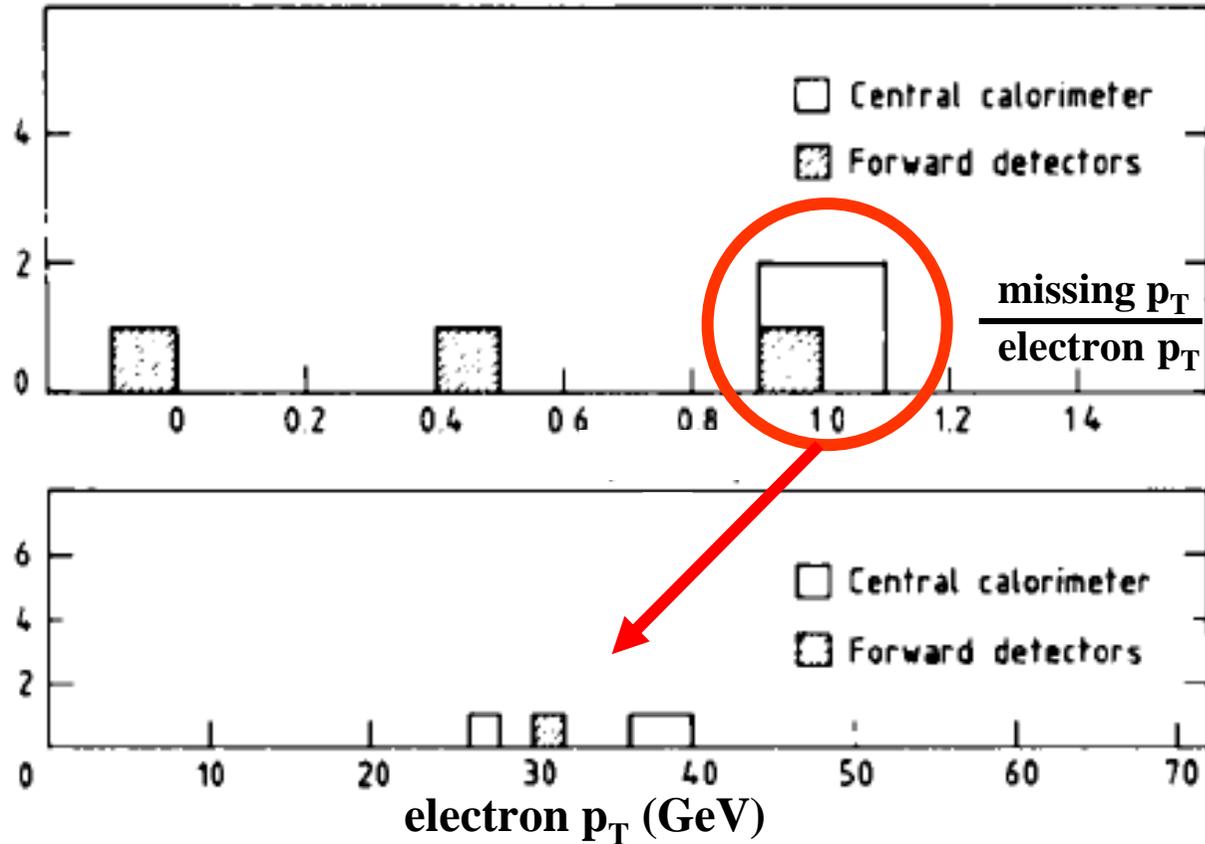
EVENT 4017. 838.



electron track

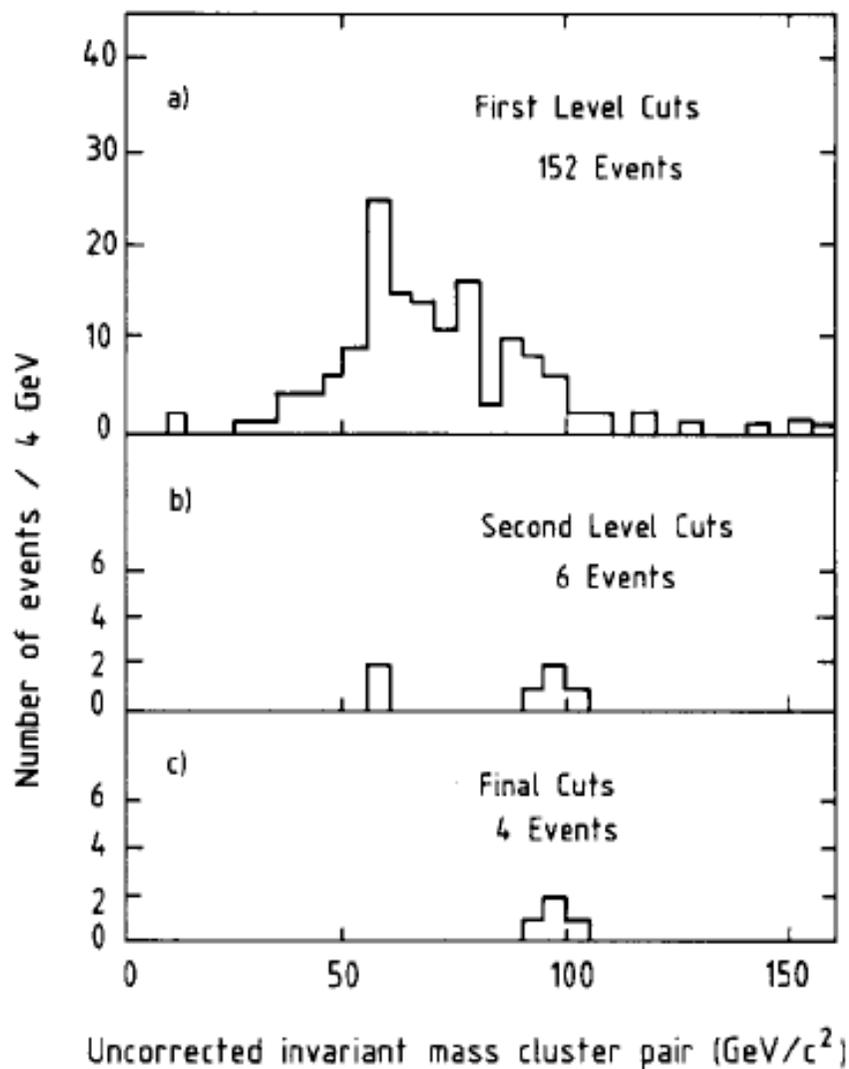
UA2: results presented at a CERN seminar on January 21, 1983

Six events containing an electron with $p_T > 15$ GeV



UA1: observation of $Z \rightarrow e^+ e^-$

(May 1983)

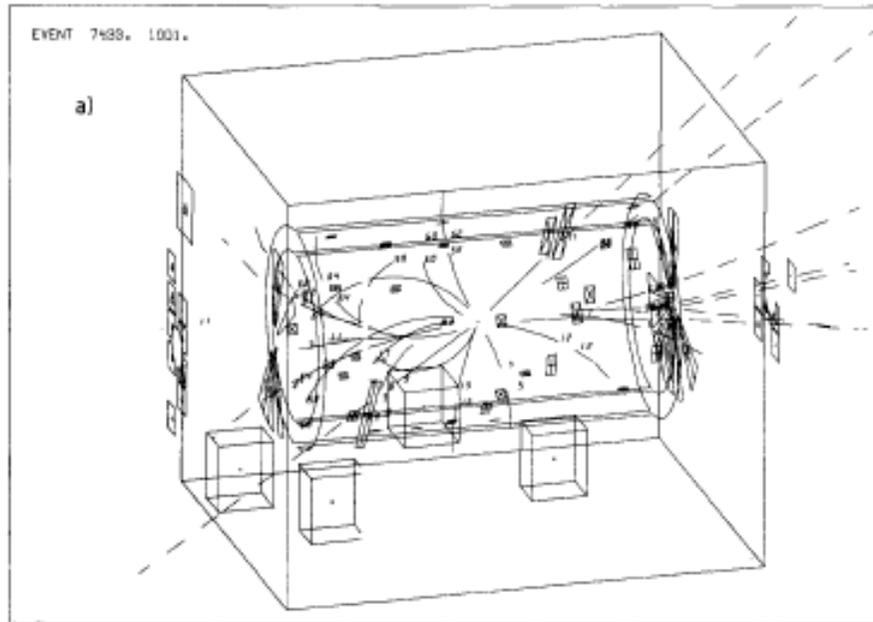


**Two energy clusters ($p_T > 25$ GeV)
in electromagnetic calorimeters;
energy leakage in hadronic calorimeters
consistent with electrons**

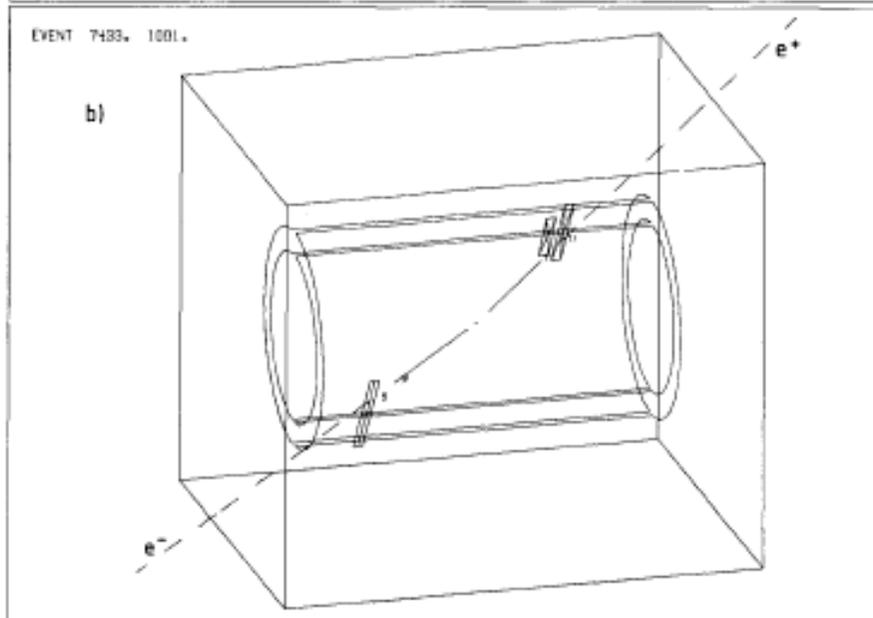
**Isolated track with $p_T > 7$ GeV
pointing to at least one cluster**

**Isolated track with $p_T > 7$ GeV
pointing to both clusters**

UA1 $Z \rightarrow e^+ e^-$ event

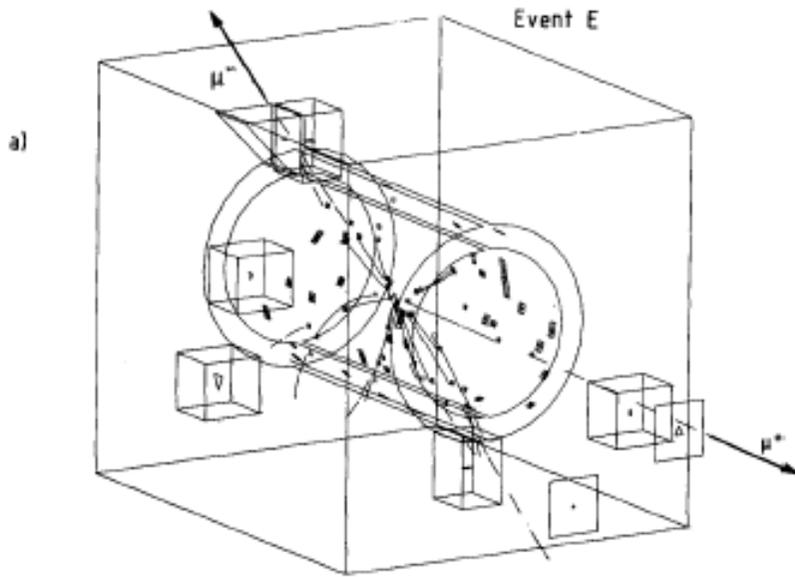


**Display of all reconstructed tracks
and calorimeter hits**



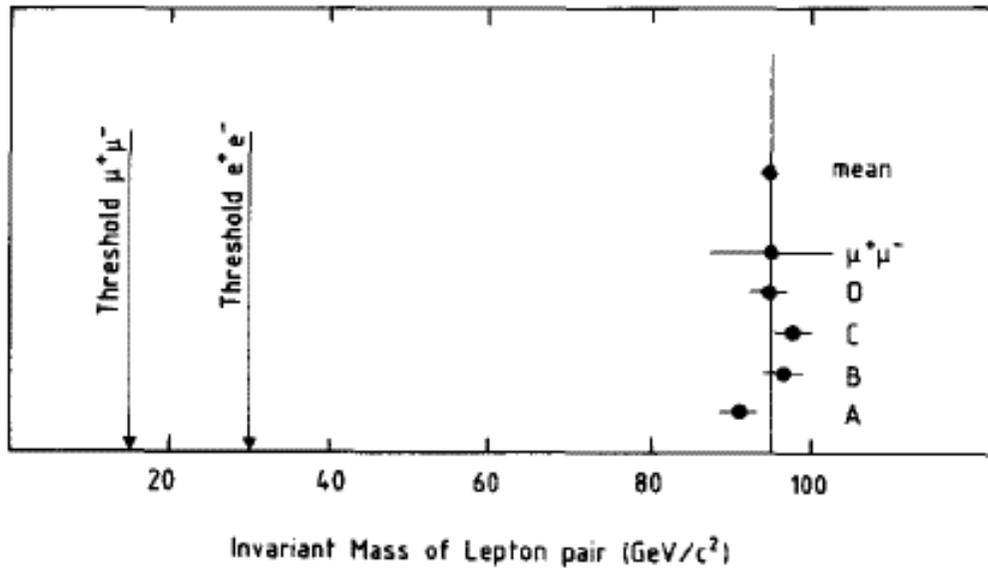
**Display of tracks and calorimeter hits
with $p_T > 2$ GeV**

EVENT 6600. 222.



**UA1 $Z \rightarrow \mu^+ \mu^-$ event
(May 1983)**

**The only $\mu^+ \mu^-$ pair observed
during the 1983 collider run**



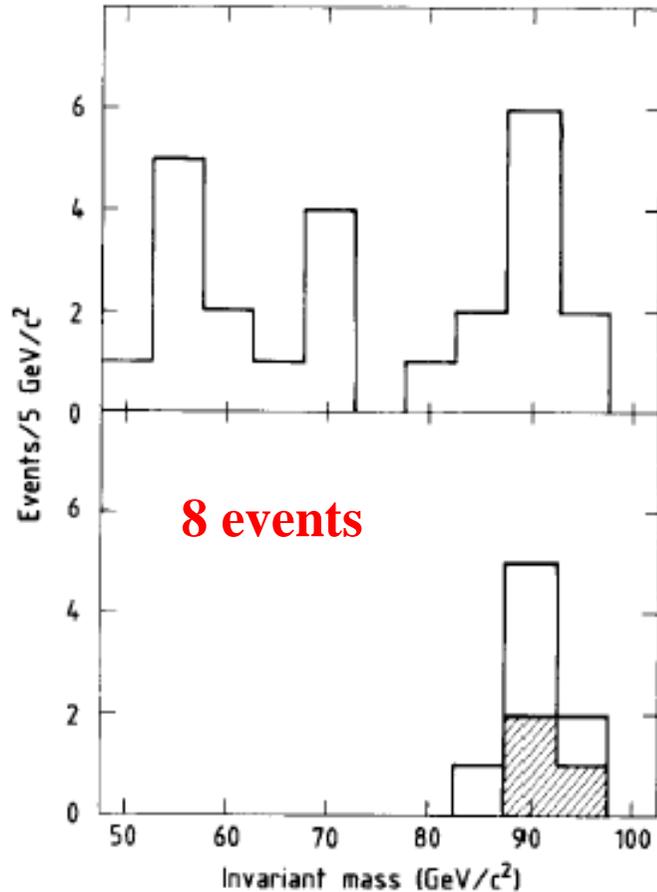
**UA1:
all lepton pairs
from the 1983 run**

$$m_Z = 95.2 \pm 2.5 \pm 3.0 \text{ GeV}$$

(stat) (syst)

UA2: observation of $Z \rightarrow e^+ e^-$

(June 1983)



Two energy clusters with $p_T > 25$ GeV
in electromagnetic calorimeters;
energy leakage in hadronic calorimeters
consistent with electrons

A track identified as an isolated electron
pointing to at least one of the two clusters

 Track identified as an isolated electron
pointing to both energy clusters

$$m_Z = 91.9 \pm 1.3 \pm 1.4 \text{ GeV}$$

(stat) (syst)



The Nobel Prize in Physics 1984

"For their decisive contributions to the large project, which led to the discovery of the field particles W and Z, communicators of weak interaction."



Carlo Rubbia

1/2 of the prize

Italy

CERN
Geneva, Switzerland
b. 1934



Simon van der Meer

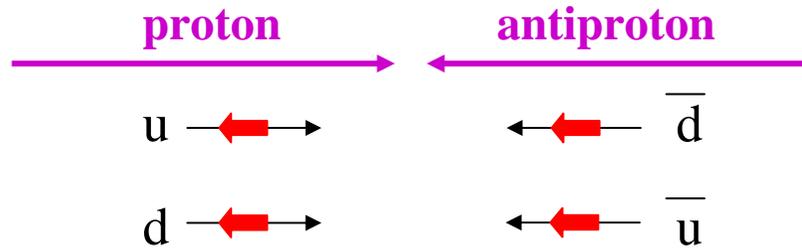
1/2 of the prize

the Netherlands

CERN
Geneva, Switzerland
b. 1925

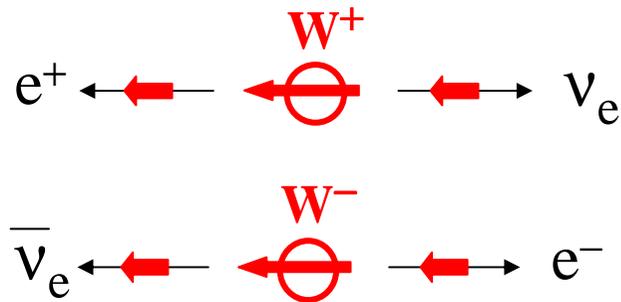


Charge asymmetry in $W \rightarrow e \nu$ decay



W^\pm polarization along antiproton direction
(consequence of V - A coupling)

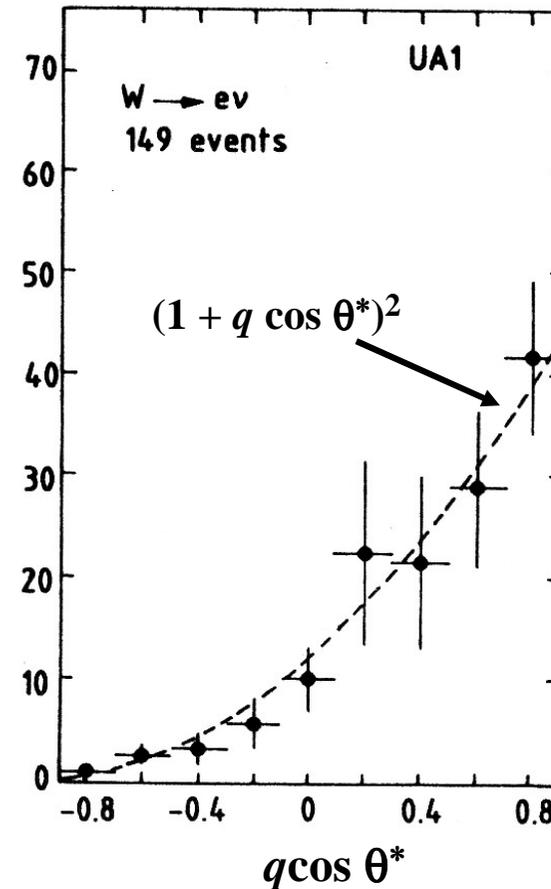
In the W rest frame:



Electron (positron) angular distribution:

$$\frac{dn}{d \cos \theta^*} \propto (1 + q \cos \theta^*)^2$$

$q = +1$ for positrons; $q = -1$ for electrons
 $\theta^* = 0$ along antiproton direction

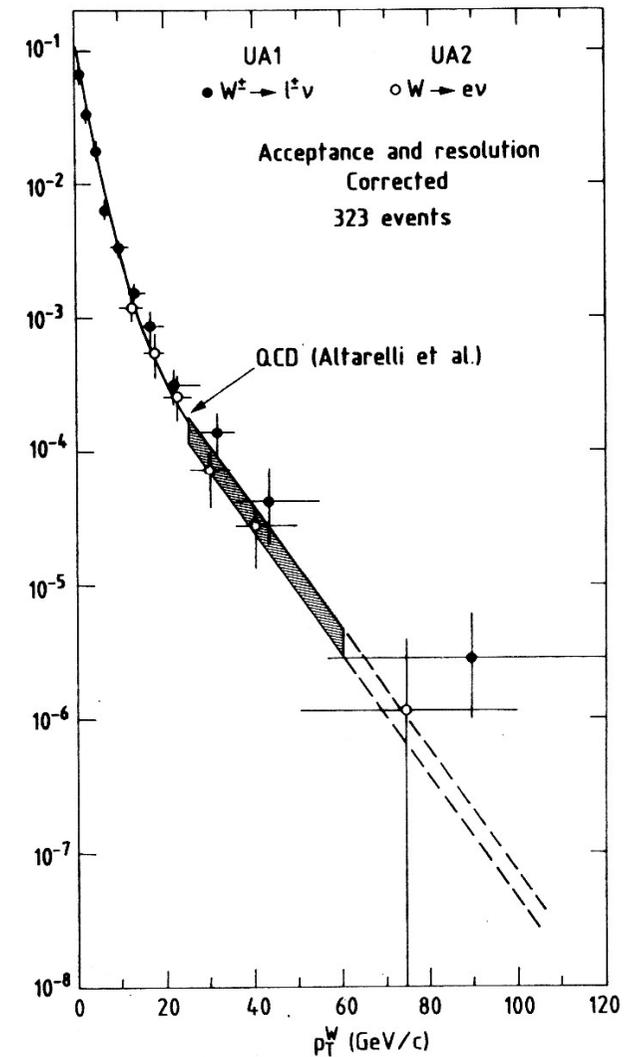


W transverse momentum (\vec{p}_T^W)

- $p_T^W \neq 0$ because of initial-state gluon radiation
- \vec{p}_T^W equal and opposite to total transverse momentum carried by all hadrons produced in the same collision:

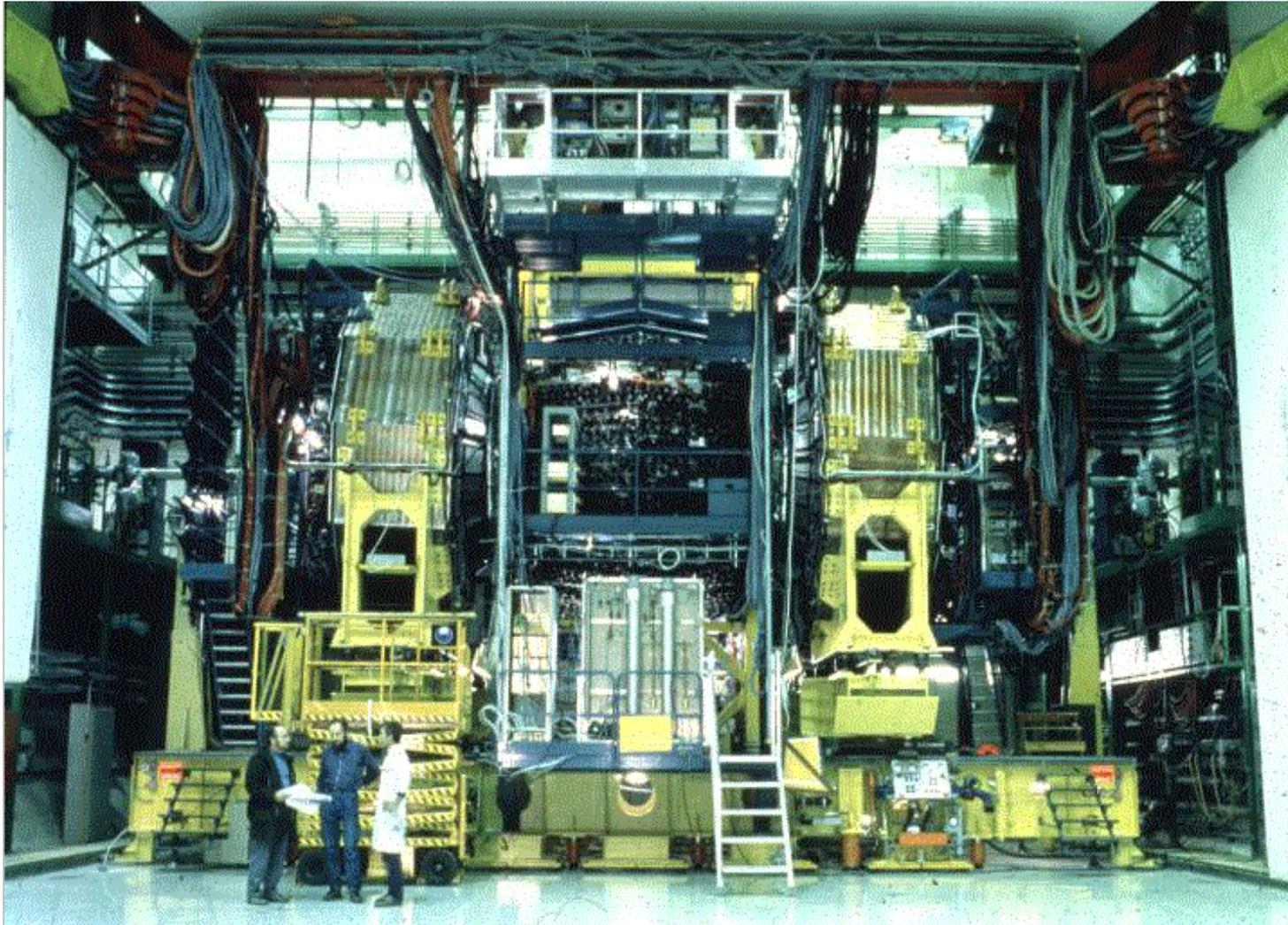
$$\vec{p}_T^W = - \sum_{hadrons} \vec{p}_T$$

- p_T^W distribution can be predicted from QCD



UA2 detector 1987 – 90

- **Tenfold increase of collider luminosity**
- **Full calorimetry down to $\sim 5^\circ \Rightarrow$ improved measurement of missing p_T**
- **No magnetic field, no muon detectors**

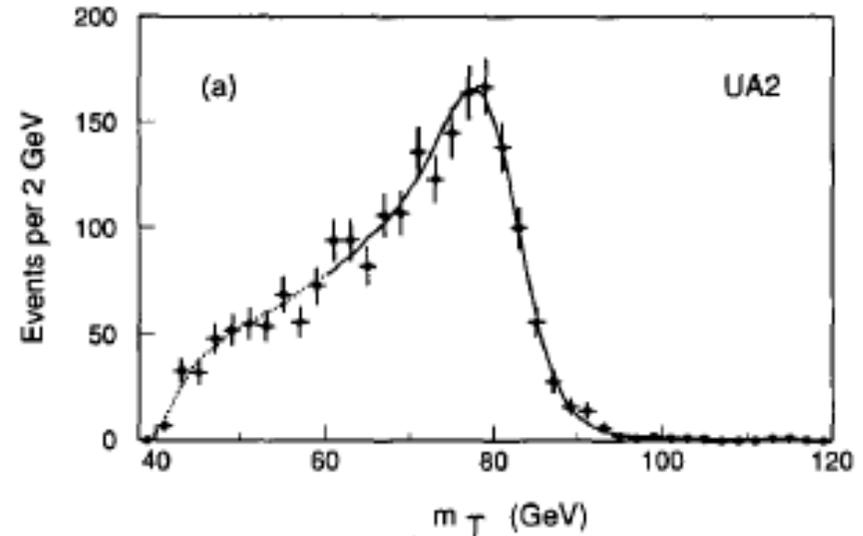


UA2: precise measurement of $\frac{m_W}{m_Z}$

(mass ratio has no uncertainty
from calorimeter calibration)

2065 $W \rightarrow e \nu$ events with the electron
in the central calorimeter ($\theta = 90^\circ \pm 50^\circ$)

Distribution of “transverse mass” m_T
(m_T : invariant mass using only the
e and ν momentum components
normal to beam axis – the longitudinal
component of the ν momentum cannot
be measured at hadron colliders)



Fit of the distribution with m_W as fitting parameter:

$$m_W = 80.84 \pm 0.22 \text{ GeV}$$

Two samples of $Z \rightarrow e^+ e^-$ events :

- both electrons in central calorimeter (95 events)

$$m_Z = 91.65 \pm 0.34 \text{ GeV}$$

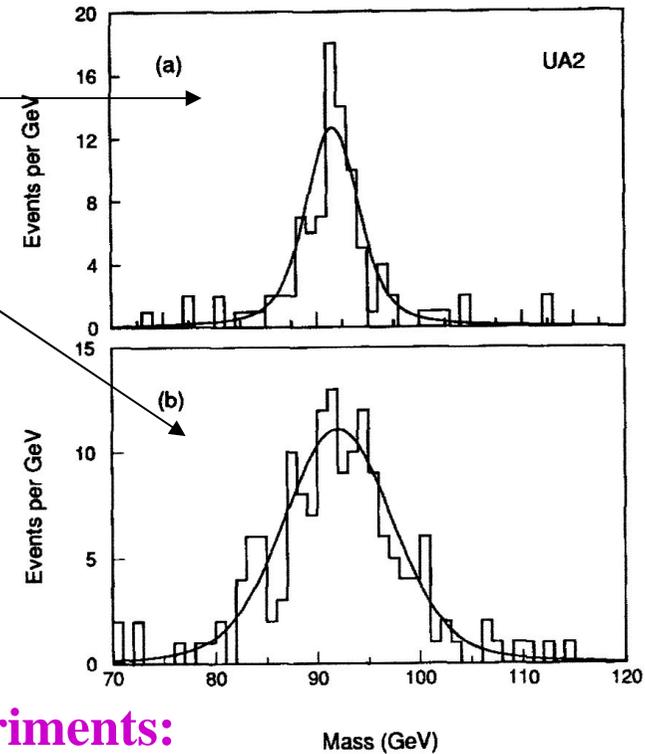
- only one electron in central calorimeter (156 events)

$$m_Z = 92.10 \pm 0.48 \text{ GeV}$$

- combined samples:

$$m_Z = 91.74 \pm 0.28 \text{ GeV}$$

→ $\frac{m_W}{m_Z} = 0.8813 \pm 0.0036 \pm 0.0019$



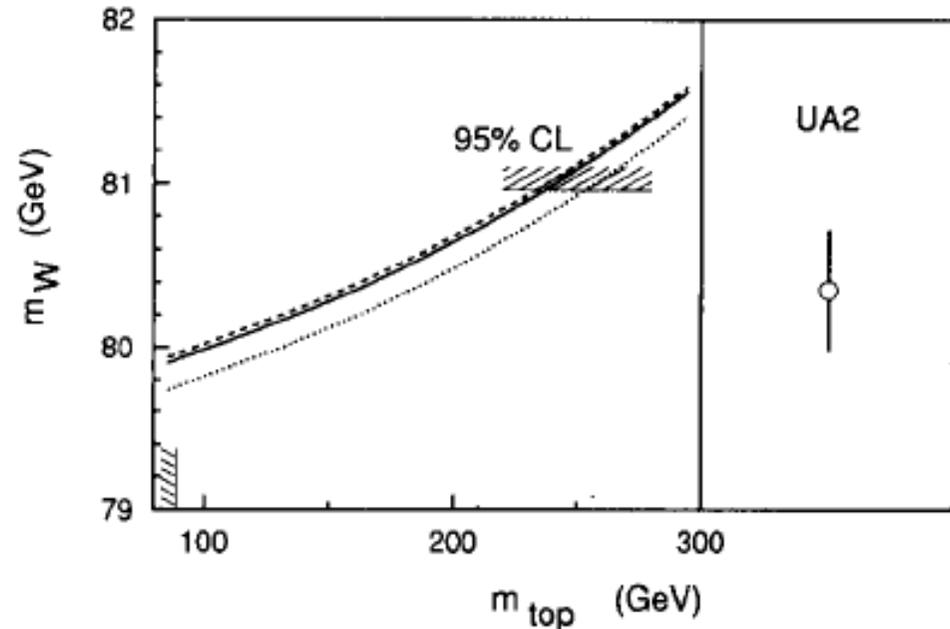
Using precise measurement of m_Z from LEP experiments:

$$m_W = 80.35 \pm 0.33 \pm 0.17 \text{ GeV}$$

- bounds on the mass of the top quark in the frame of the Standard Model:**

$$m_{top} = 160_{-60}^{+50} \text{ GeV}$$

(five years before the top quark discovery at Fermilab)



Jet production at high transverse momentum

Proton – antiproton collisions at high energy \equiv collisions between two broad-band beams of quarks, antiquarks, gluons (“partons”)

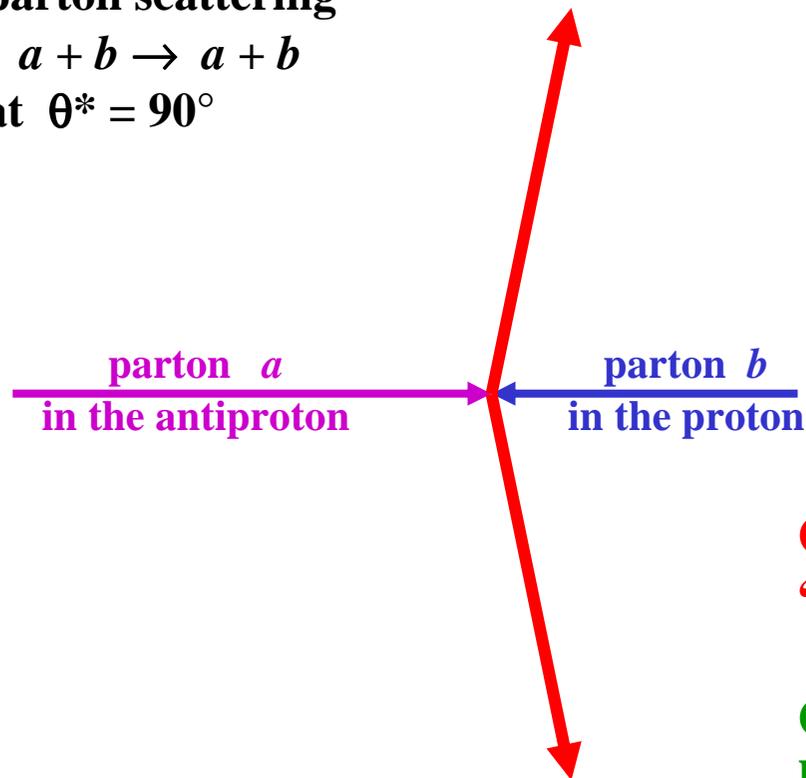
Parton – parton scattering at large angles \Rightarrow two hadronic jets at large p_T

EXAMPLE:

parton scattering

$$a + b \rightarrow a + b$$

at $\theta^* = 90^\circ$



Elementary processes :

$$\bar{q} + q \rightarrow \bar{q} + q$$

$$\bar{q} + g \rightarrow \bar{q} + g$$

$$g + q \rightarrow g + q$$

$$g + g \rightarrow g + g$$

$$\bar{q} + q \rightarrow g + g$$

$$g + g \rightarrow \bar{q} + q$$

Outgoing partons \Rightarrow hadronic jets
“Hadronization”: an effect of
parton confinement

Outgoing jets are COPLANAR ($\Delta\phi = 180^\circ$)
but not COLLINEAR

UA2

- For each calorimeter cell define “transverse energy”:

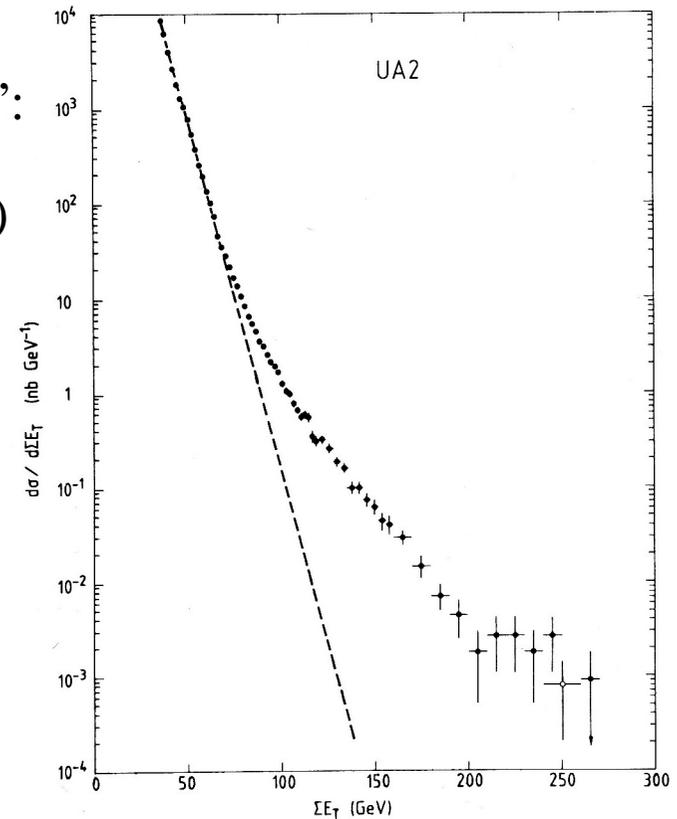
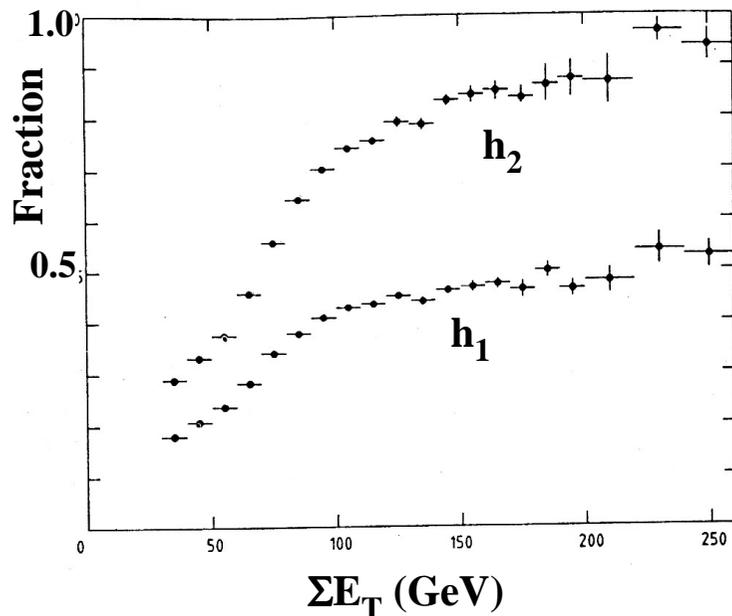
$$\mathbf{E}_T = \mathbf{E} \sin \theta$$

(\mathbf{E} : energy deposition; θ : polar angle of cell centre)

- On-line selection of events (“trigger”):

$$\Sigma \mathbf{E}_T > \text{threshold} \quad (\text{sum over all calorimeter cells})$$

- Build “clusters”: groups of adjacent cells, each with $\mathbf{E}_T > 0.4 \text{ GeV}$
- Define cluster transverse energy by adding \mathbf{E}_T over all cells in cluster
- Order clusters: $\mathbf{E}_T^1 > \mathbf{E}_T^2 > \dots > \mathbf{E}_T^n$

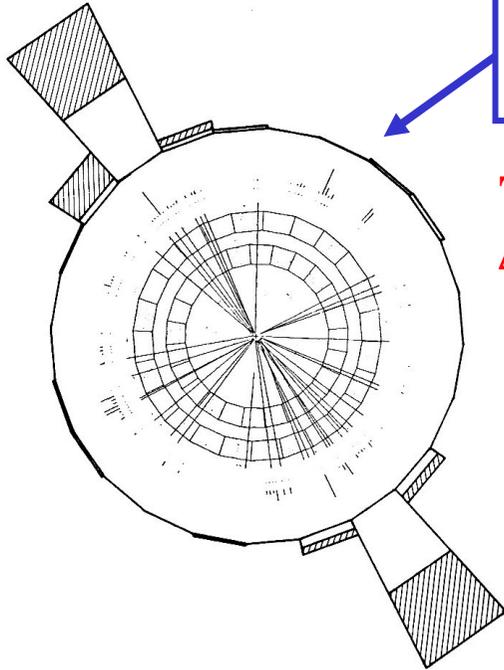


$$h_1 = \frac{\mathbf{E}_T^1}{\Sigma \mathbf{E}_T} \quad \text{fraction of total } \mathbf{E}_T \text{ carried by leading cluster}$$

$$h_2 = \frac{\mathbf{E}_T^1 + \mathbf{E}_T^2}{\Sigma \mathbf{E}_T} \quad \text{fraction of total } \mathbf{E}_T \text{ carried by the two leading clusters}$$

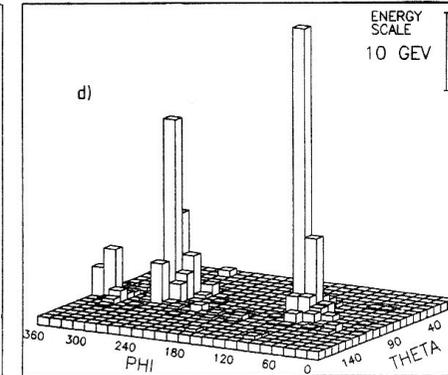
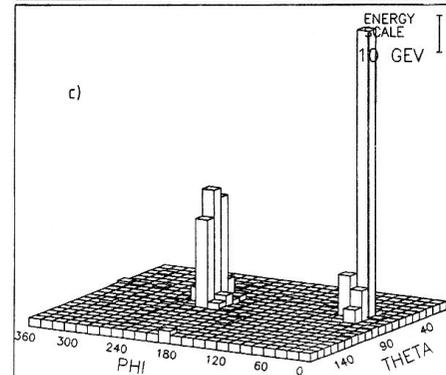
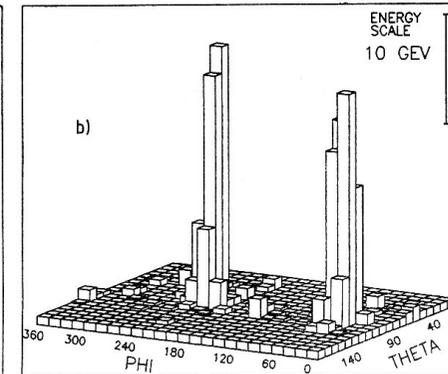
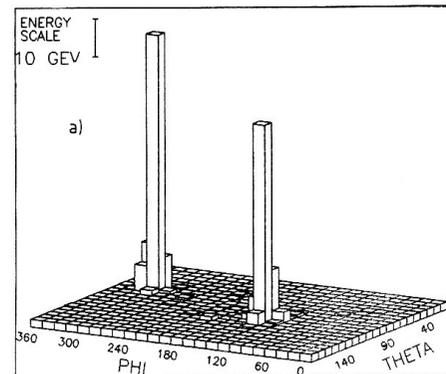
**For large values of ΣE_T ,
the total transverse energy
is mainly concentrated into two clusters**

View of a typical event with large total transverse energy in a plane perpendicular to the beam axis



The two leading clusters consist of a small number of cells $\Delta\phi \approx 180^\circ$ as expected for two-jet production

Transverse energy distribution
in the $\phi - \theta$ plane
for four typical events
with large total transverse energy



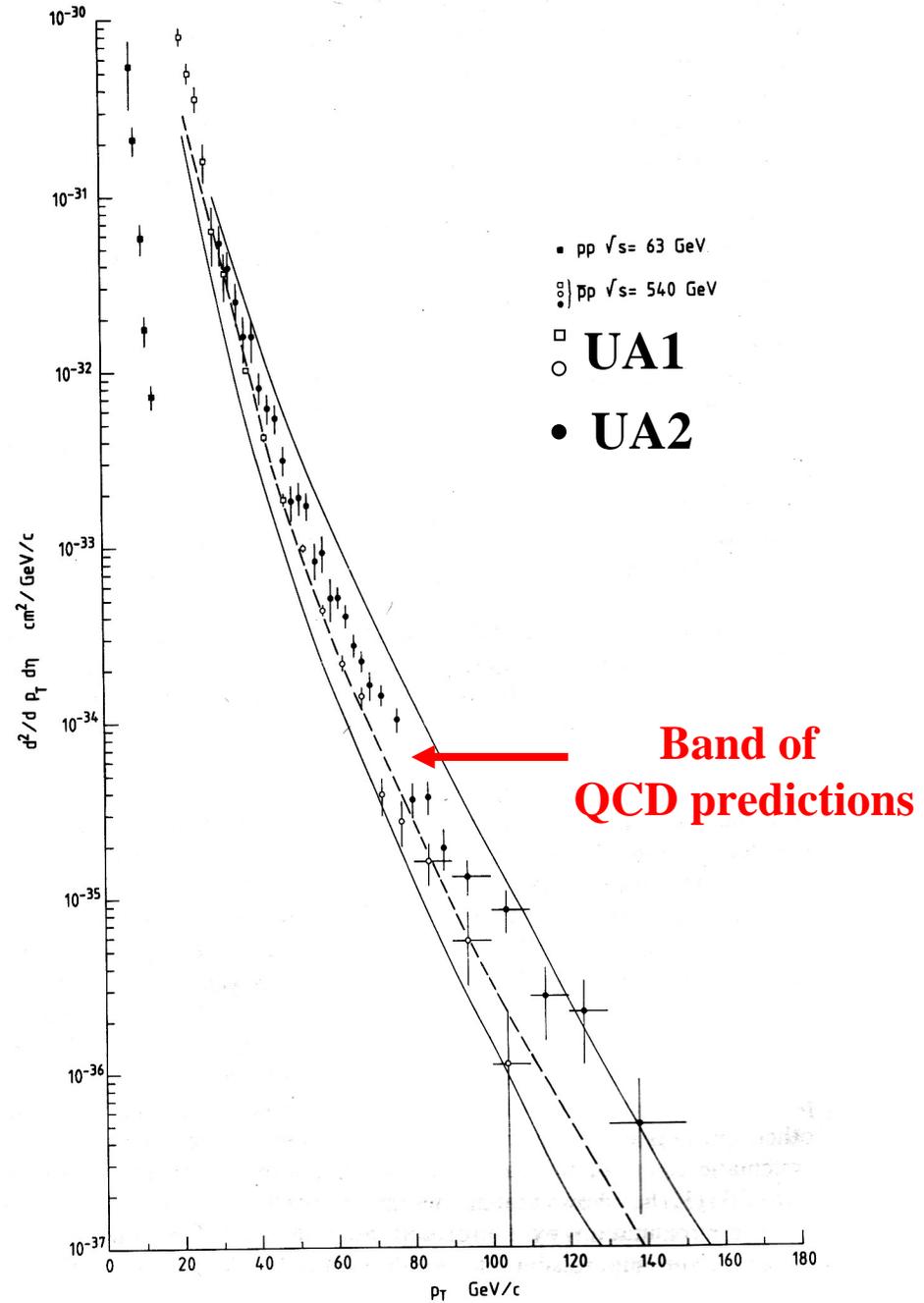
Differential cross-section for inclusive jet production at 90°

Relative contributions of
parton scattering processes:

$$\bar{q} + q \rightarrow \bar{q} + q \quad 1.0$$

$$\left. \begin{array}{l} \bar{q} + g \rightarrow \bar{q} + g \\ g + q \rightarrow g + q \end{array} \right\} 1.2$$

$$g + g \rightarrow g + g \quad 6.0$$



Angular distribution of parton-parton scattering

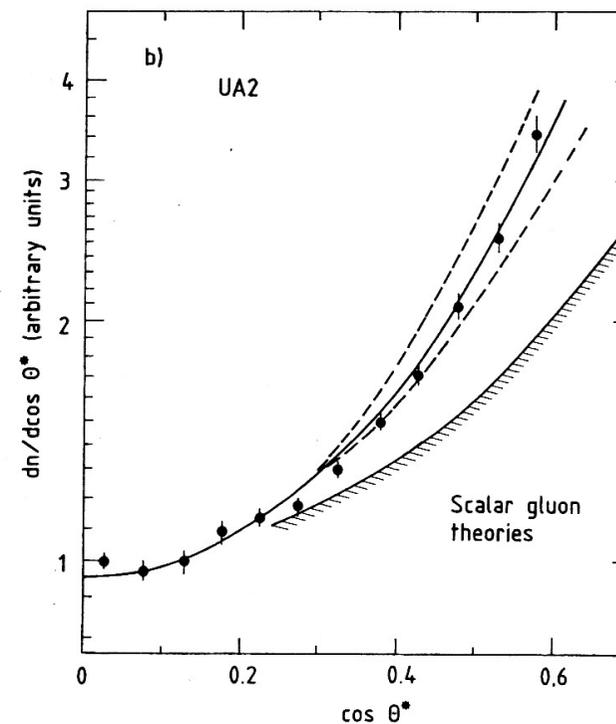
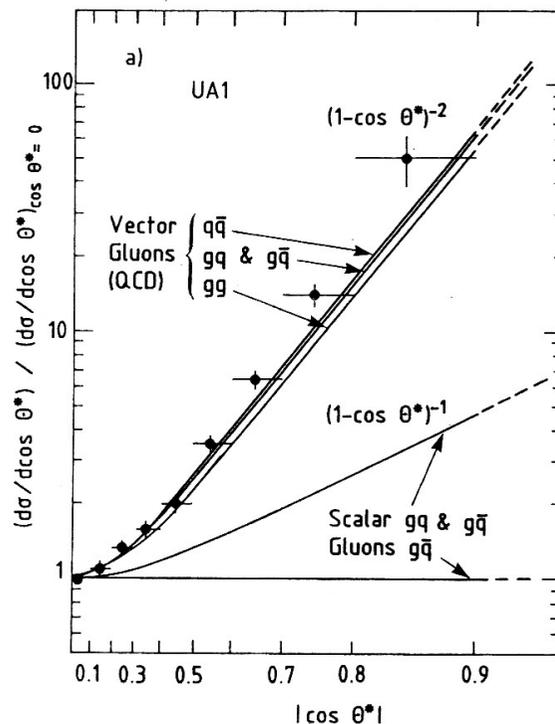
Transform the two jets to the centre-of-mass of the two-jet system

⇒ can measure the parton-parton scattering angle θ^*

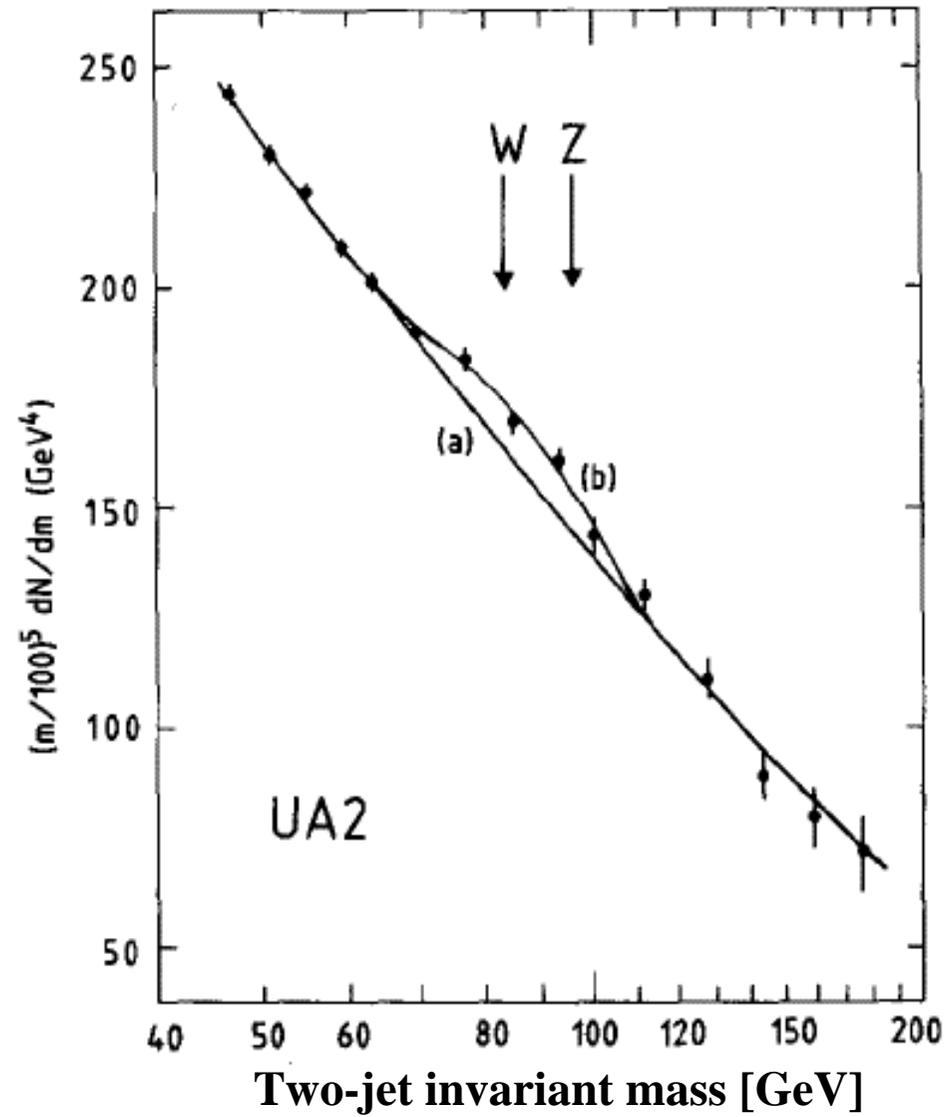
(without distinguishing between θ^* and $\pi - \theta^*$)

All elementary parton scattering processes are dominated by gluon exchange in the t – channel ⇒ expect “Rutherford formula” for spin-1 gluons:

$$\frac{d\sigma}{d\cos\theta^*} \propto \frac{1}{\sin^4(\theta^*/2)} \propto \frac{1}{(1-\cos\theta^*)^2} \quad (\text{as for one-photon exchange})$$

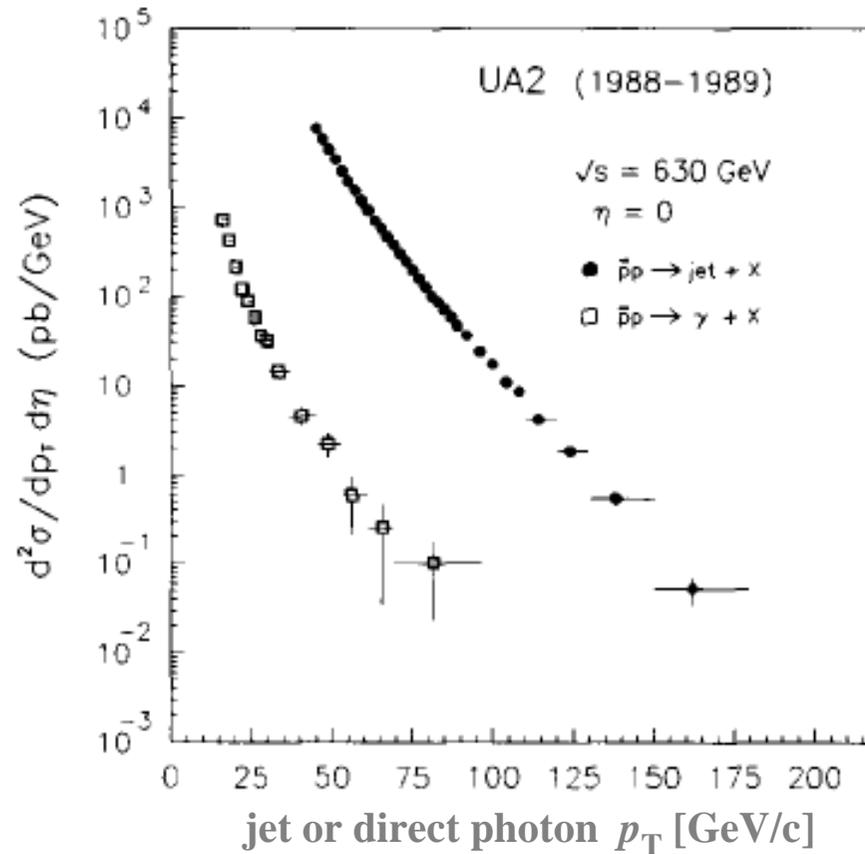
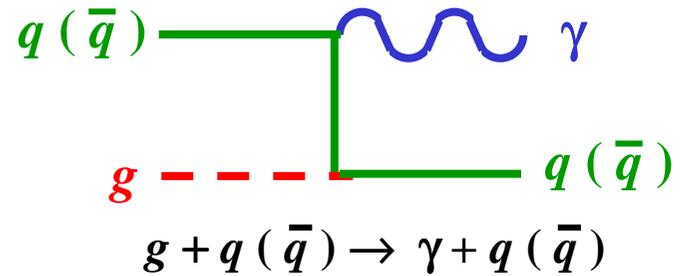
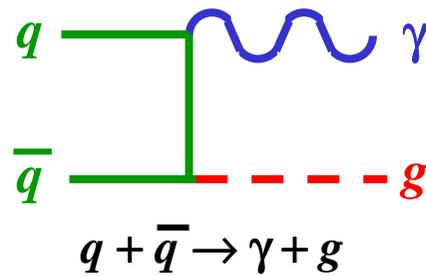


UA2: search for $W^\pm \rightarrow q \bar{q}'$ and $Z \rightarrow q \bar{q} \Rightarrow 2$ jets



UA2: comparison of inclusive jet and direct photon production

Diagrams for direct photon production:



CONCLUSIONS

The CERN Proton – Antiproton Collider:

**initially conceived as an experiment to detect the W^\pm and Z bosons;
in the end, a general – purpose accelerator facility exploring hadron collisions
at centre-of-mass energies an order of magnitude larger than those previously available.**

Among the main physics results:

- **W^\pm and Z detection and studies (tests of the electroweak theory)**
- **study of hadronic jets and photons at high p_T (tests of perturbative QCD)**
- **heavy flavour physics (first indirect evidence of $B^0 - \bar{B}^0$ mixing)**

**The prevailing opinion before the first operation of the CERN $\bar{p} p$ Collider:
proton – proton (and antiproton – proton) collisions are “DIRTY”, “COMPLICATED”
and “DIFFICULT TO INTERPRET”**

**The physics results (and those from the Fermilab $\bar{p} p$ collider at 1.8 TeV) have shown
that this pessimistic view is wrong if the experiments are designed to look at the
basic “physics building blocks”:**

- **hadronic jets at large p_T (representing quarks, antiquarks, gluons)**
- **leptons**
- **photons**
- **missing transverse momentum (neutrinos, other possible weakly interacting particles)**

**THE SUCCESS OF THE CERN PROTON – ANTIPROTON COLLIDER
HAS OPENED THE ROAD TO THE LHC**